UNIVERSITY OF CAMBRIDGE

Pseudizing the Hamiltonian

Gareth Conduit
Pascal Bugnion
Richard Needs
Pablo López Ríos
Jonathan Lloyd-Williams

TCM Group, Department of Physics

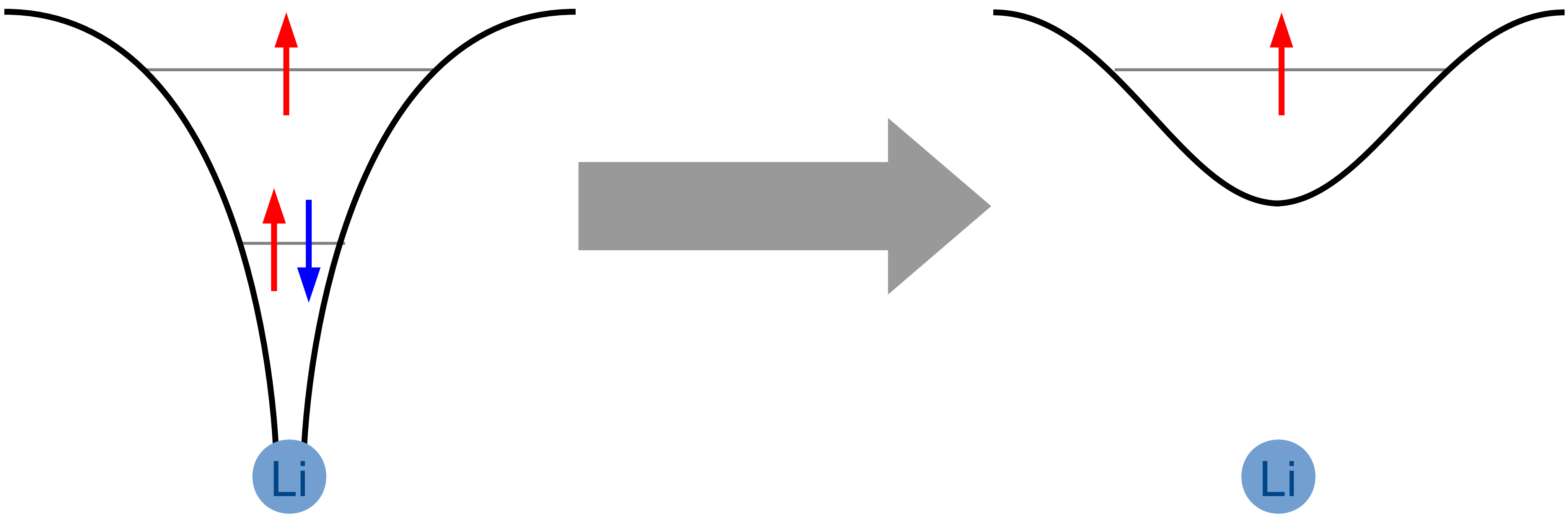
Pseudopotentials

$$H = KE + V_{\text{e-i}} + V_{\text{e-e}}$$

$$E = \frac{\int \bar{\psi} H \psi \, \mathrm{d}\boldsymbol{r}}{\int \bar{\psi} \psi \, \mathrm{d}\boldsymbol{r}}$$

Pseudopotentials

$$H = KE + V_{e\text{-}i} + V_{e\text{-}e}$$

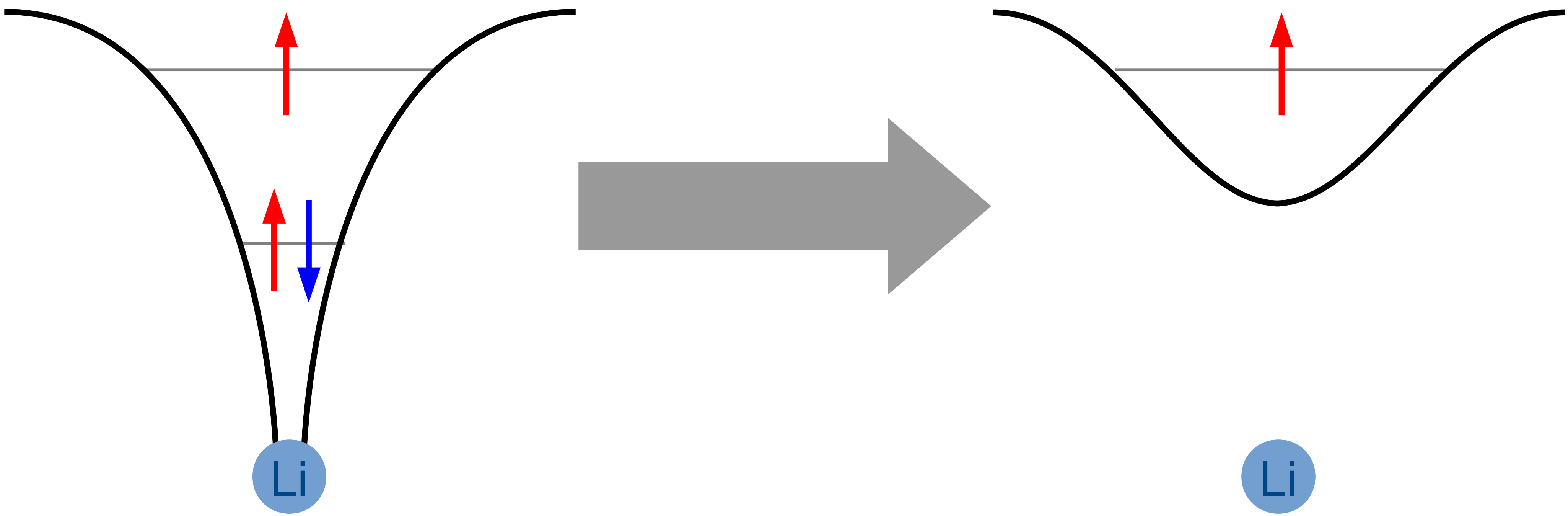


Pseudopotentials

$$H = KE + V_{e\text{-}i} + V_{e\text{-}e}$$

Smooth background

Fewer electrons



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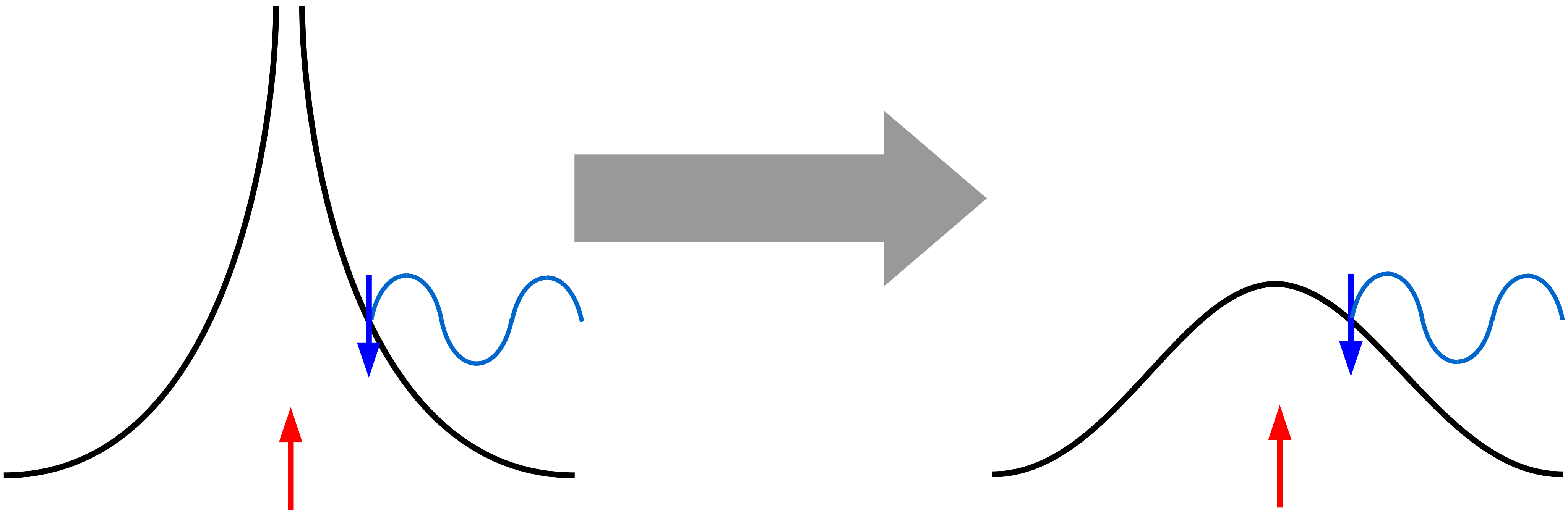
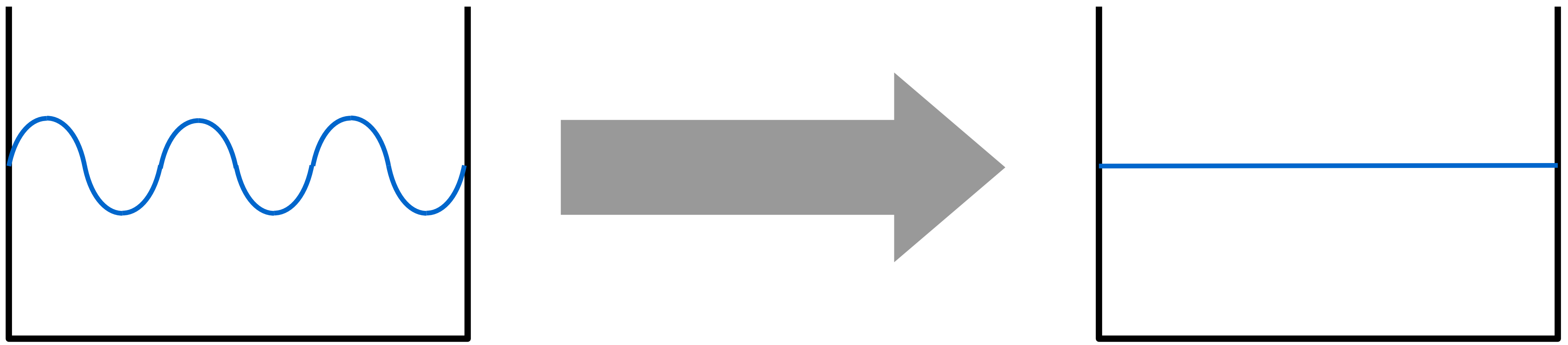


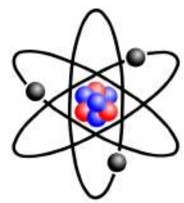
Pseudopotentials

$$H = KE + V_{e\text{-}i} + V_{e\text{-}e}$$

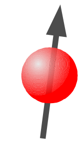
Smooth integrand

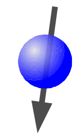


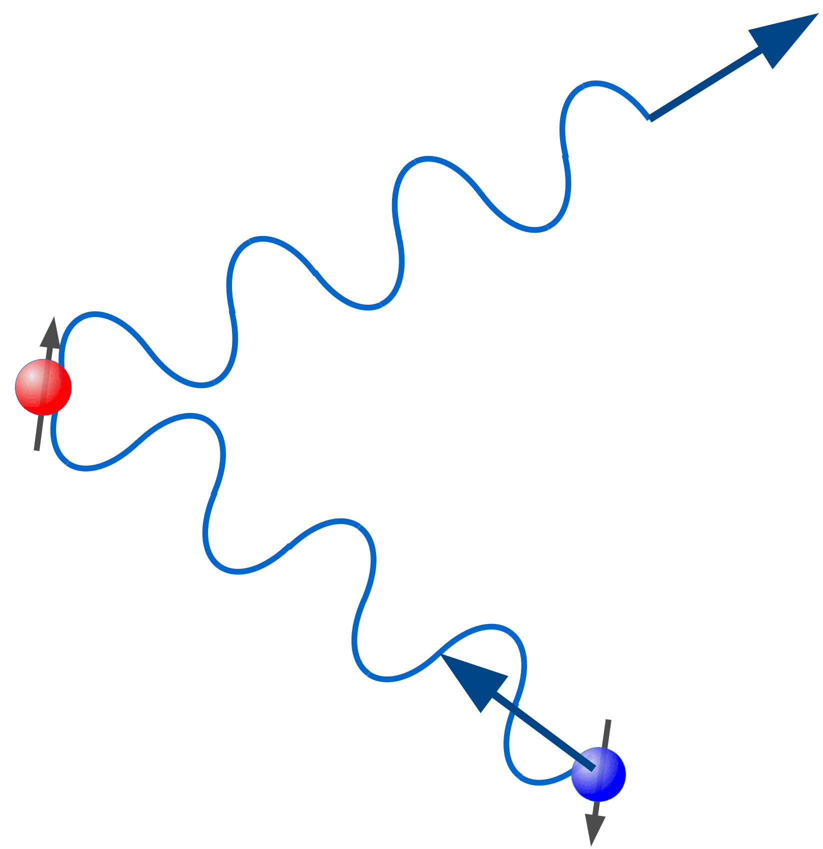
Scattering in ultracold atom gases



^{6}Li atom

$|F = 1/2, m_F = 1/2\rangle$ →  Up spin electron

$|F = 1/2, m_F = -1/2\rangle$ →  Down spin electron



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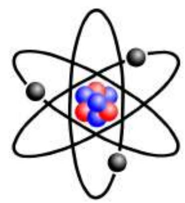
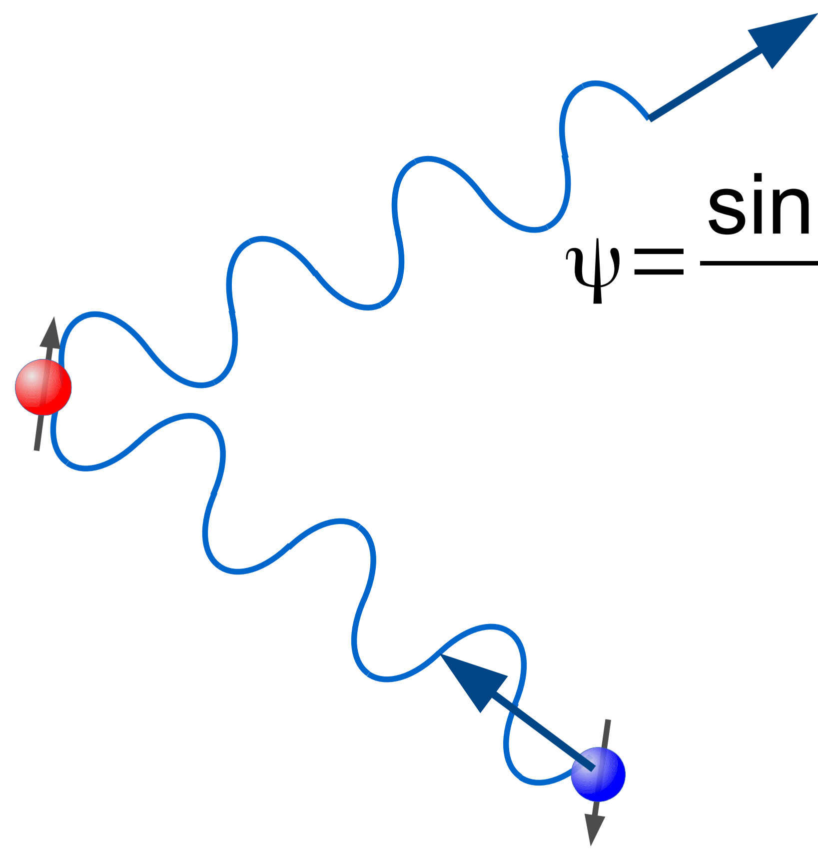


^{6}Li atom

$|F = 1/2, m_F = 1/2\rangle$ → Up spin electron

$|F = 1/2, m_F = -1/2\rangle$ → Down spin electron

$$V(r) = \frac{a}{2r^2}\delta(r)$$

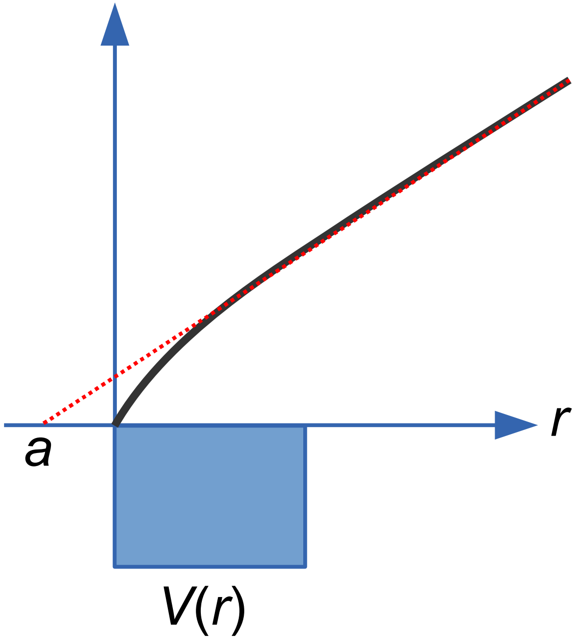


$$\psi = \frac{\sin(kr + \delta)}{kr}$$

Scattering potentials

Underlying attractive

Effective attractive



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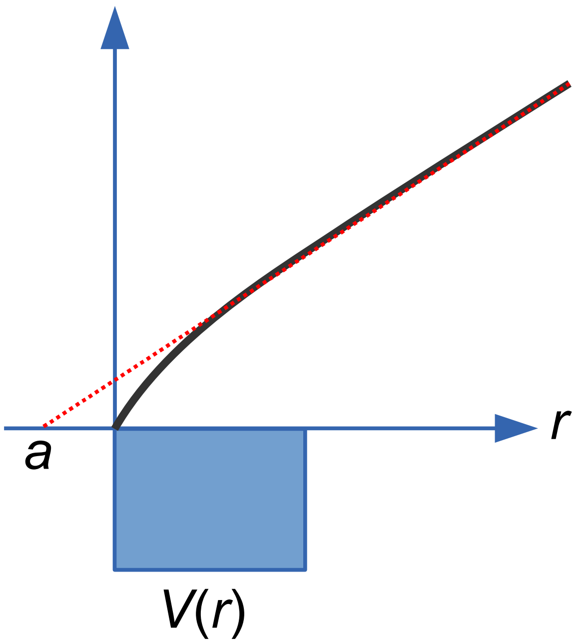
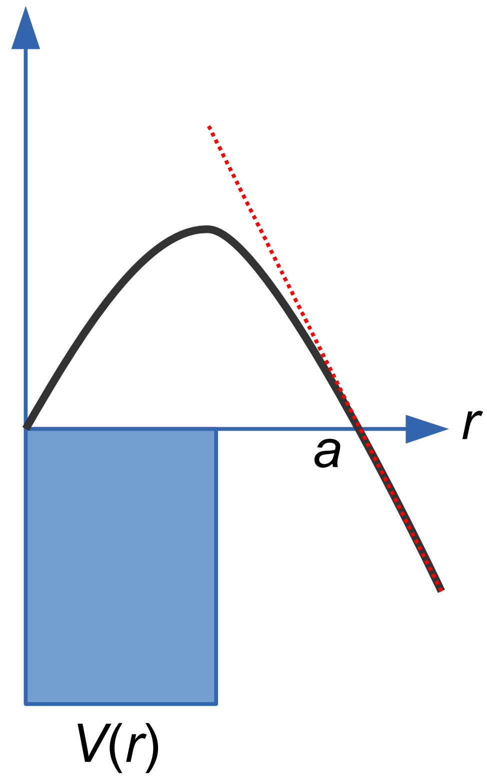


Underlying attractive

Effective repulsive



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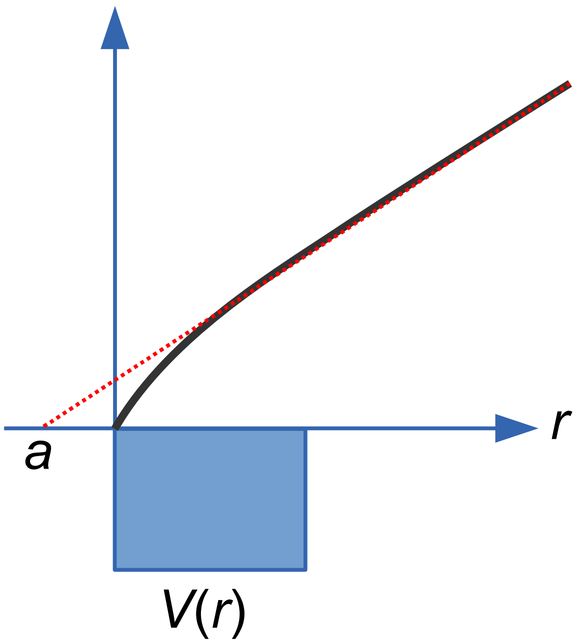

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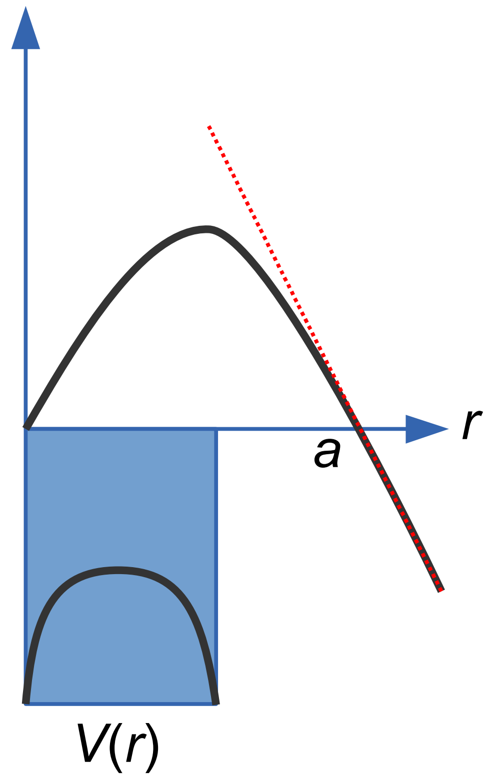

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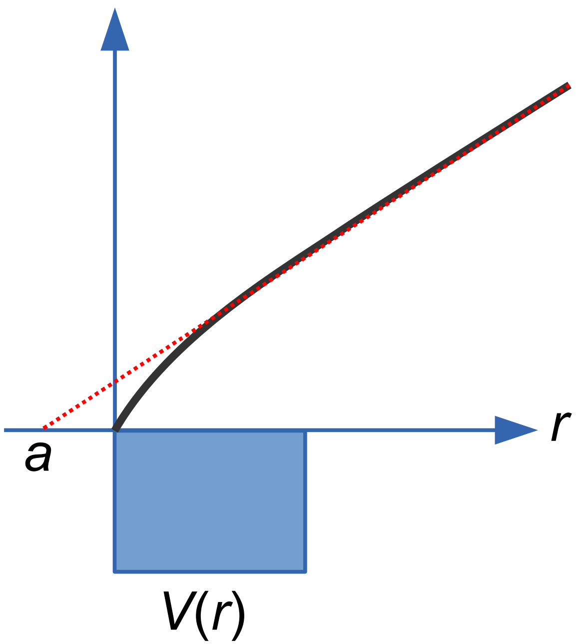

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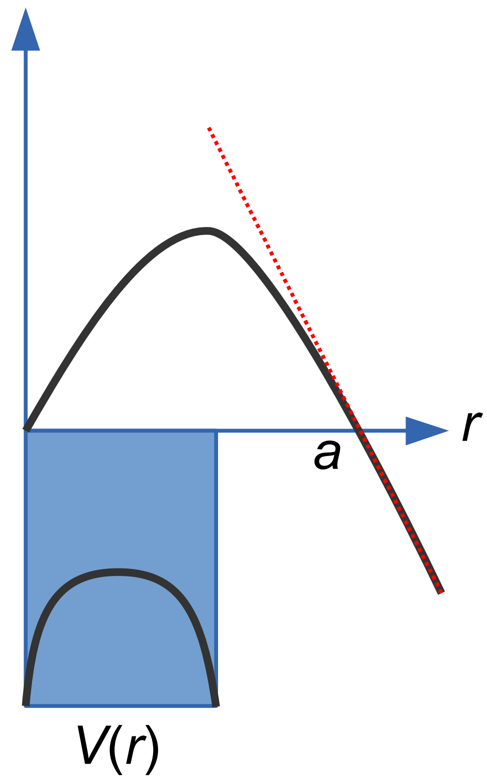
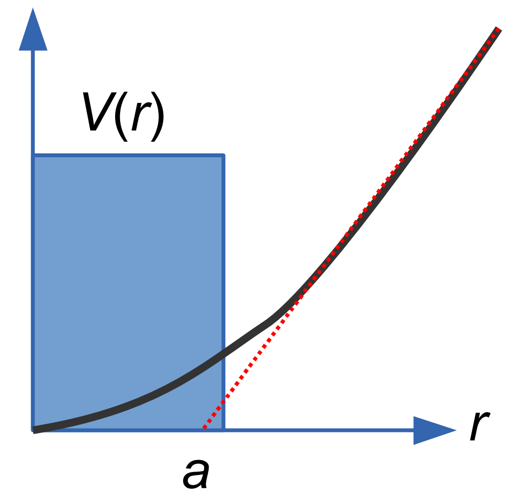
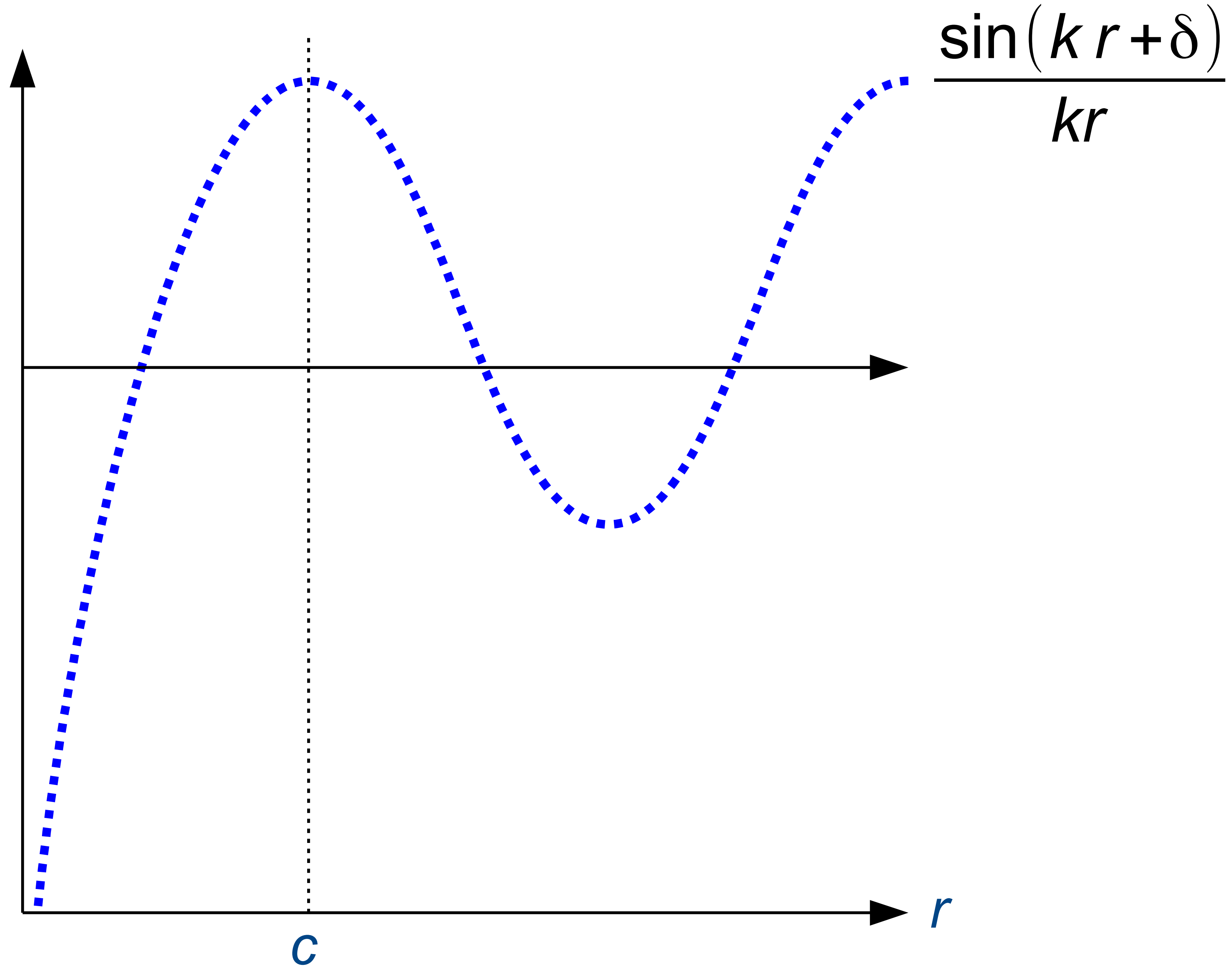


Underlying repulsive

Effective repulsive



Construction of a pseudopotential



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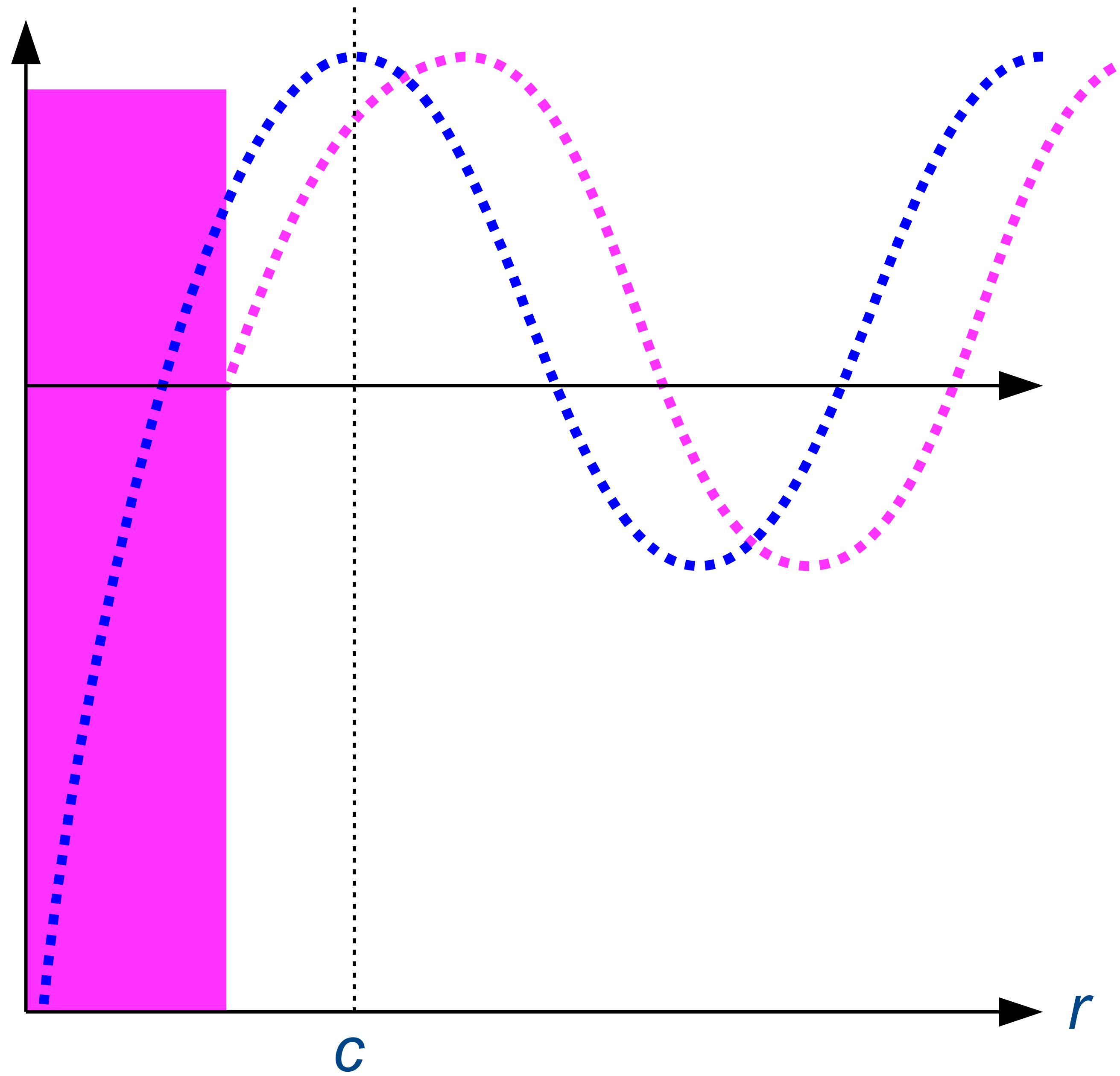

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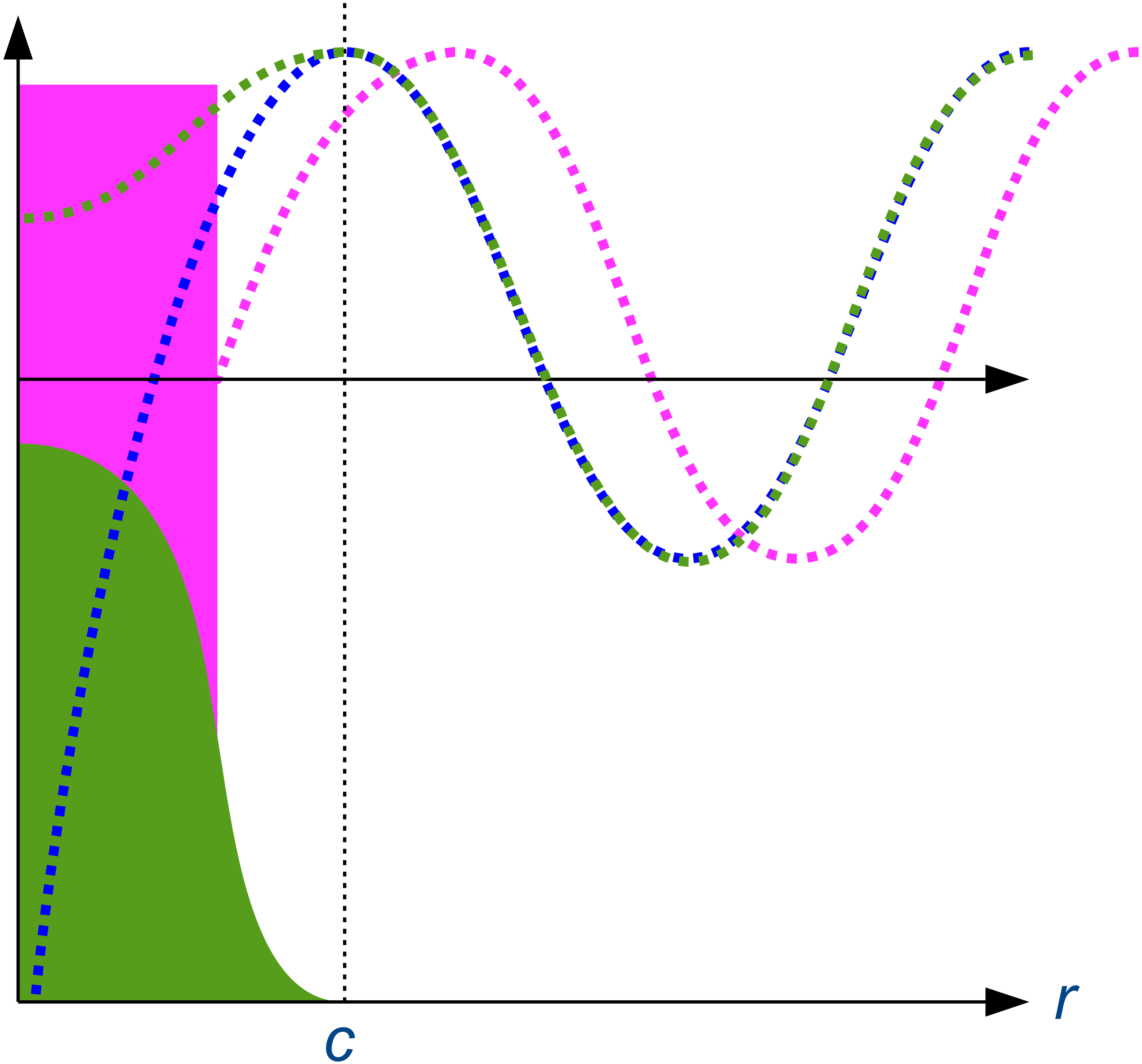
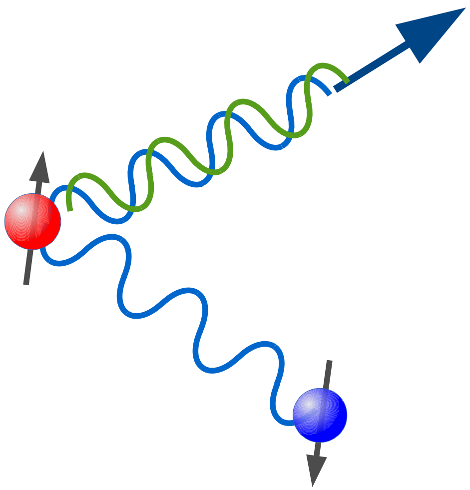


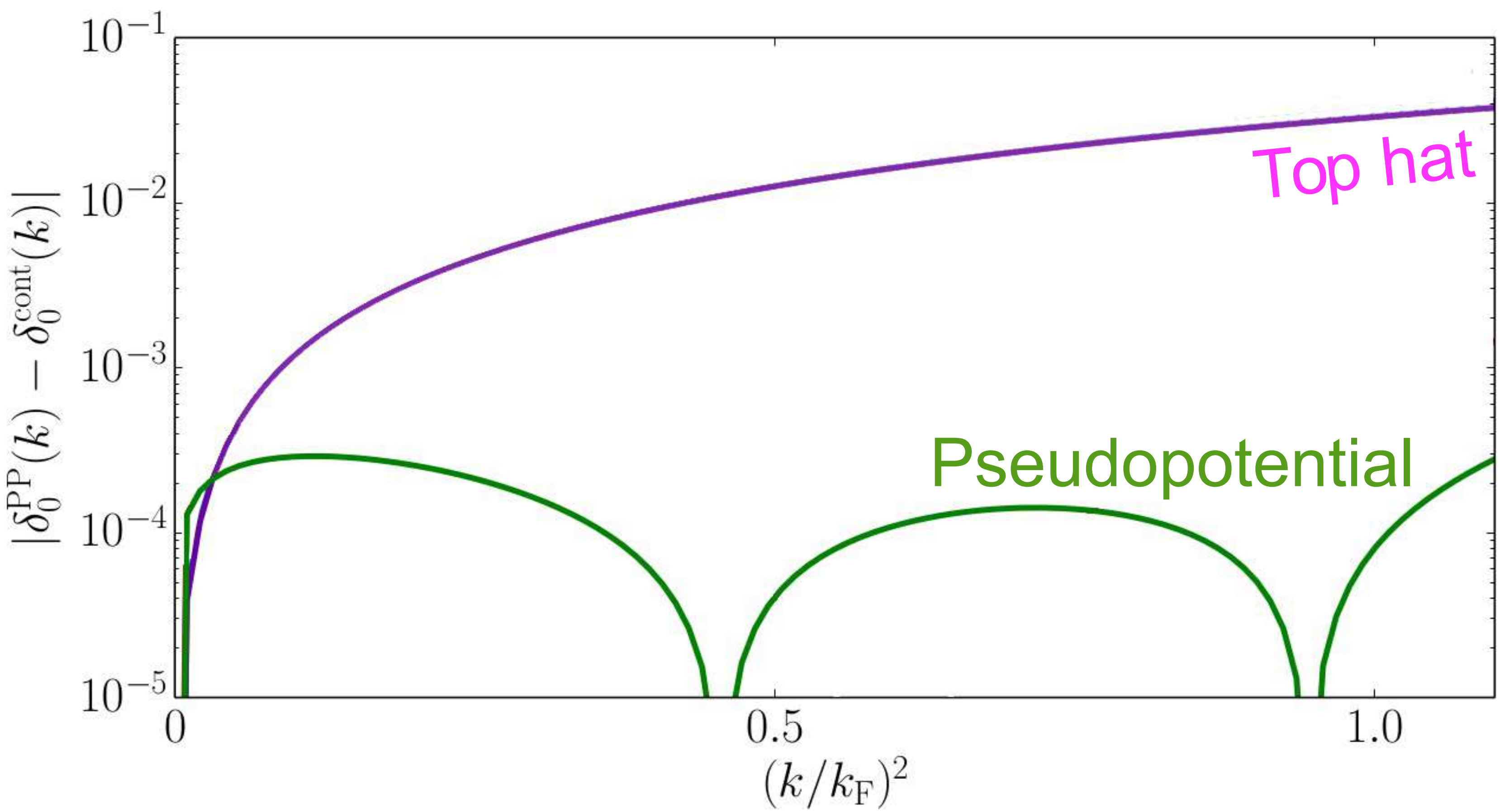
Construction of a pseudopotential

$$V_{\mathrm{PP}}(r)=\begin{cases}\dfrac{1}{c}+\left(1-\dfrac{r}{c}\right)^2\left[v_1\left(\dfrac{1}{2}+\dfrac{r}{c}\right)+\sum_{i=2}^{N_{\mathrm{v}}} v_i\left(\dfrac{r}{c}\right)^i\right] & r<c\\ \dfrac{1}{r} & r>c\end{cases}$$

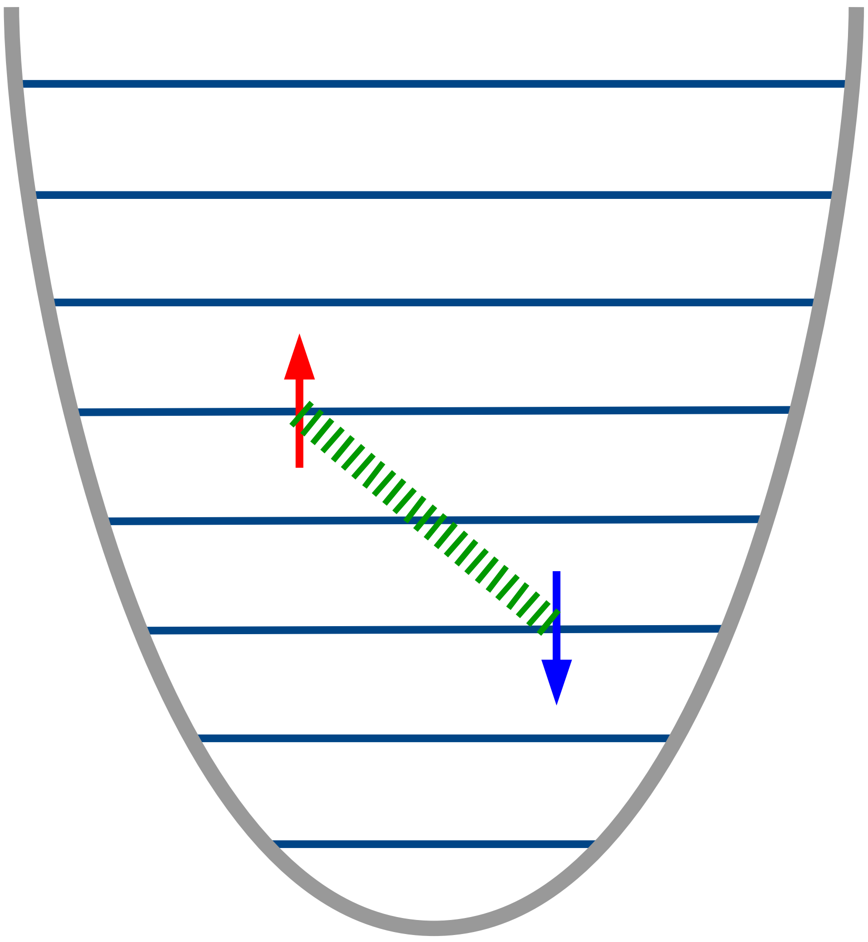
$$\sum_{l=0}^{l_{\max}}\int_0^{k_{\mathrm{F}}}\left[\left.\frac{\mathrm{d}\ln\psi_{\mathrm{PP}}(k,l)}{\mathrm{d}r}\right|_c-\left.\frac{\mathrm{d}\ln\psi_{\mathrm{cont}}(k,l)}{\mathrm{d}r}\right|_c\right]^2\mathrm{d}k$$

Pseudopotential: scattering properties





Pseudopotential: two atoms in a trap



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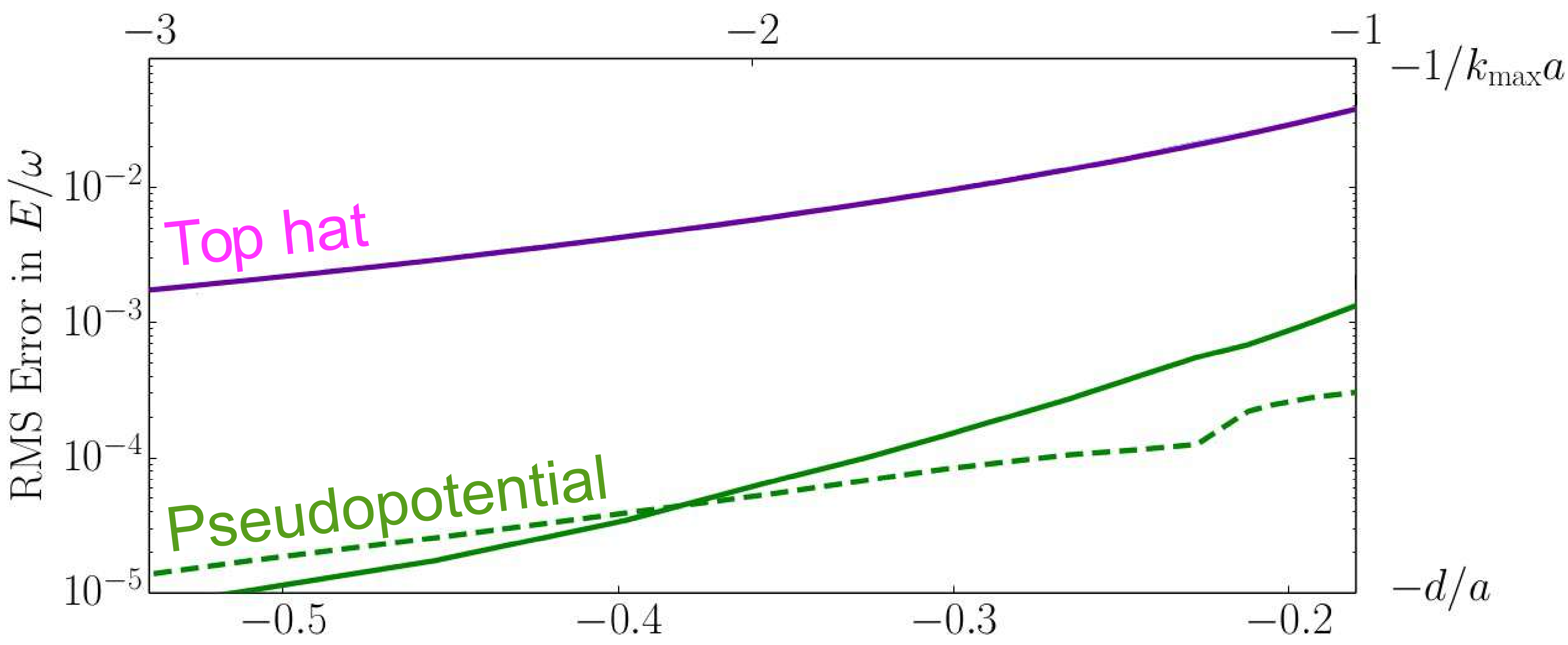


Pseudopotentials summary

Repulsive & attractive state: 100 times more accurate, 1000 times faster

Bound state: 1000 times more accurate, 1000 times faster

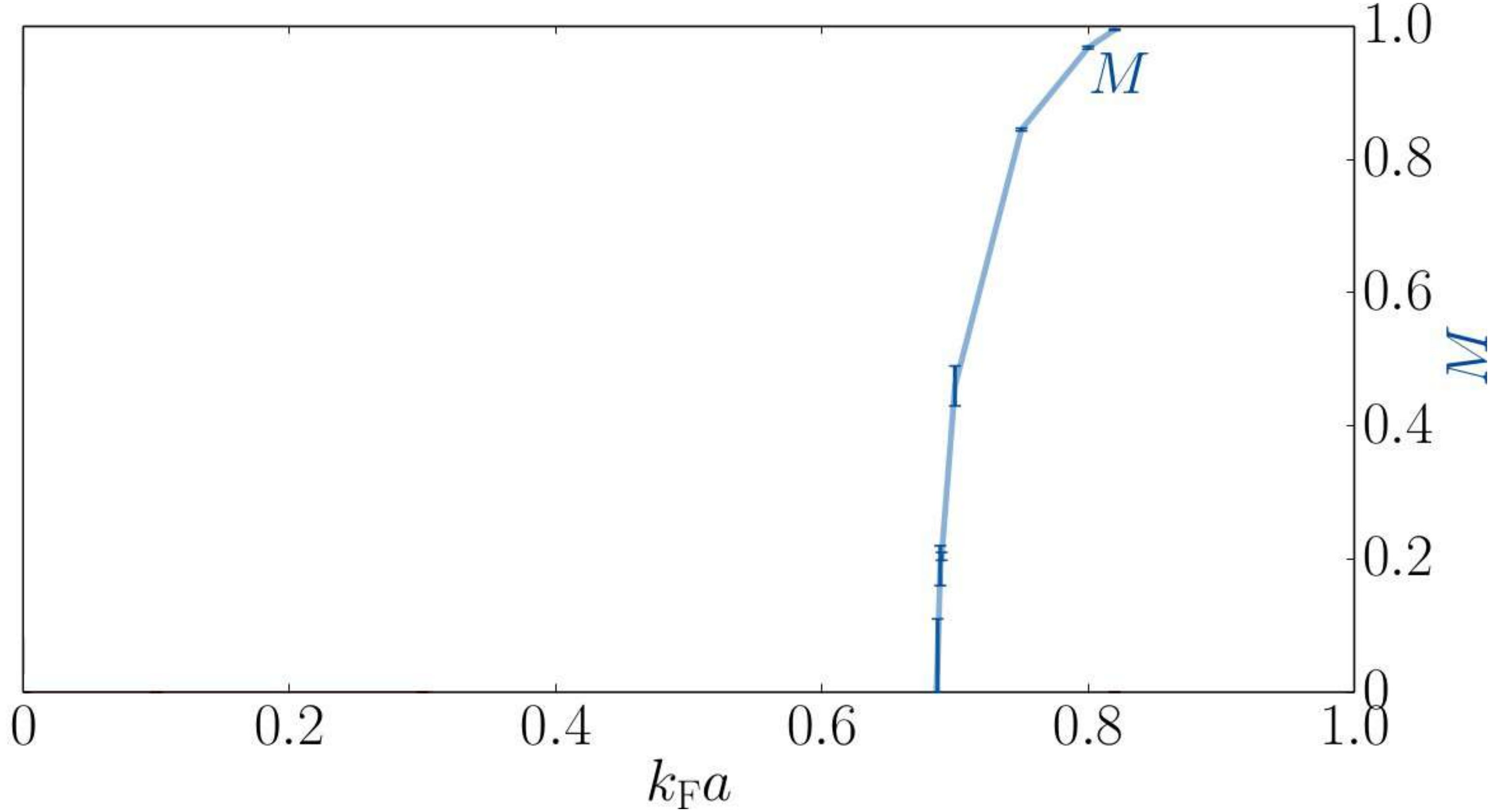
Stoner Hamiltonian

$$H = \frac{-\nabla^2}{2} + 4\pi a \delta(\boldsymbol{r}_{\uparrow} - \boldsymbol{r}_{\downarrow})$$

Theories of ferromagnetism

Stoner mean-field theory	Second order	$k_F a$=1.57
Fluctuations beyond Hertz-Millis	First order	-
Polaron theory	First order	-
Field theory	First order	$k_F a$=1.054
Tan relations	No magnetism	-
DMC hard sphere	First order	$k_F a$=0.81(2)
Hartree Fock MC	First / second order	$k_F a$=0.83(2)

Stoner Hamiltonian



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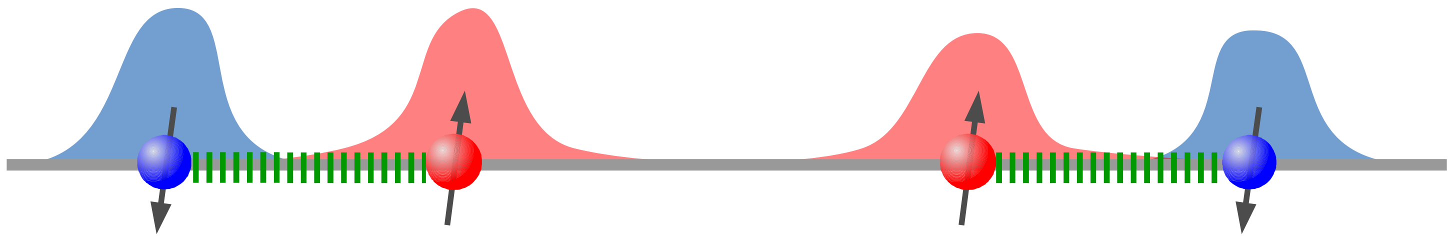
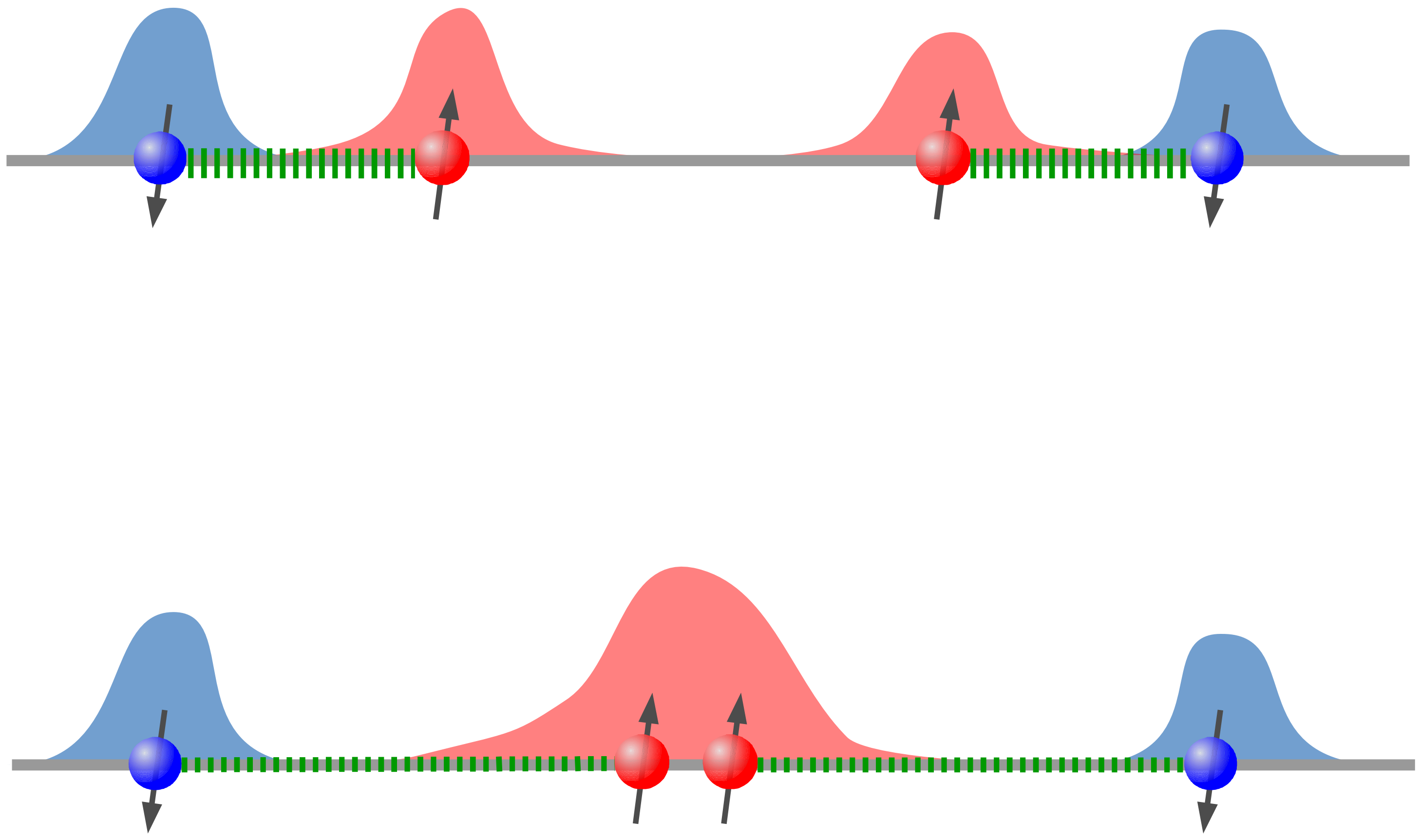
DMC hard sphere	First order	$k_F a$=0.81(2)
Hartree Fock MC	First / second order	$k_F a$=0.83(2)
DMC pseudopotential	Second order	$k_F a$=0.683(1)

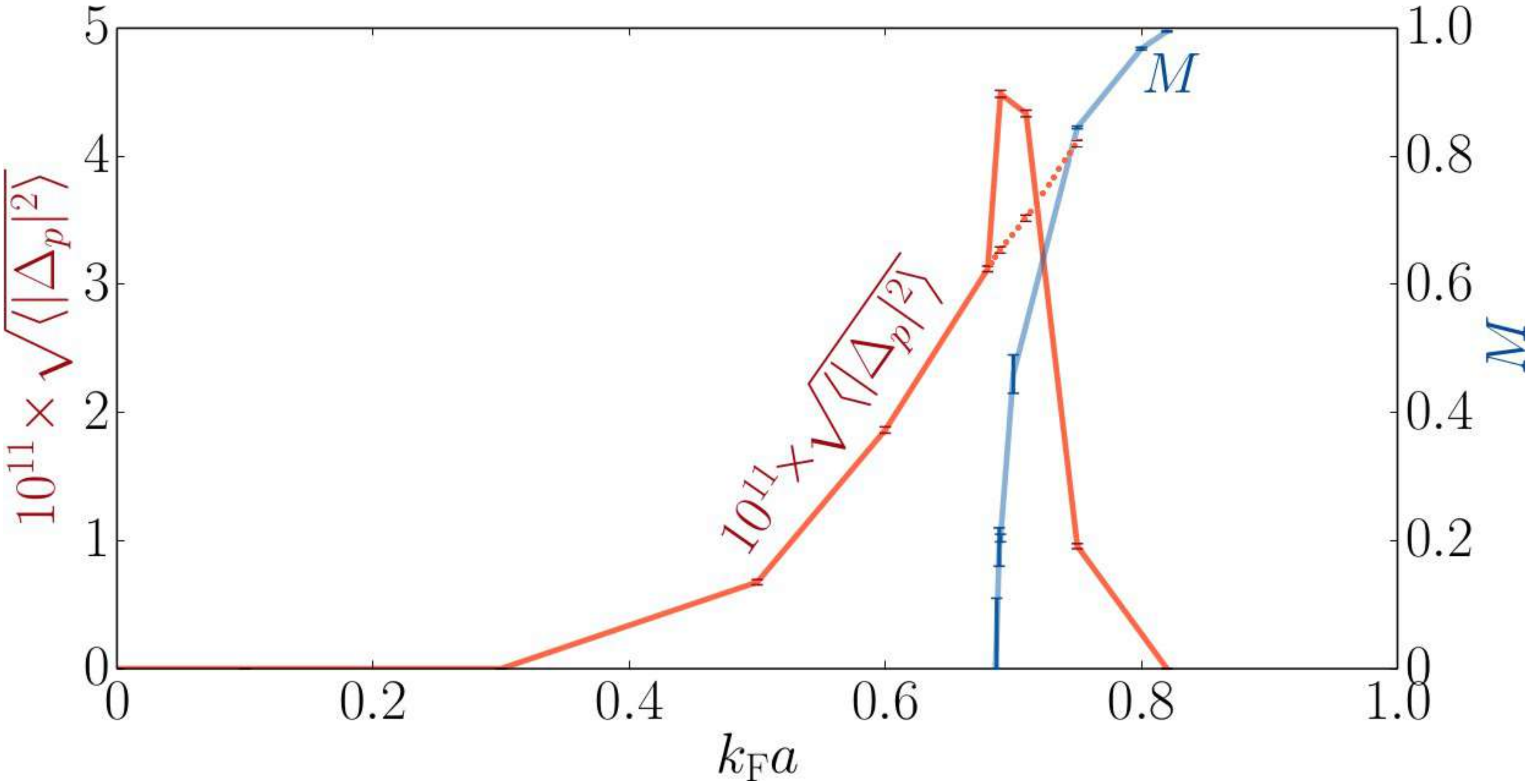
Fluctuation contributions



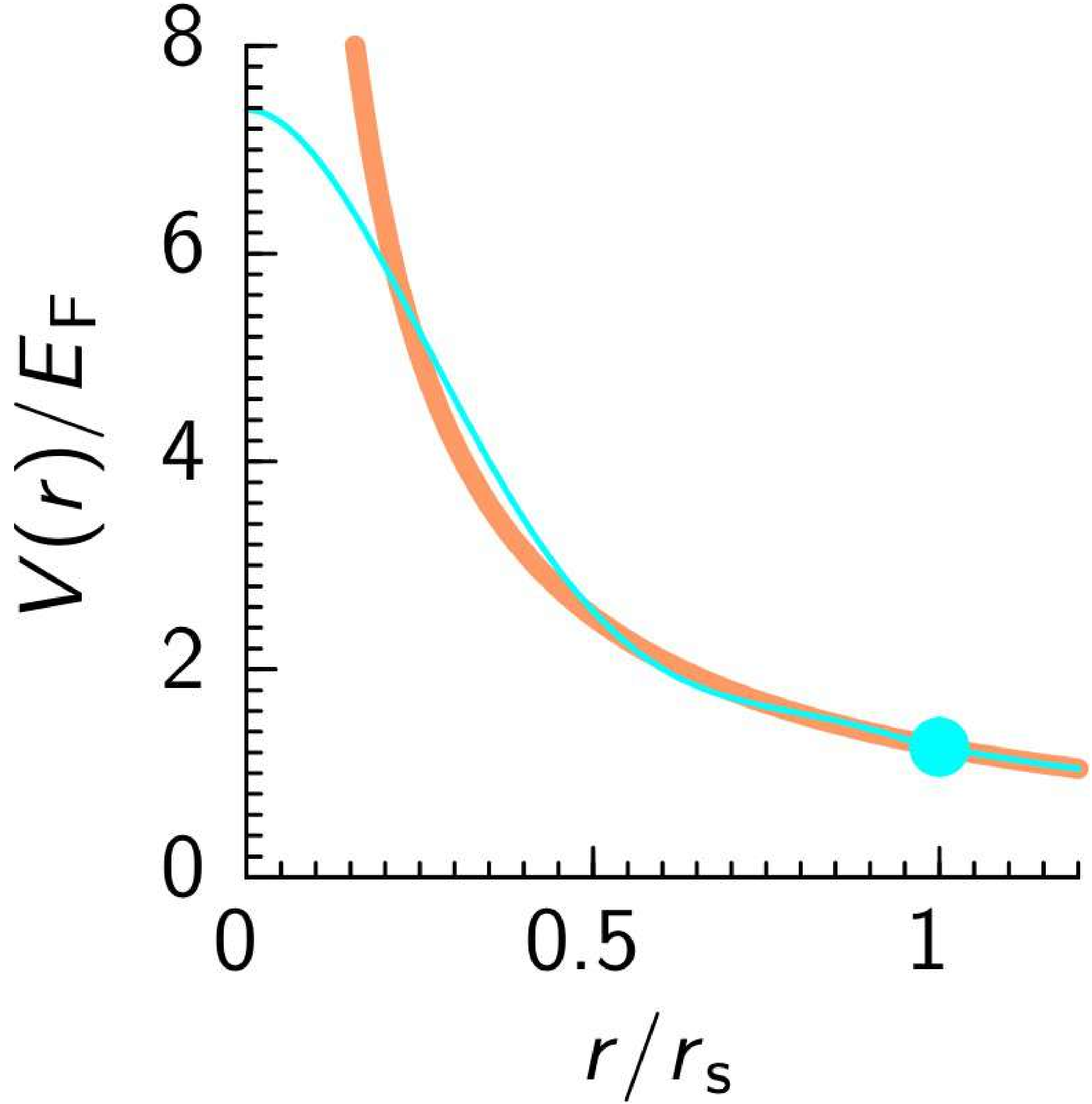
Fluctuation contributions



Stoner Hamiltonian



Coulomb pseudopotential



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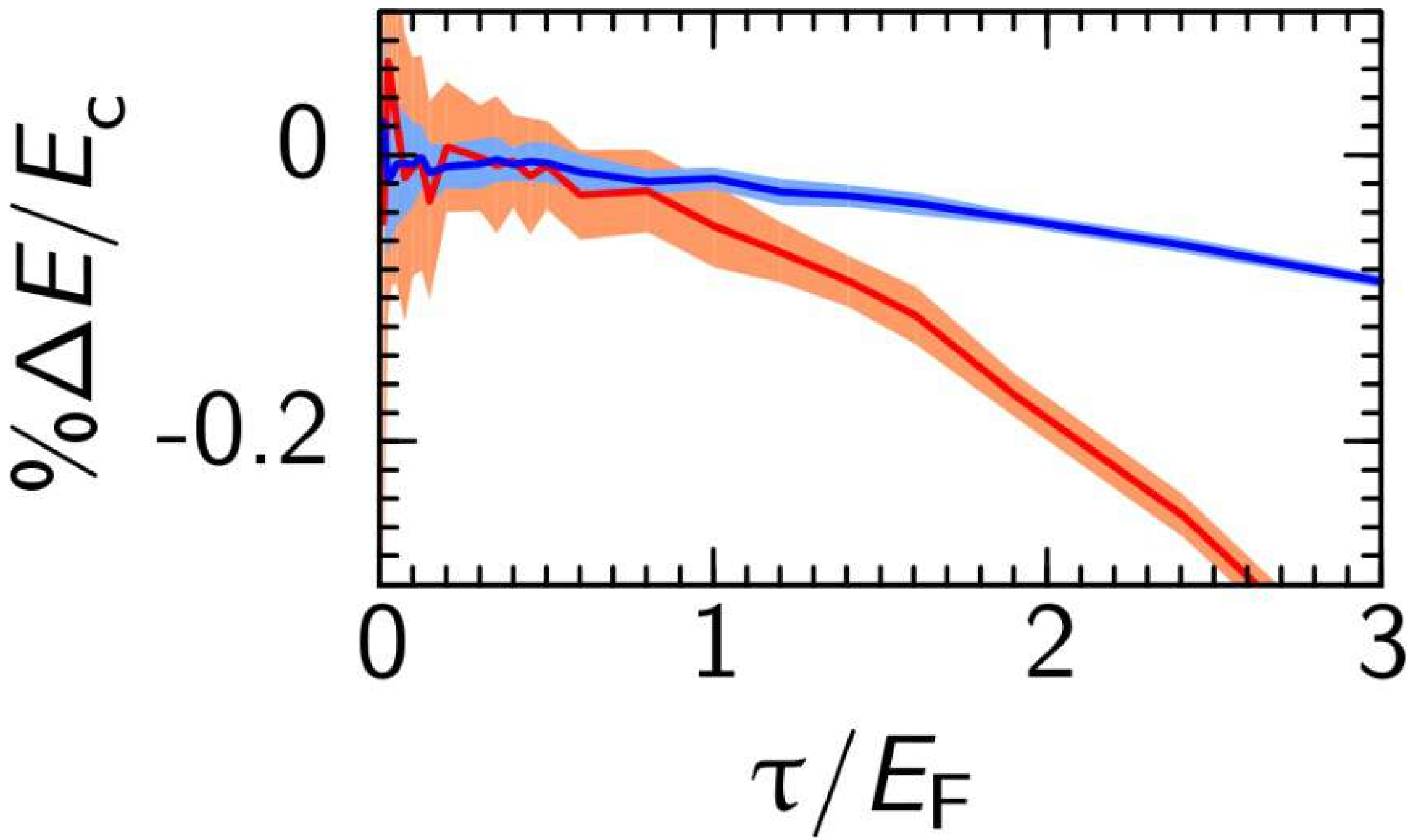
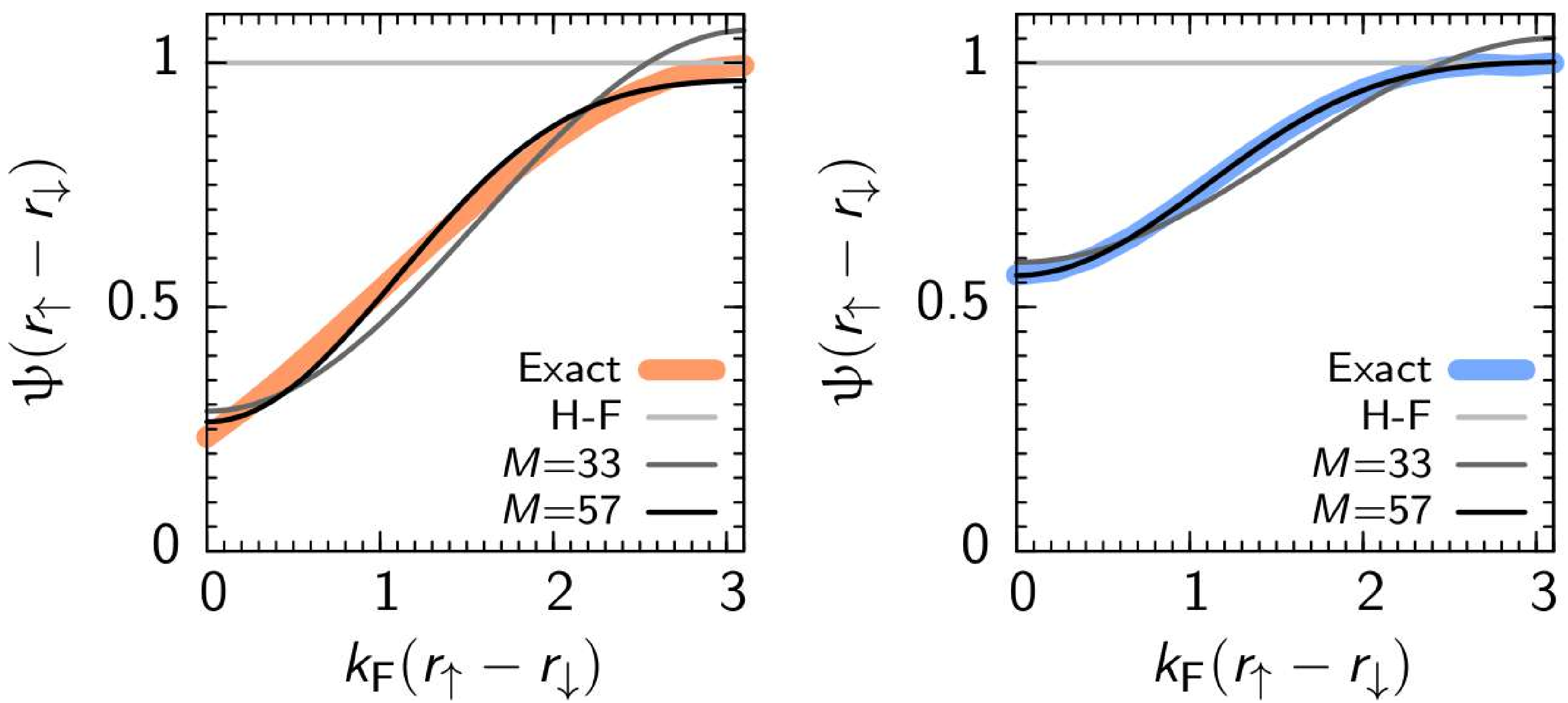


Coulomb pseudopotential



Pseudopotentials summary

Created a pseudopotential for the contact interaction that is 100 times more accurate, 1000 times faster

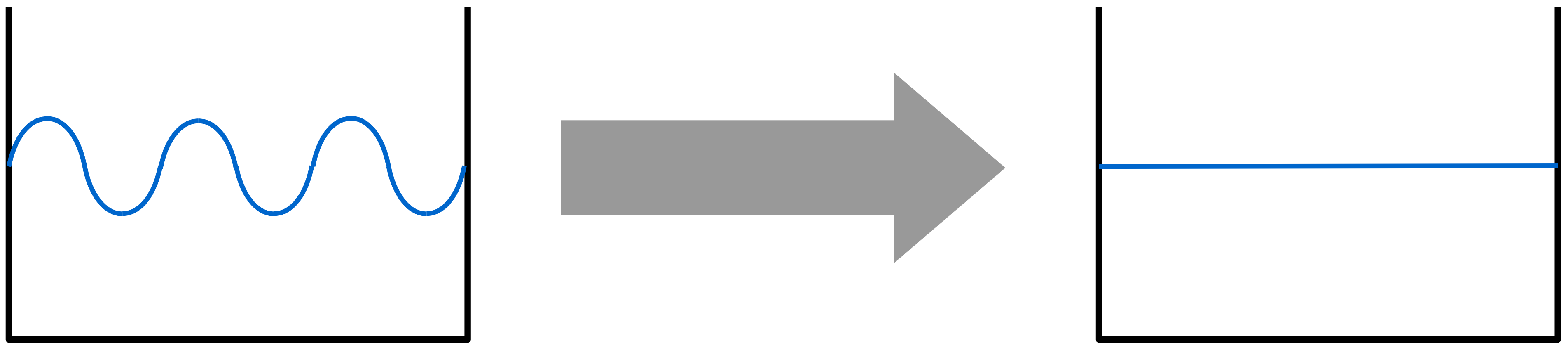
Stoner Hamiltonian displays second order ferromagnetic phase transition and p-wave ordering

Created a pseudopotential for the Coulomb interaction that is 30 times faster

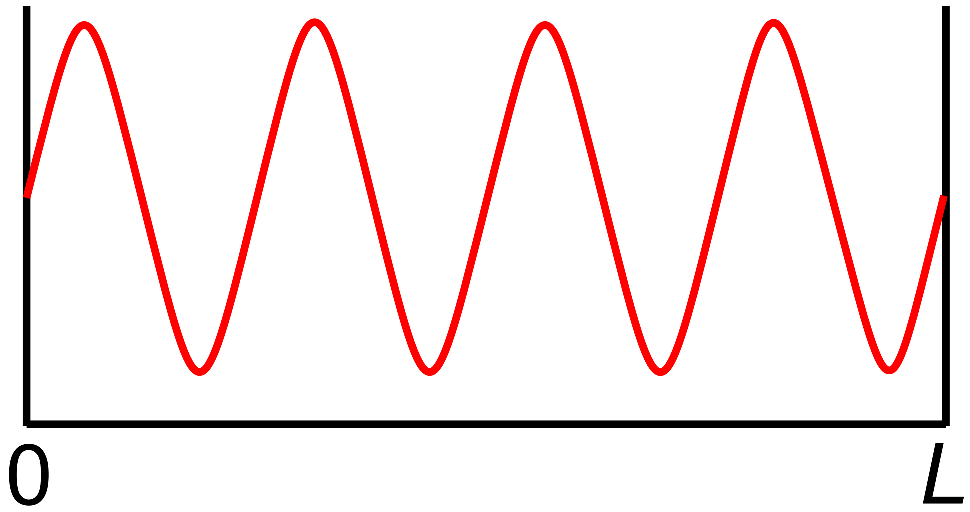
Pseudopotentials

$$H = KE + V_{e\text{-}i} + V_{e\text{-}e}$$

Smooth integrand



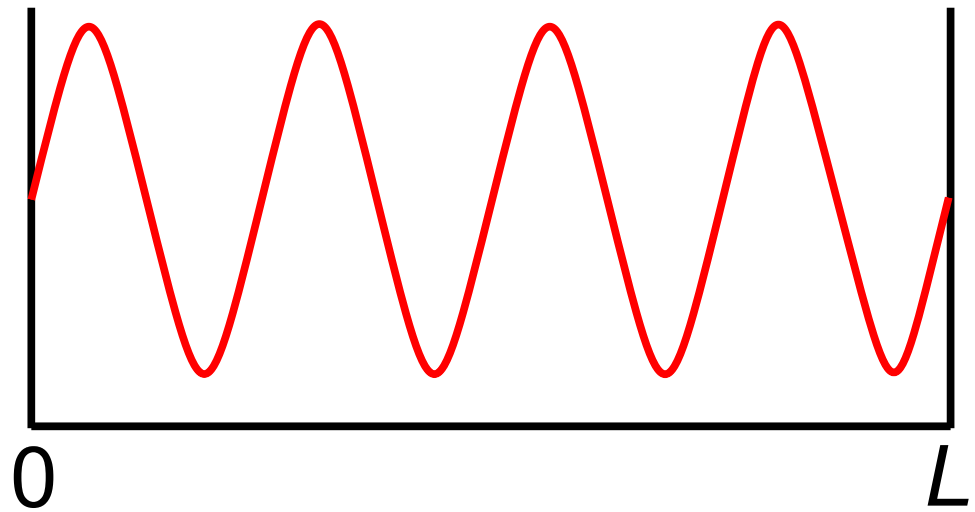
Particle in a box



$$\psi = A\sin\left(8\pi\frac{x}{L}\right)$$

$$E = -\int \bar{\psi}\frac{\nabla^2}{2}\psi\,\mathrm{d}x$$

Particle in a box

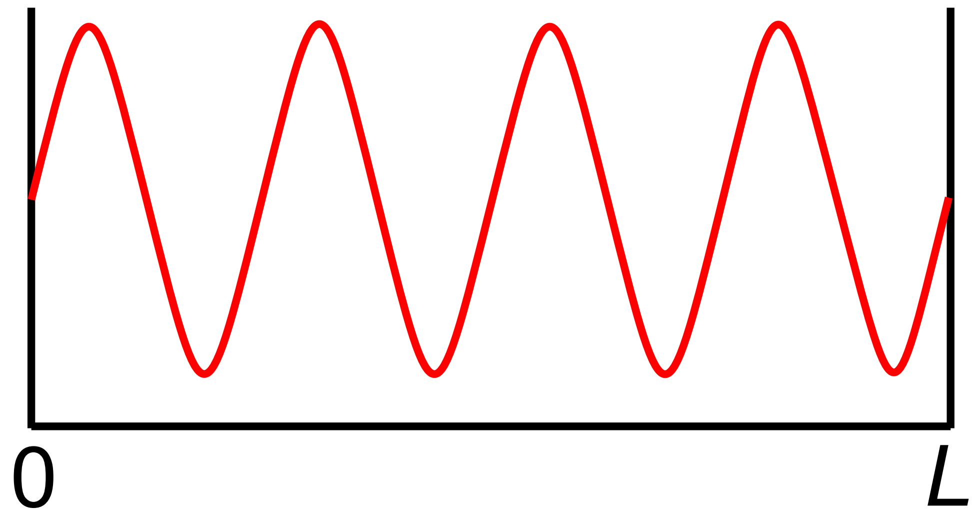


$$\psi = A \sin\left(8\pi \frac{x}{L}\right)$$

$$E = -\int \bar{\psi} \frac{\nabla^2}{2} \psi \, dx$$

$$= \frac{32\pi^2 A^2}{L^2} \int \sin^2\left(8\pi \frac{x}{L}\right) dx$$

Particle in a box



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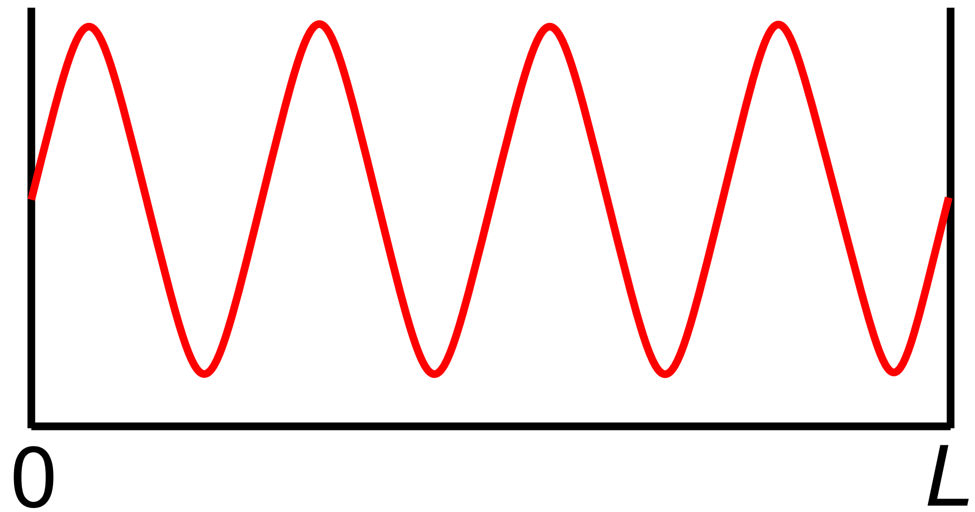
$$\langle \sin^2 \rangle = \frac{1}{2}$$

Particle in a box



$$\psi = A \sin\left(8\pi \frac{x}{L}\right)$$

$$E = -\int \bar{\psi} \frac{\nabla^2}{2} \psi \, \mathrm{d}x$$

$$= \frac{32\pi^2 A^2}{L^2} \int \sin^2\left(8\pi \frac{x}{L}\right) \mathrm{d}x$$

$$= \frac{16\pi^2 A^2}{L}$$

Particle in a box

$$E = \langle \sin^2 \rangle$$

$$\langle \sin^2 \rangle = \frac{1}{2}$$

Particle in a box

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$$\cos^2 + \sin^2 = 1$$

Particle in a box

$$E = \langle \sin^2 \rangle$$

$$\langle \sin^2 \rangle = \frac{1}{2}$$

$$\cos^2 + \sin^2 = 1$$

$$\frac{d}{dx} \sin x = \cos x$$

Particle in a box

$$-\bar{\psi}\frac{\nabla^2}{2}\psi \rightarrow \frac{k^2}{2}A^2\sin^2(kx)$$

$$\frac{(\nabla\bar{\psi})(\nabla\psi)}{2} \rightarrow \frac{k^2}{2}A^2\cos^2(kx)$$

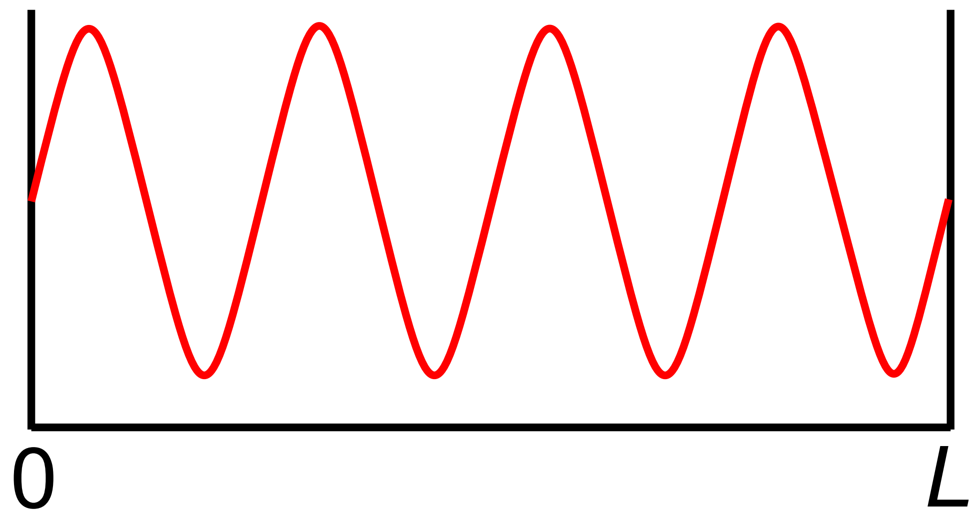
Particle in a box

$$-\bar{\psi}\frac{\nabla^2}{2}\psi \rightarrow \frac{k^2}{2}A^2\sin^2(kx)$$

$$\frac{(\nabla\bar{\psi})(\nabla\psi)}{2} \rightarrow \frac{k^2}{2}A^2\cos^2(kx)$$

$$E = \frac{1}{2}\int -\bar{\psi}\nabla^2\psi \,\mathrm{d}x \rightarrow \frac{1}{4}\int (\nabla\bar{\psi})(\nabla\psi) - \bar{\psi}\nabla^2\psi \,\mathrm{d}x$$

Particle in a box



$$\psi = A \sin\left(8\pi \frac{x}{L}\right)$$

$$E = \frac{1}{4}\int (\nabla \bar{\psi})(\nabla \psi) - \bar{\psi}\nabla^2 \psi \, dx$$

$$= \frac{16\pi^2 A^2}{L^2}\int \cos^2\left(8\pi \frac{x}{L}\right) + \sin^2\left(8\pi \frac{x}{L}\right) dx$$

$$= \frac{16\pi^2 A^2}{L}$$

Wave function normalization

$$N = \int \bar{\psi}\psi \, \mathrm{d}x \rightarrow \frac{1}{2}\int \bar{\psi}\psi + \frac{(\nabla\bar{\psi})(\nabla\psi)}{k^2} \, \mathrm{d}x$$

$$k^2 = \frac{-\nabla^2\psi}{\psi}$$

Pseudizing the Hamiltonian

$$E = \int \hat{KE}\, \bar{\psi}_{\vec{r}}\, \psi_{\vec{r}}$$
$$+ V_{\text{e-i}}(\vec{r} - \vec{r}\,')\, n_{\vec{r}}\, \bar{\psi}_{\vec{r}}\, \psi_{\vec{r}}$$
$$+ V_{\text{e-e}}(\vec{r} - \vec{r}\,')\, \bar{\psi}_{\vec{r}}\, \bar{\psi}_{\vec{r}\,'}\, \psi_{\vec{r}\,'}\, \psi_{\vec{r}}$$

$$N = \int \bar{\psi}_{\vec{r}}\, \psi_{\vec{r}}$$

Pseudizing the Hamiltonian

$$E=\frac{\int \frac{-\bar{\psi}\nabla^2\psi+(\nabla\bar{\psi})(\nabla\psi)}{2}+V\left(\bar{\psi}\psi+\frac{(\nabla\bar{\psi})(\nabla\psi)}{-\nabla^2\psi/\psi}\right)\mathrm{d}r}{\int \bar{\psi}\psi+\frac{(\nabla\bar{\psi})(\nabla\psi)}{-\nabla^2\psi/\psi}\mathrm{d}r}$$

$$E=\frac{\int \left(\frac{-\bar{\psi}\nabla^2\psi}{2}+V\bar{\psi}\psi\right)\left(1+\frac{(\nabla\bar{\psi})(\nabla\psi)}{\bar{\psi}\nabla^2\psi}\right)\mathrm{d}r}{\int \bar{\psi}\psi\left(1+\frac{(\nabla\bar{\psi})(\nabla\psi)}{\bar{\psi}\nabla^2\psi}\right)\mathrm{d}r}$$

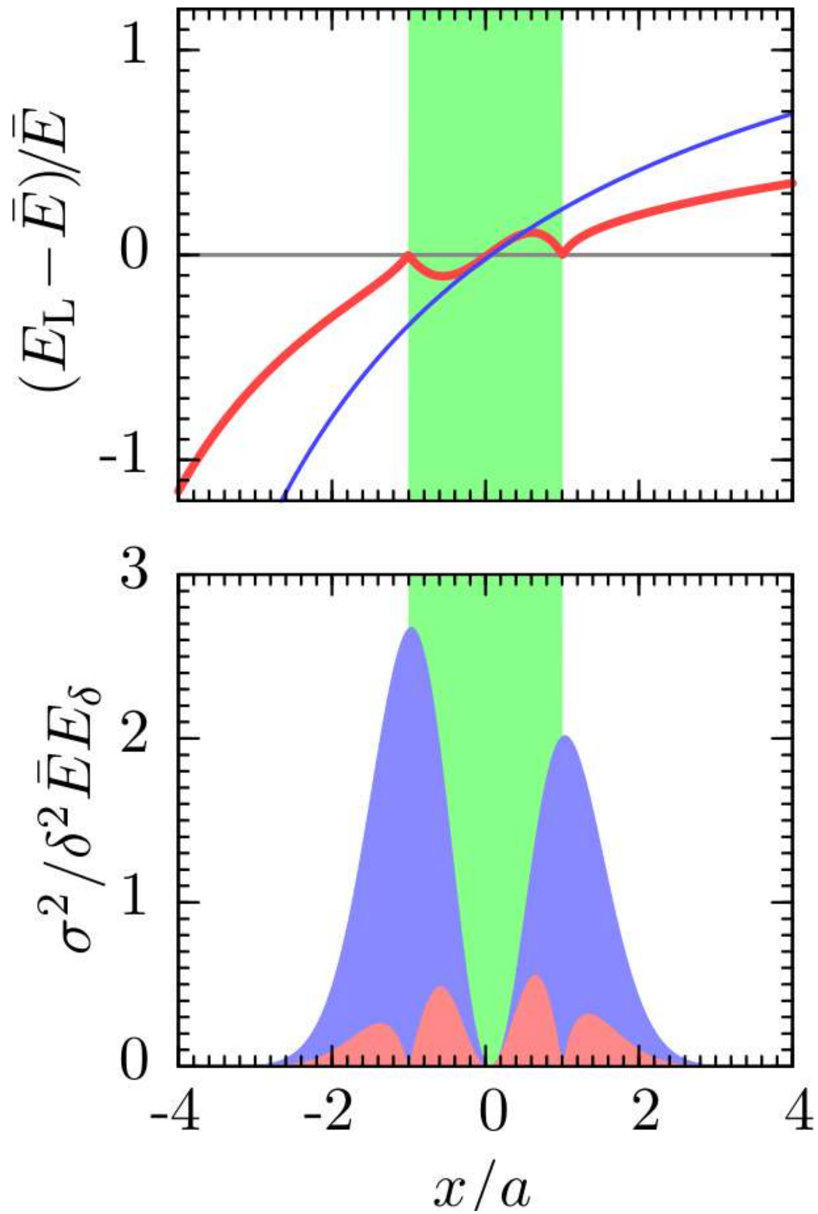
$$E=\frac{\int E\bar{\psi}\psi\left(1+\frac{(\nabla\bar{\psi})(\nabla\psi)}{\bar{\psi}\nabla^2\psi}\right)\mathrm{d}r}{\int \bar{\psi}\psi\left(1+\frac{(\nabla\bar{\psi})(\nabla\psi)}{\bar{\psi}\nabla^2\psi}\right)\mathrm{d}r}=E$$

Pseudizing the Hamiltonian

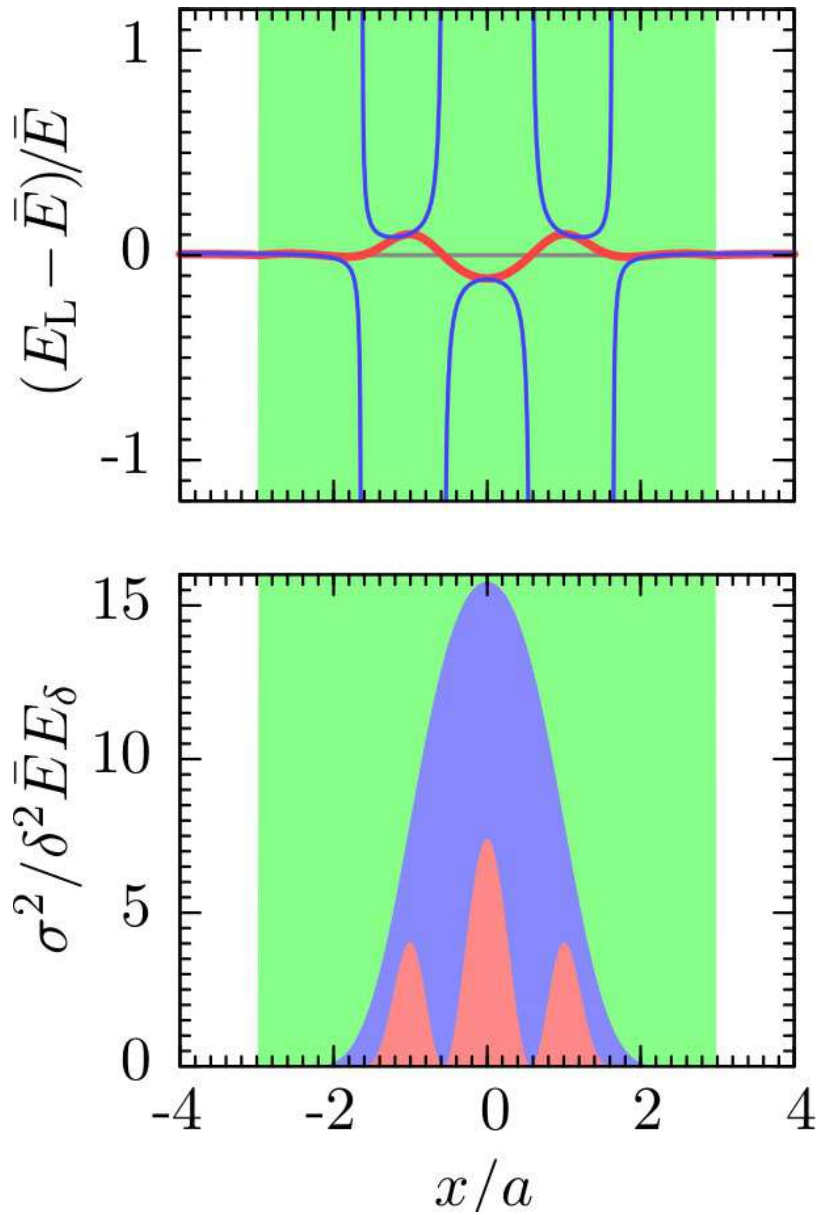
$$E = \frac{\int \bar{\psi} \hat{H} \psi \, \mathrm{d}r}{\int \bar{\psi} \psi \, \mathrm{d}r} = \int_{\bar{\psi}\psi} \frac{\bar{\psi} \hat{H} \psi}{\bar{\psi} \psi} \, \mathrm{d}r$$

$$\frac{\sigma^2_{\mathrm{PP}}}{\sigma^2_{\mathrm{std}}} = \frac{1}{4}$$

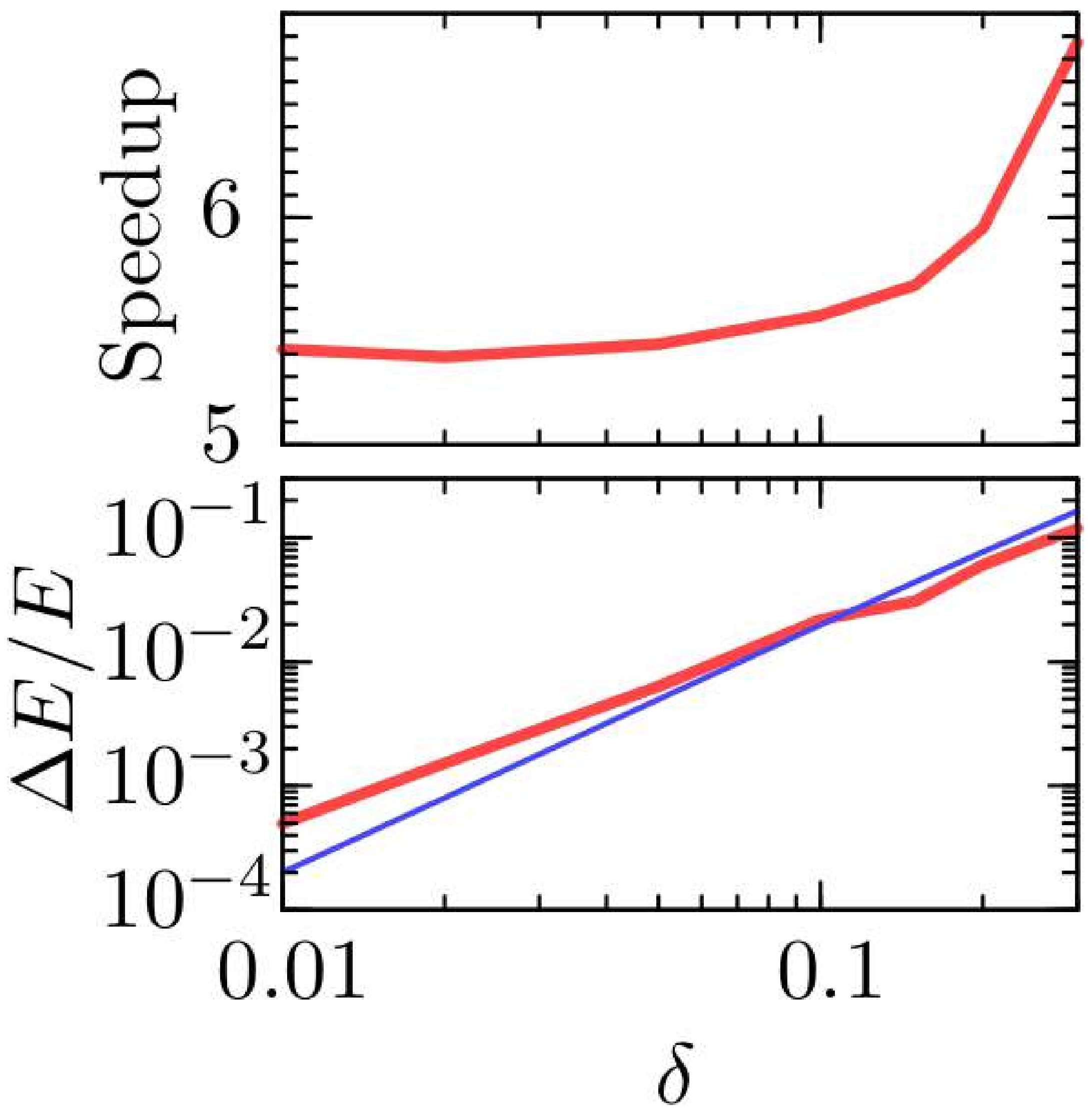
Results: Ψ_0+0.1Ψ_1



Results: $\Psi_4 + 0.1\Psi_0$



Results: $\Psi_0+\delta\ \Psi_1$



Summary

Developed a pseudopotential for the contact and Coulomb interactions

Stoner Hamiltonian displays second order ferromagnetic phase transition, and p-wave ordering

Proposed the formalism to pseudize the kinetic energy and wave function