

Neutral Fermion Excitations in the Moore-Read State at Filling Factor $\nu = 5/2$

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We present evidence supporting the weakly paired Moore-Read phase in the half-filled second Landau level, focusing on some of the qualitative features of its excitations. Based on numerical studies, we show that systems with odd particle number at the flux $N_\phi = 2N - 3$ can be interpreted as a neutral fermion mode of one unpaired fermion, which is gapped. The mode is found to have two distinct minima, providing a signature that could be observed by photoluminescence. In the presence of two quasiparticles the same neutral fermion excitation is shown to be gapless, confirming expectations for non-Abelian statistics of the Ising model with degenerate fusion channels 1 and ψ .

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Previous studies of the $\nu = 5/2$ quantum Hall effect [1] have accumulated mounting evidence in favor of the Moore-Read state [2–4] of weakly paired composite fermions (CFs) [5]. Because of its potential application in topological quantum computation [6], scrutinizing the physical realization of this phase at $\nu = 5/2$ is a great challenge of fundamental and technological importance [6–8]. Theoretically, the main evidence comes from adiabatic connection [9,10] and significant overlaps of the exact ground state in finite model systems with the Moore-Read (MR) state [4] or more general weakly paired wave functions [9]. The key expectation is that, by extension, the quasiparticles of this gapped topological phase are described by the same underlying (Ising) conformal field theory and obey non-Abelian statistics.

In this Letter, we provide an approach that tests the *qualitative* properties of the $\nu = 5/2$ quasiparticles directly, using numerical analysis in the spherical geometry [11]. Our results support the predictions of Moore and Read [2], without reference to any trial wave function. First, as a consequence of the pairing, it is expected that odd-numbered configurations should be disfavored [3]. We investigate this effect for a selection of simple two-body Hamiltonians such as the Coulomb interaction in the second Landau level (LL). (We focus on qualitative features, and exclude detailed modeling of effects such as finite width [12] or LL mixing [13,14], and assume full spin polarization [15–17].) We show that systems with odd particle number at the flux of the MR ground state possess a dispersing band of low-lying excitations which we interpret as a neutral fermion (NF) mode arising from an unpaired electron. This NF mode has an energy gap Δ_{NF} of the order of the charge gap Δ_c [18]. Second, we study the energetics of a NF in the presence of charged quasiparticles (QPs): positive quasiholes (QHs) or negative quasielectrons (QEs). In this case, our thermodynamic extrapolations of the energy are consistent with a gapless NF. This confirms one of the core features of the non-Abelian statistics of the MR state: the topological degeneracy of two

possible fusion channels 1 and ψ of a pair of two distant QPs, corresponding to the absence or presence of an additional fermion. Furthermore, we determine the NF dispersion and propose an experiment to probe this dispersion directly. We also give evidence that the QHs and QEs of the $\nu = 5/2$ state (for microscopic two-body Hamiltonians) fuse in the ψ channel.

For our studies, we perform exact diagonalizations of model Hamiltonians for $N \leq 20$ spin-polarized, quasi two-dimensional (2D) electrons on a sphere of radius R pierced by $N_\phi = 2N - \sigma$ magnetic flux quanta. The MR Pfaffian state is at the shift of $\sigma = 3$. We consider three model Hamiltonians: First, the Coulomb interaction \mathcal{H}_C in the second LL, as defined by the pseudopotential coefficients either for a 2D layer of effective width $w = 0$ or $w = 3\lambda$ (λ being the magnetic length). Second, a modified Coulomb interaction \mathcal{H}_1 , with the short-range pseudopotential V_1 (for pairs with relative angular momentum $m = 1$) increased by $\delta V_1 = 0.04e^2/\lambda$. This increase in V_1 is known to yield maximum overlap with the MR Pfaffian [9,19], and was found to mimic LL mixing [20] in the perturbative analysis of Bishara-Nayak [14,21]. Third, and finally, the “Pfaffian (Pf) model” given by the projector on triplets with the minimal allowed relative angular momentum ($m = 3$), $\mathcal{H}_{\text{Pf}} = \sum_{i,j,k} P_{ijk}^{(3)}$. Although we focus on the Pfaffian state, the particle-hole symmetry of the two-body interactions \mathcal{H}_C , \mathcal{H}_1 makes our conclusions equally valid for its conjugate, the “anti-Pfaffian.”

We now analyze the excitation spectra of these Hamiltonians. A typical set of raw data is shown in Fig. 1 for $N = 15$ particles. Here we first focus on Figs. 1(d)–1(f), in which we confirm the existence of a dispersive mode associated with a single fermionic QP for finite systems with an odd N [3]. All these three spectra feature a ground state at nonzero angular momentum within a low-lying band of collective excitations. Empirically, we find that this well-separated band extends up to the angular momentum $L = N/2$, as seen most

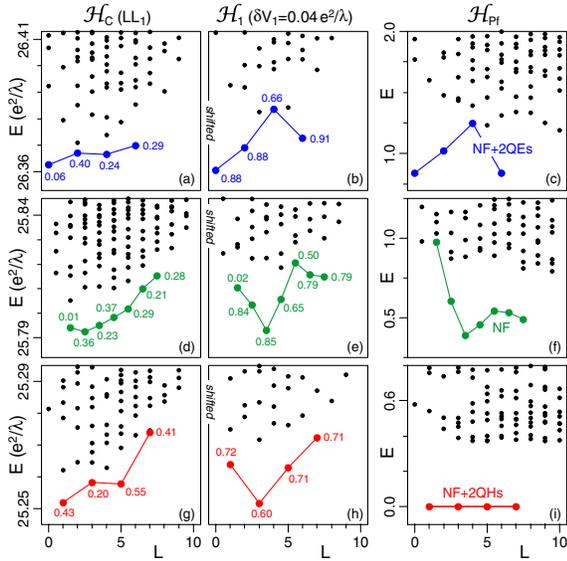


FIG. 1 (color online). Energy spectra (bare interaction energy E versus angular momentum L) of $N = 15$ electrons in a half-filled second LL, interpreted in terms of the neutral fermion (NF) excitation in the Pfaffian (Pf) ground state. Different values of the magnetic flux are $N_\phi = 26$ (top), 27 (center), and 28 (bottom), corresponding to the NF with additional pair of charged quasielectrons (QEs), NF alone, and NF with additional pair of charged quasiholes (QHs), respectively. Different interactions are \mathcal{H}_C = pure Coulomb (left), \mathcal{H}_1 = Coulomb with an additional enhancement of the $m = 1$ pair pseudopotential (center), and \mathcal{H}_{Pf} = three-body Pfaffian Hamiltonian with the only triplet pseudopotential at $m = 3$ (right). Labels indicate squared overlaps $|\langle \mathcal{H}_{C1} | \mathcal{H}_{Pf} \rangle|^2$ for the low-lying bands.

clearly for \mathcal{H}_{Pf} . The minima of the dispersion are far below the continuum. Eigenstates within the band are spaced by $\Delta L = 1$, as expected for a single mobile NF in the background of an underlying quantum liquid.

In Fig. 2 we collected data from energy spectra for different $N \leq 19$ to estimate the NF dispersion, which is compared to the magnetoroton mode in Fig. 2(c). To reduce finite-size effects, for each N we used polynomial interpolation of our data to locate the wave vector k_0 and

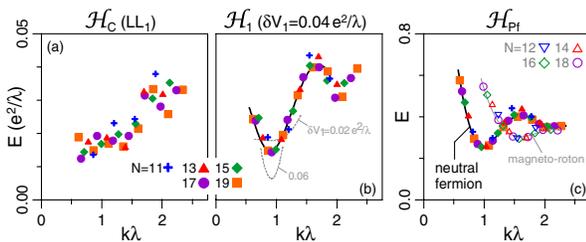


FIG. 2 (color online). Dispersions (energy E versus wave vector k) of the neutral fermion (NF) collective modes of the half-filled second LL, estimated from the systems of $N \leq 19$ electrons at the magnetic flux $N_\phi = 2N - 3$, for the different Hamiltonians of Fig. 1. Gray dashed lines in (b) show the evolution of Δ_{NF} with δV_1 . For comparison, (c) also shows the magnetoroton mode for \mathcal{H}_{Pf} .

the energy E_0 corresponding to the minimum of the dispersion. We find that $k_0(N)$ is essentially constant, and the minimum E_0 is well described by $E_0(N) \simeq \Delta_{NF} + \beta/N$, converging to a finite NF gap Δ_{NF} measured with respect to the MR ground state energy [22]. Subtracting the finite-size scaling, our data reduce to one well-defined curve (most accurately for \mathcal{H}_{Pf}).

The shape of the NF mode varies significantly between our model Hamiltonians. A general feature emerging for all spectra is a NF dispersion with two minima [these are best seen in panel (c)]: a deeper “NF₁” near $k_0\lambda \simeq 1$ and a more shallow “NF₂” near $k\lambda \simeq 2$. The dispersion of \mathcal{H}_C shows strong finite-size effects, which we interpret as a consequence of the proximity to a phase transition into a charge-density wave phase [19]. Indeed, \mathcal{H}_1 which is known to yield a state well inside the weakly paired phase also produces cleanly defined dispersions, particularly for the NF. (The magnetoroton dispersion for \mathcal{H}_C or \mathcal{H}_1 —not shown—has stronger finite-size effects than the NF, as it involves two interacting QPs instead of one.) Comparing different panels, it is remarkable that as soon as the two minima of the NF dispersion actually form (which for \mathcal{H}_1 seems to require $\delta V_1 \gtrsim 0.02e^2/\lambda$), they remain located at virtually unchanged wave vectors, while the bandwidth depends significantly on the particular model (e.g., on δV_1). The NF₁ lies slightly below the Fermi surface of CFs, i.e., $k_0 \simeq k_F$ (for a half-filled LL of spinless fermions, $k_F = 1/\lambda$). This confirms the expectation of Bogoliubov theory, that in a weakly paired phase of CFs, and for weak coupling, the minimum is close to k_F [4]. The presence of the second minimum NF₂ is more surprising. It could arise as a superposition of a NF with additional magnetoroton excitations. However, the combined energy for a NF and magnetoroton is found to be larger than NF₂. Tentatively, this feature could be related to a bound state of these objects. In any case, we conclude that NF₂ cannot decay into a NF and a magnetoroton, so it is a genuine feature of the NF dispersion describing a long-lived excitation, and can be tested in experiment.

Direct observation of the NF requires a probe changing the fermion number of the second LL. One such probe is photoluminescence (PL), in which an electron in this LL recombines with a photoexcited valence band hole (the “1,0” or “1,1” PL lines in Ref. [23]). The PL spectrum depends on the nature of the state into which the hole relaxes prior to recombination. Often for fractional quantum Hall systems, this is an “excitonic” state in which the hole binds a charge e to form a neutral exciton moving in the background incompressible liquid [24]. Interestingly, for the $\nu = 5/2$ state there are two distinct excitonic states [25] with different parity of the electron number N_e [26]. Which of these two states has the lower energy is difficult to predict: this amounts to determining the fusion channel of four QEs in the presence of the hole. However, if such fusion can be viewed as pairwise, then the two pairs will fuse either both to 1 or both to ψ , giving an overall even N_e . PL recombination removes one electron, so it leaves a

final state with opposite parity to the initial state. For an odd initial N_e , recombination can occur to the MR ground state, yielding a sharp PL line (symmetrically broadened by disorder). For even initial N_e , recombination leaves an odd final N_e , and so must involve the creation of a NF. Since the excitonic state is typically easily localized by disorder [24], the resulting PL spectrum will probe the density of states of the NF band. The minima in the NF band will appear as two asymmetrically broadened peaks [27]. The observation of this double-peak structure in PL would allow direct measurements of the energies of the minima of the NF band.

The evolution of the dispersion minimum NF_1 as a function of δV_1 , sketched as gray lines in Fig. 2(b), gives some insight into the nature of phase transitions. At small δV_1 , the NF gap Δ_{NF} remains nonzero, so we expect the pairing nature of the phase to survive up to the transition into a charge-density wave [19]. The collective NF mode flattens at small δV_1 ; however, the spectrum is dominated by finite-size effects below $V_1 \approx 0.02$. At large δV_1 (approaching interactions resembling the lowest LL) a smooth decay of Δ_{NF} indicates a continuous weakening of pairing (see also [9]). The minimum becomes steeper, and the gap collapses near $k \approx k_F = 1/\lambda$, consistent with a crossover into the CF Fermi liquid state.

We now turn to investigate the physics of a neutral fermion in the presence of two QEs or QHs. For a pair of QHs in the MR state (even N), \mathcal{H}_{Pf} features a band of zero-energy states spaced by $\Delta L = 2$ and terminating at $L = N/2$ [28]. We find that when a neutral fermion is added to the system, low-lying states are found at the same angular momenta [see Figs. 1(g)–1(i)]. Empirically, we find that quasielectron states behave similarly. First, we identify a band of low-lying states for \mathcal{H}_{Pf} at even N , again with $\Delta L = 2$ but terminating at $L = N/2 - 2$. For 2QE + NF configurations, a band with the same angular momenta is obtained by removing one particle and two flux from the system, so the lowest energy NF states can be thought of as holelike. Figures 1(a)–1(c) show example 2QE + NF spectra.

In the presence of QPs the low-lying excitations are not as well separated in the spectrum as for the NF alone, and more significant finite-size effects are expected. Thus, we proceed carefully in analyzing the energetics. The presence of the NF can affect the angular momentum, relative positions, and shape of the QPs, changing their interaction energy in an unknown way. Since we are unable to subtract these effects systematically, instead we average the energy over all states in the low-energy band associated with the two QP (+ NF) excitations. Thus, we evaluate the (properly normalized) average $\langle E_\alpha \rangle \sim \sum_L (2L + 1) E_\alpha(L)$. The energy of each eigenstate $E_\alpha(L)$ is measured with respect to the ground state energy [22] (where α indicates a set of QP numbers, $\alpha = 2QE, 2QH, 2QE + NF$, and $2QH + NF$). For Coulomb Hamiltonians, we apply standard corrections to the energies $E_\alpha(L)$, including using rescaled magnetic length, applying an electrostatic charging correction of the energies, and correcting for the

Coulomb interactions of QPs δV_{QP} in the excited configurations [29–31].

The values $\langle E_\alpha \rangle$ obtained after applying the charging corrections are shown in Figs. 3(a)–3(c) and reflect the total energy of a system with two QEs or QHs, with or without an additional NF [31]. Importantly, for odd N it contains the energy cost for adding a NF in the presence of a pair of QEs (Δ_{NF}^-) or QHs (Δ_{NF}^+), that includes the interaction of the NF with these QPs. Since we average over all positions of the QPs, in finite-size systems one expects a nonzero splitting Δ_{NF}^\pm between the ψ and 1 channels from configurations with overlapping QPs. Estimates from trial states of QHs at close separation [32] suggest that for Coulomb interactions the splitting is $\approx 0.01 e^2/\lambda \equiv \Delta_{max}$. When averaged over all possible QP positions, the contribution would be significantly smaller, due in part to its oscillatory behavior. For finite systems, we find the (average) splitting of fusion channels, including its finite-size effects, satisfies $\Delta_{NF}^\pm \lesssim \Delta_{max}$. [For \mathcal{H}_{Pf} in Fig. 3(c), the splitting Δ_{NF}^+ vanishes by construction [28].]

If the QPs are non-Abelian Ising anyons, as in the Moore-Read phase, then the energy splitting Δ_{NF}^\pm should scale to zero in the thermodynamic limit. The extrapolations in Figs. 3(a)–3(c) are consistent with the vanishing of both splittings Δ_{NF}^\pm (in each case, the extrapolated value is considerably smaller than its standard deviation). Although we cannot prove that the splittings vanish exactly, we emphasize that their best estimates are at least an order of magnitude smaller than the charge gap Δ_c or Δ_{NF} . It is highly nontrivial to find a near degeneracy on this scale. We have examined the behavior as a function of δV_1 in our model Hamiltonian \mathcal{H}_1 . We find that the splitting remains similarly small over the same range of interactions for which the $L = 0$ ground state has a large gap and a high overlap with the Moore-Read Pfaffian [9,10], i.e.,

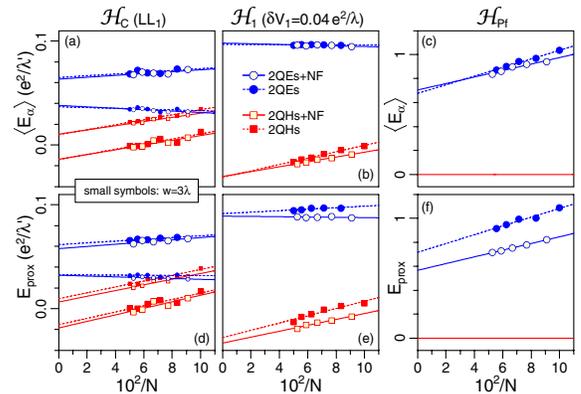


FIG. 3 (color online). Comparison of the total energies of a pair of QEs or QHs, with and without an additional NF (see text for the precise definition), for systems of different size N and for the different Hamiltonians of Fig. 1. Top: energies $\langle E_\alpha \rangle$ averaged over all 2QE/2QH (+ NF) states; bottom: energies E_{prox} taken for the smallest average QE-QE or QH-QH distance. While $\langle E_\alpha \rangle$ extrapolate to similar values with or without NF, the differences are significant for E_{prox} .

extending upwards in δV_1 from about \mathcal{H}_C towards the point of collapse of the NF gap. We take these results as evidence that, over this range of δV_1 , the QEs and the QHs in these realistic systems have non-Abelian exchange statistics of the form of the Moore-Read phase.

To investigate which fusion channel is preferred at short distance, we have estimated the splitting Δ_{NF}^{\pm} for QPs at near-coincident points. In Figs. 3(d)–3(f), we show $E_{\text{prox}} = E_{\alpha}(L_{\text{max}})$ using the low-lying 2QP/2QP + NF states with the largest angular momentum (and closest separation of the QPs) [33]. In this case the odd-even splitting opens for each of the considered Hamiltonians. The splitting also remains in the thermodynamic limit, revealing a slightly negative Δ_{NF}^{\pm} , signaling a preference for the ψ channel [e.g., $\Delta_{\text{NF}}^+ = -0.0053(41)e^2/\lambda$ and $\Delta_{\text{NF}}^- = -0.0023(23)e^2/\lambda$ for \mathcal{H}_1 , and $\Delta_{\text{NF}}^+ = -0.15(3)$ for \mathcal{H}_{PF}]. Previously, the splitting had been known only for QHs, and was based on variational wave functions [32]. Here, we report the splittings for both QE and QH based on exact calculations in finite systems.

In conclusion, we have analyzed the neutral fermion excitations of the $\nu = 5/2$ state for microscopic two-body Hamiltonians. We showed that these exhibit similar properties to those of the Pfaffian model, for which the Moore-Read phase is the exact ground state. The neutral fermion is gapped in the ground state, but is gapless in the presence of a pair of quasiparticles. These results provide important evidence that the $\nu = 5/2$ state has the properties of the weakly paired Moore-Read phase. Our studies also elucidate additional physical properties of this phase: we predict that characteristic features of the NF dispersion can appear in photoluminescence experiments; and we show evidence for the nature of the fusion of two QEs or QHs.

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