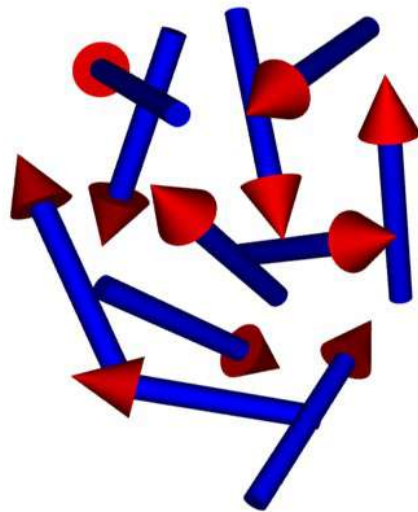
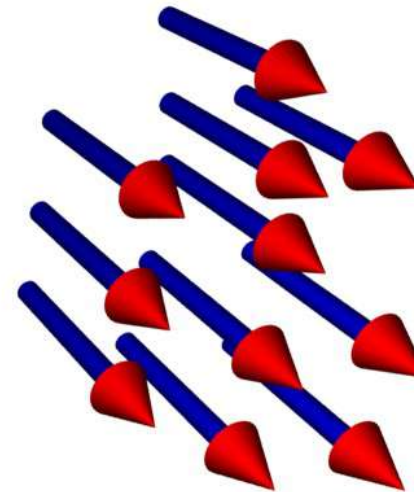


New experimental protocols to probe the Stoner transition in a Fermi gas

Weak interactions



Strong interactions



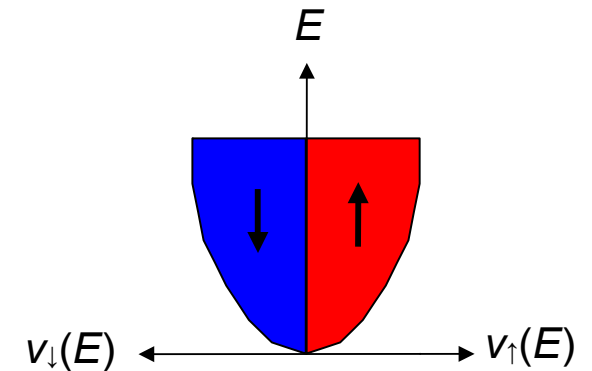
Gareth Conduit¹, **Ben Simons**², **Ehud Altman**¹ & **Curt von Keyserlingk**³

1. Weizmann Institute of Science, 2. University of Cambridge, 3. University of Oxford

Stoner instability with repulsive interactions

$$\hat{H} = \sum_{\mathbf{k}\sigma} \epsilon_{\mathbf{k}} c_{\mathbf{k}\sigma}^\dagger c_{\mathbf{k}\sigma} + g \sum_{\mathbf{k}\mathbf{k}'\mathbf{q}} c_{\mathbf{k}\uparrow}^\dagger c_{\mathbf{k}'+\mathbf{q}\downarrow}^\dagger c_{\mathbf{k}'+\mathbf{q}\downarrow} c_{\mathbf{k}\uparrow}$$

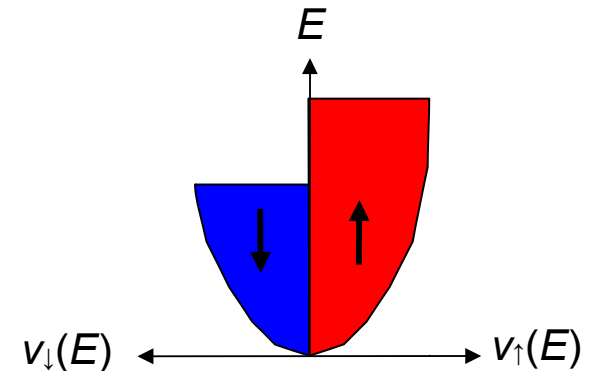
Not magnetised



- Following a mean-field approximation

$$E = \sum_{\mathbf{k},\sigma} \epsilon_{\mathbf{k}} n_{\sigma}(\epsilon_{\mathbf{k}}) + g N_{\uparrow} N_{\downarrow}$$

Partially magnetised

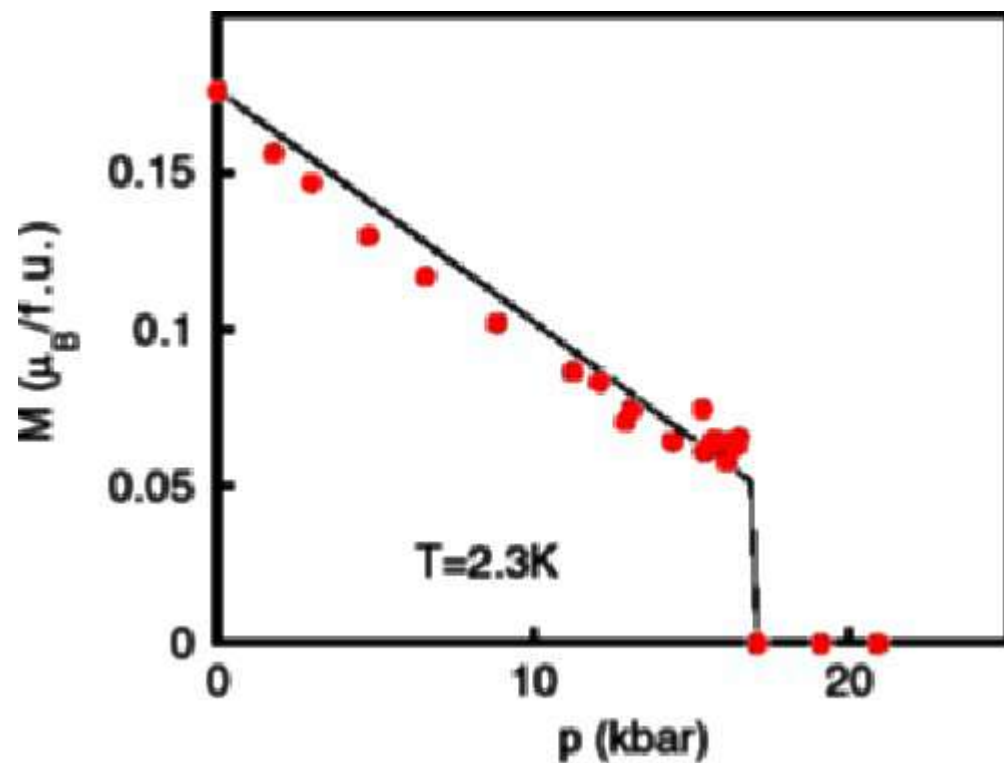
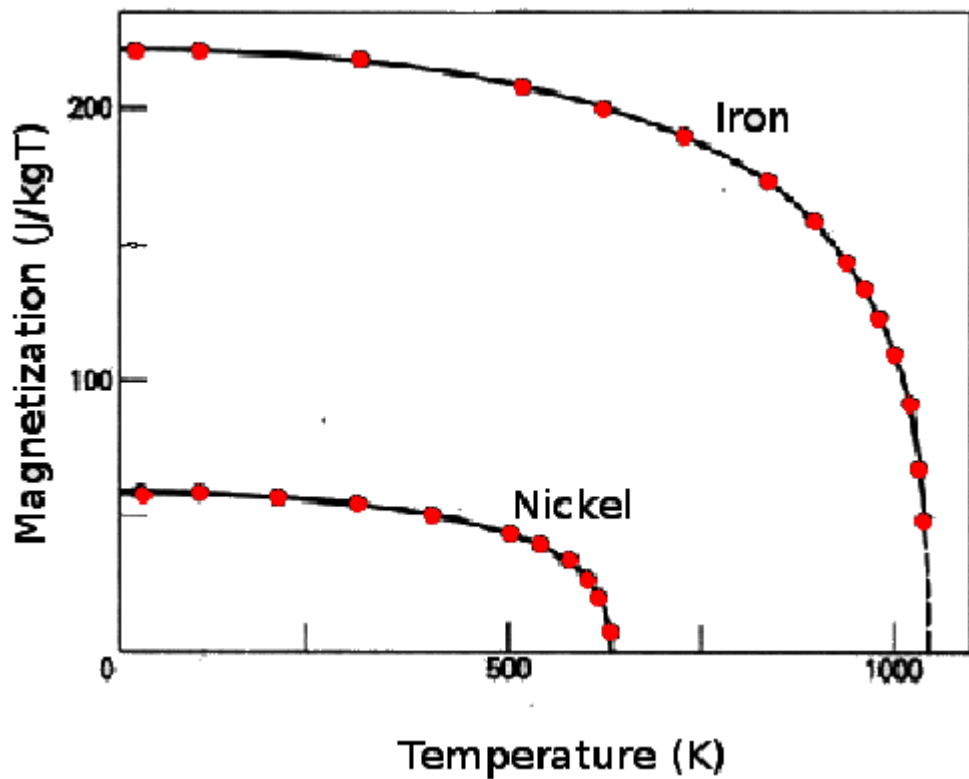


- A Fermi surface shift increases the kinetic energy and potential energy falls
- Ferromagnetic transition occurs if $g v > 1$

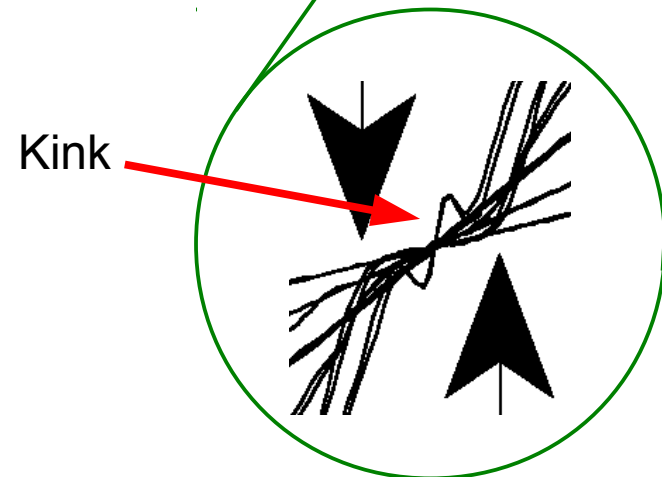
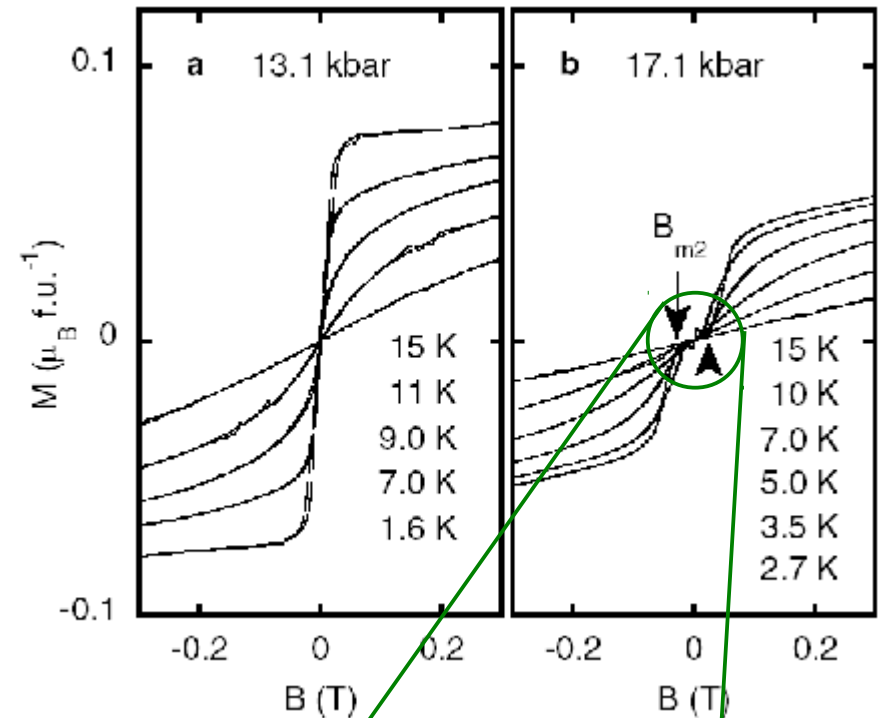
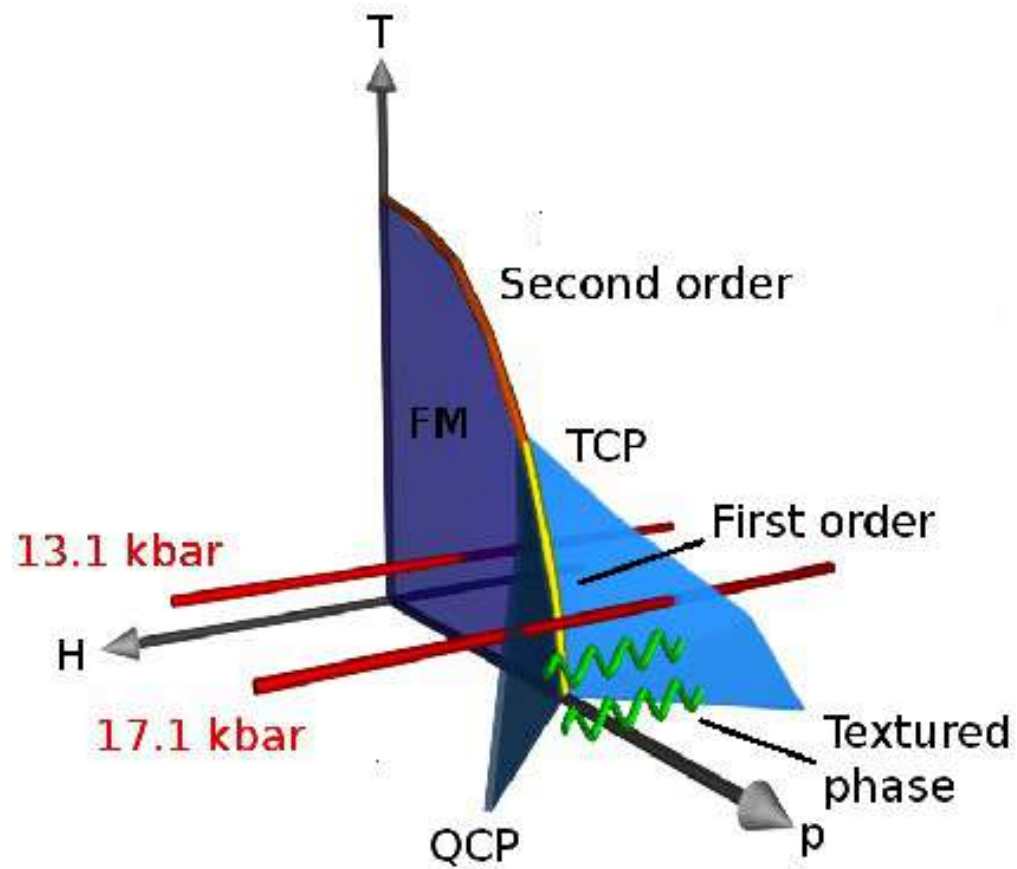
Ferromagnetism in solid state

Second order in iron & nickel

First order in ZrZn_2

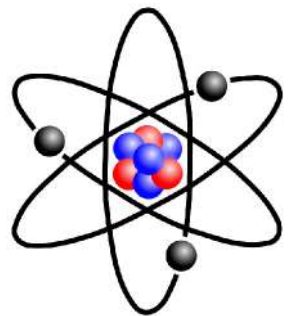


Further phase reconstruction in $ZrZn_2$



Atomic gases: a new forum for ferromagnetism

- A gas of atoms simulates electrons in a solid



${}^6\text{Li}$ atom

$$|F = 1/2, m_F = 1/2\rangle$$



Up spin electron

$$|F = 1/2, m_F = -1/2\rangle$$



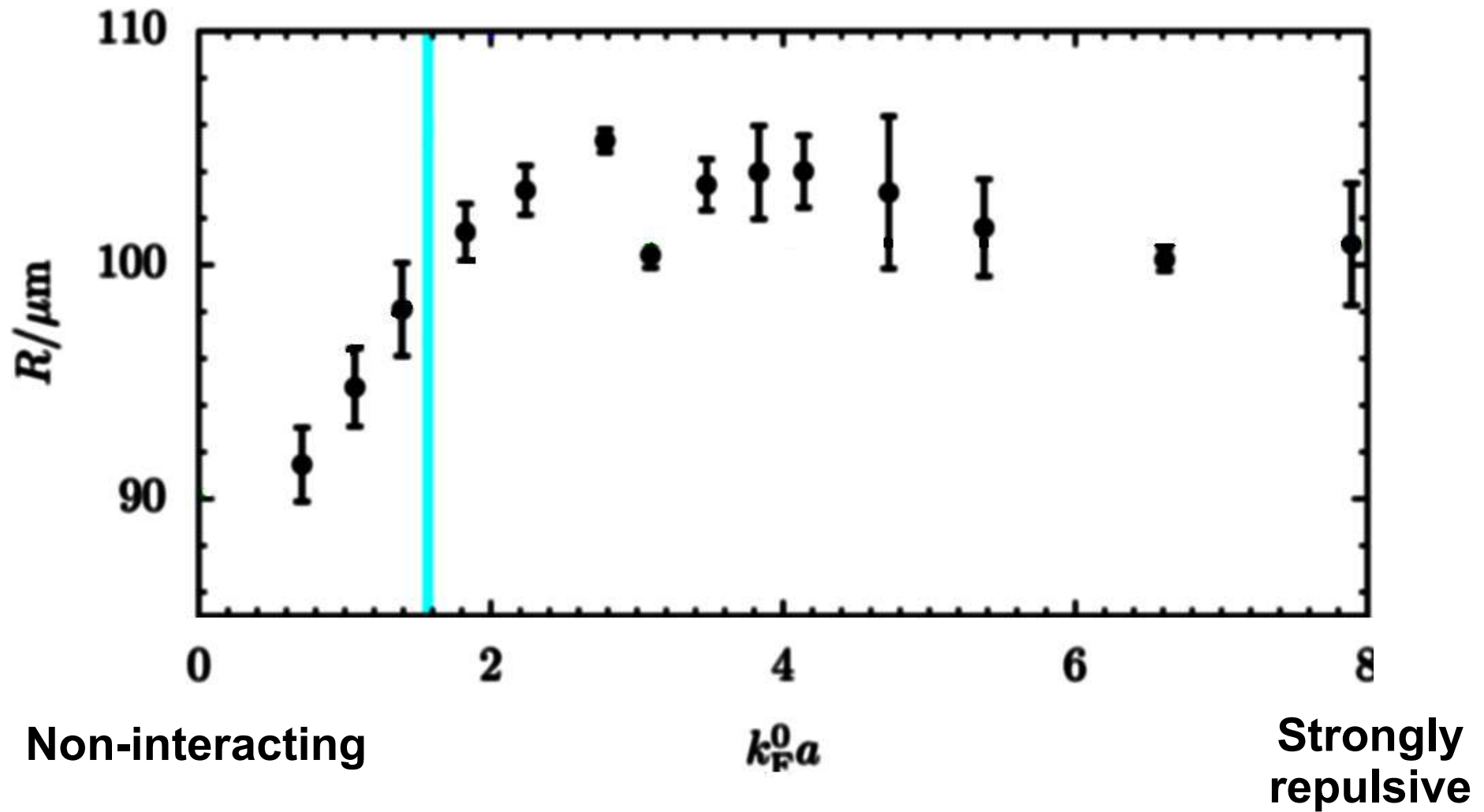
Down spin electron

- Key experimental advantages:
 - Magnetic field controls interaction strength
 - Contact interaction
 - Clean system

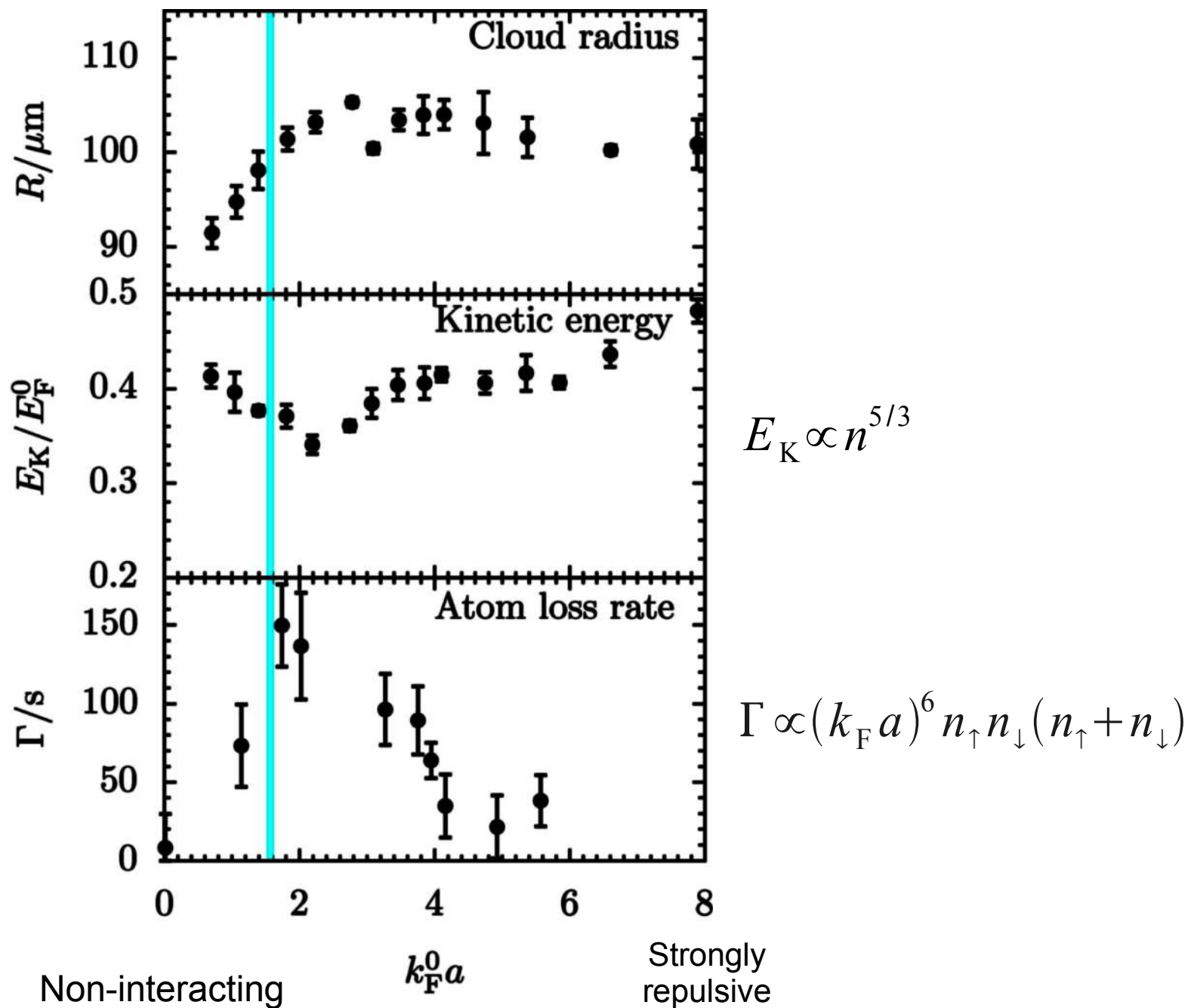
Outline

- Experimental results and mean-field analysis
- Competing many-body instabilities
- Experimental protocols that circumvent atom loss
 - Collective modes within a spin spiral
 - Ferromagnetism with mass imbalance

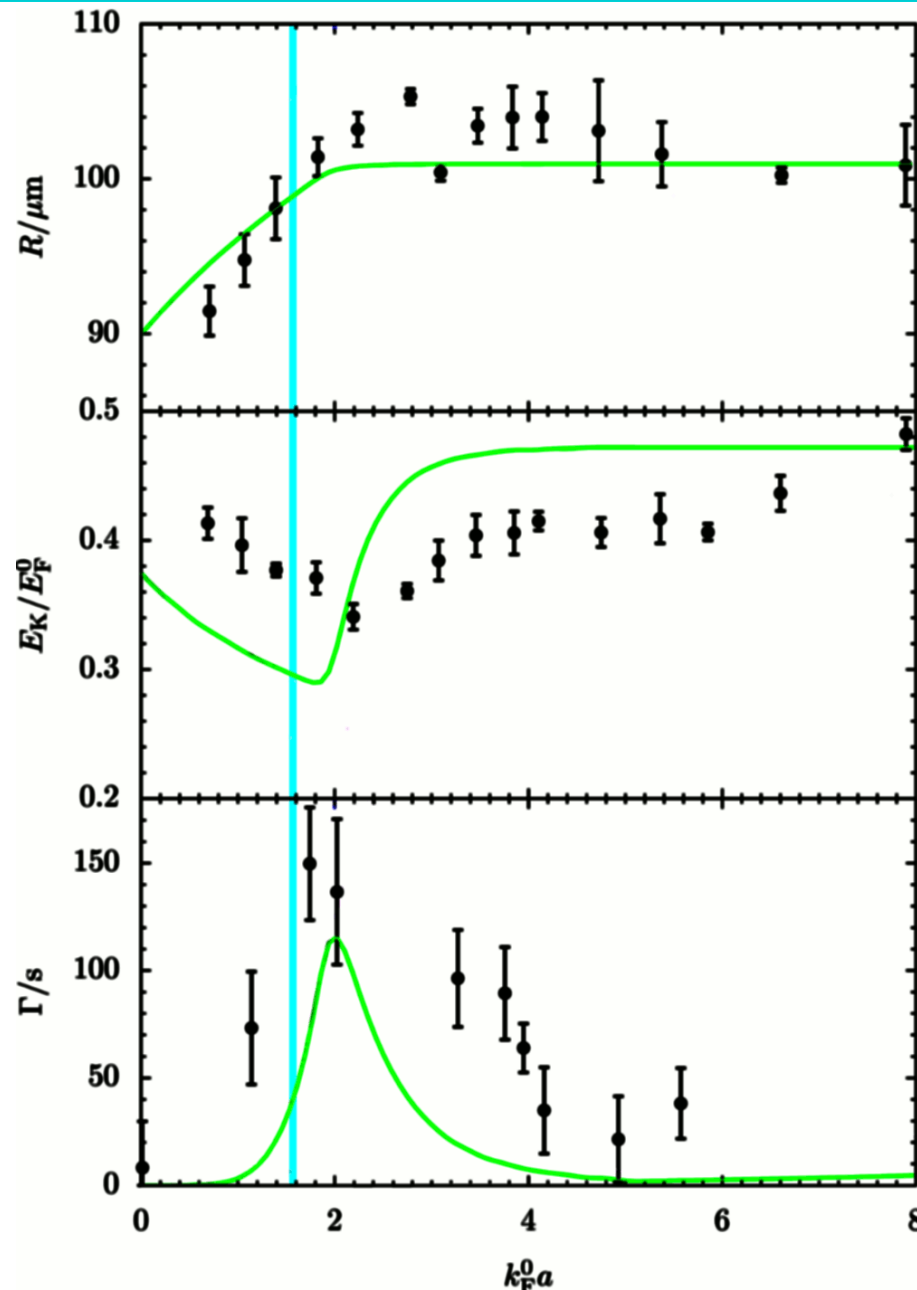
Experimental evidence for ferromagnetism



Further key experimental signatures



Mean-field analysis & consequences of trap

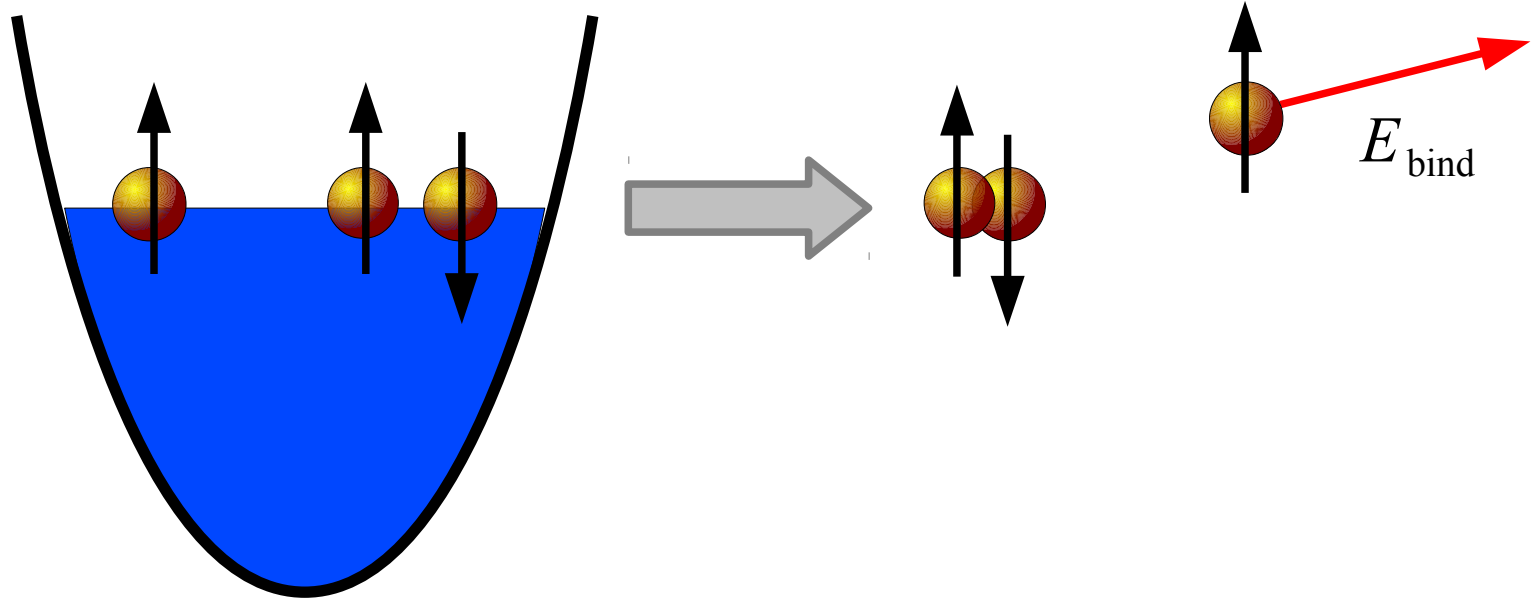


Outline: consequences of atom loss

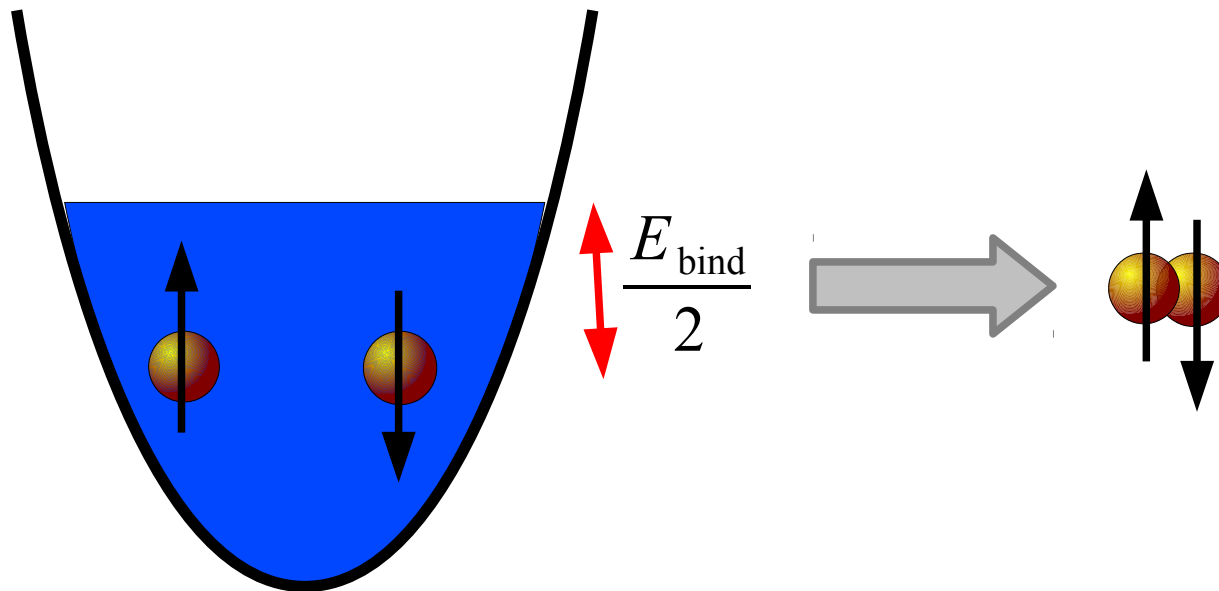
- Experimental results and mean-field analysis
- Competing many-body instabilities
- Experimental protocols that circumvent atom loss
 - Collective modes within a spin spiral
 - Ferromagnetism with mass imbalance

Two versus three-body loss

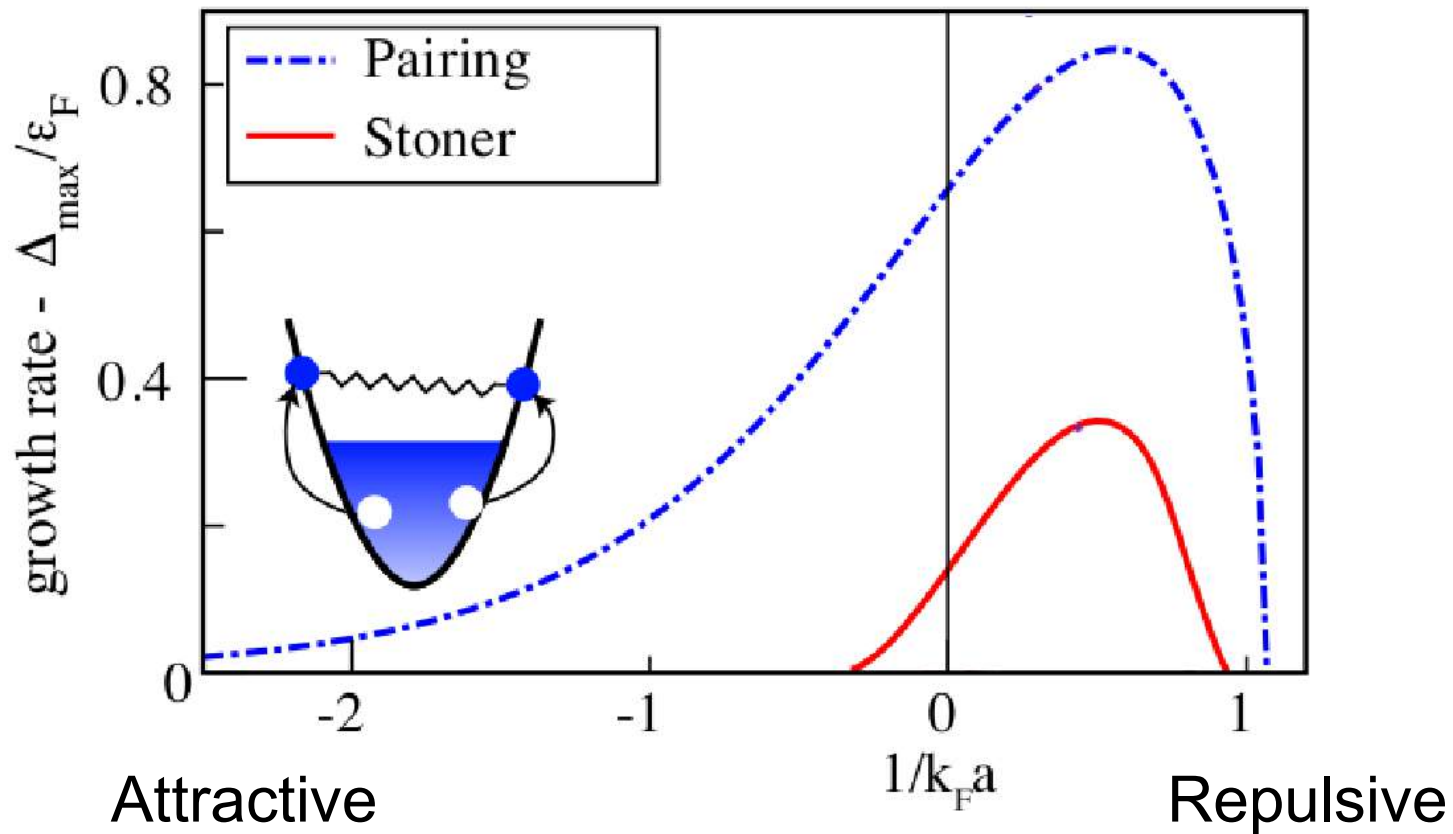
Three-body mechanism



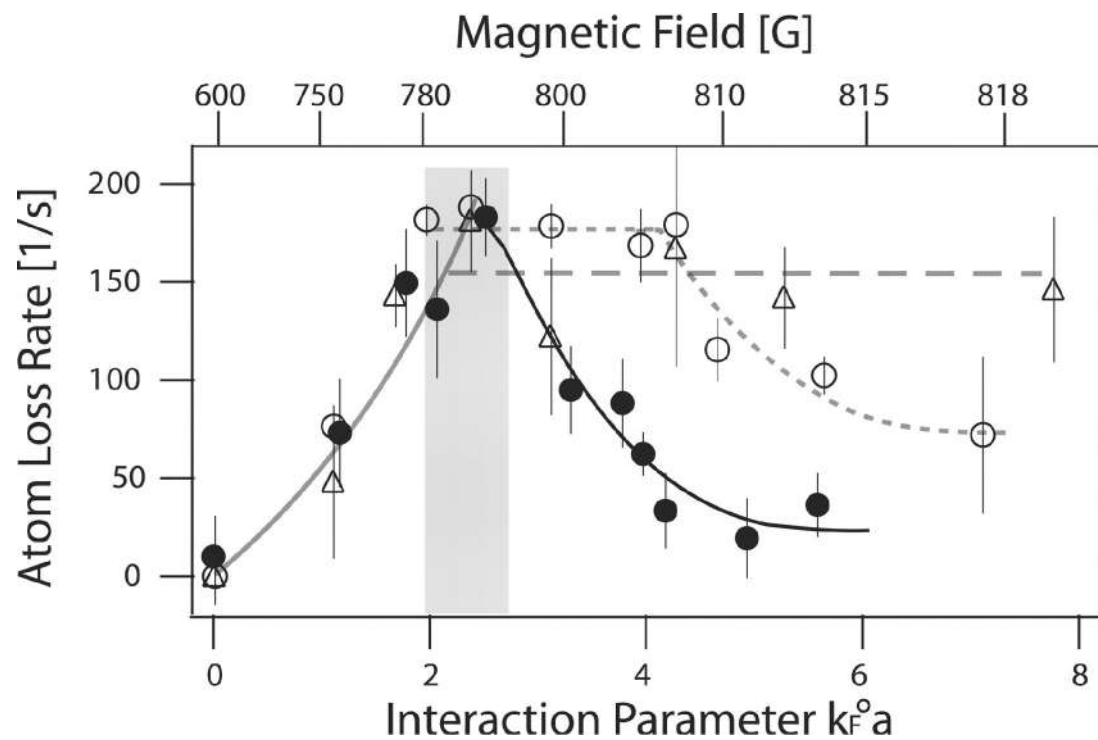
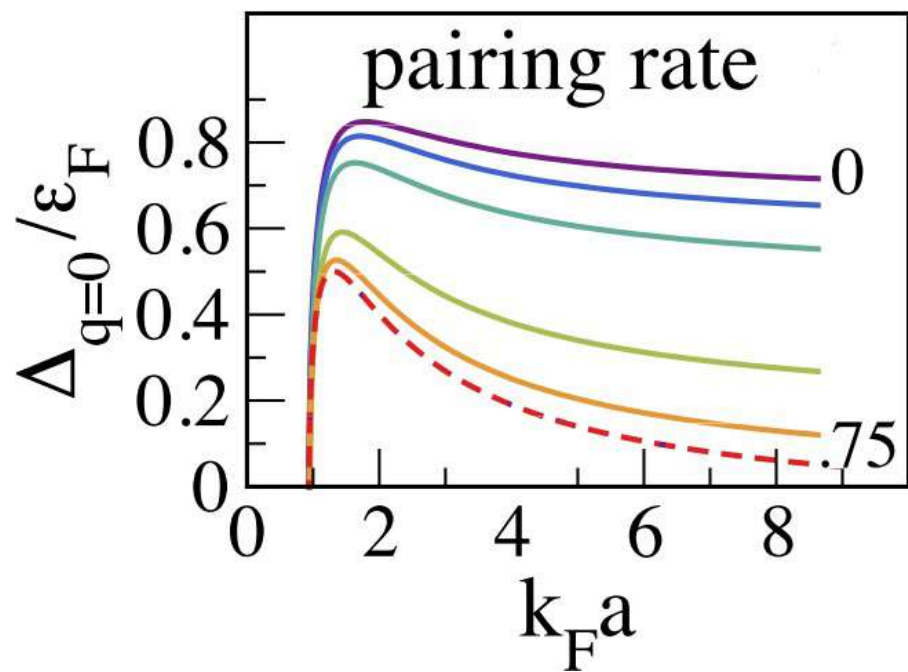
Two-body mechanism



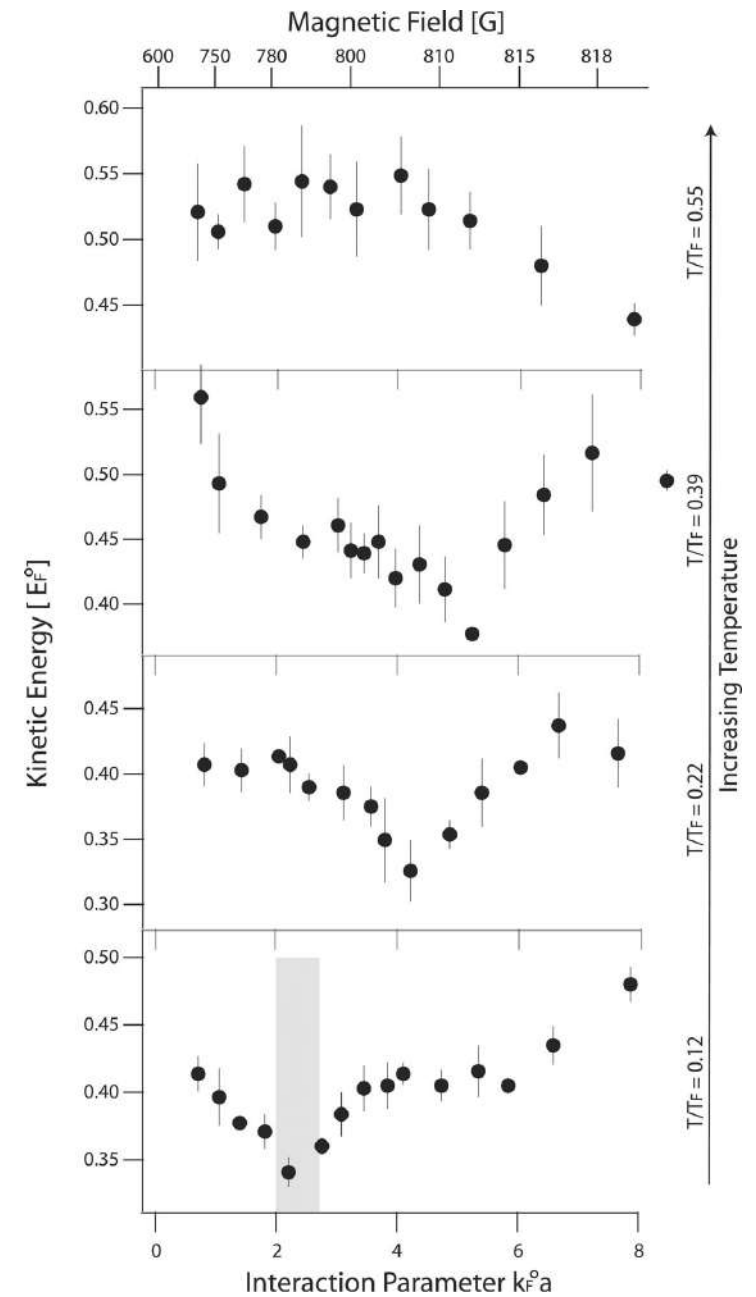
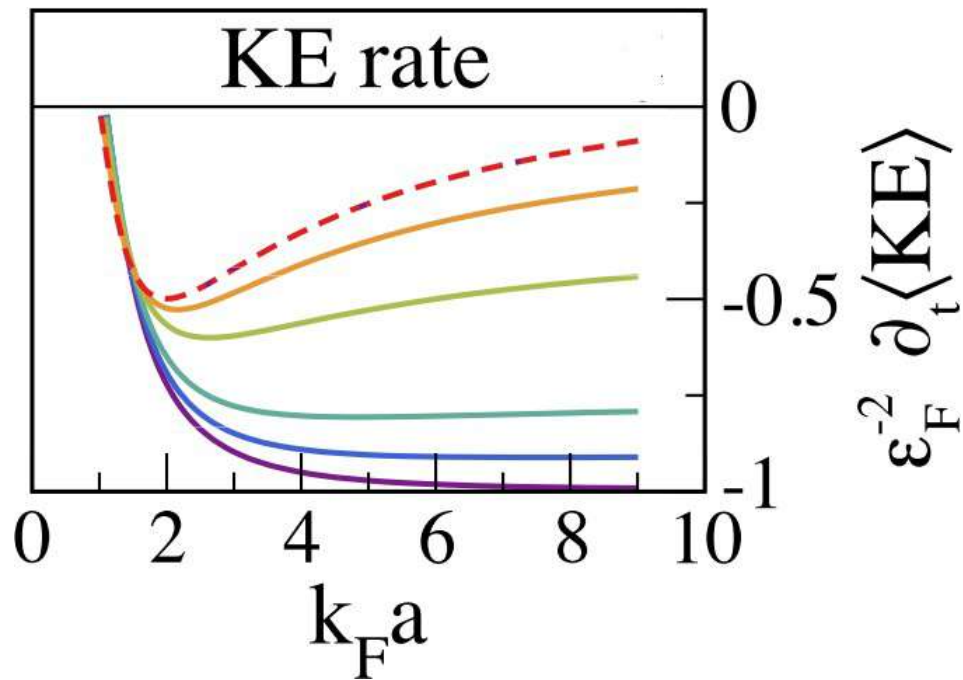
Two-body loss



Two-body loss

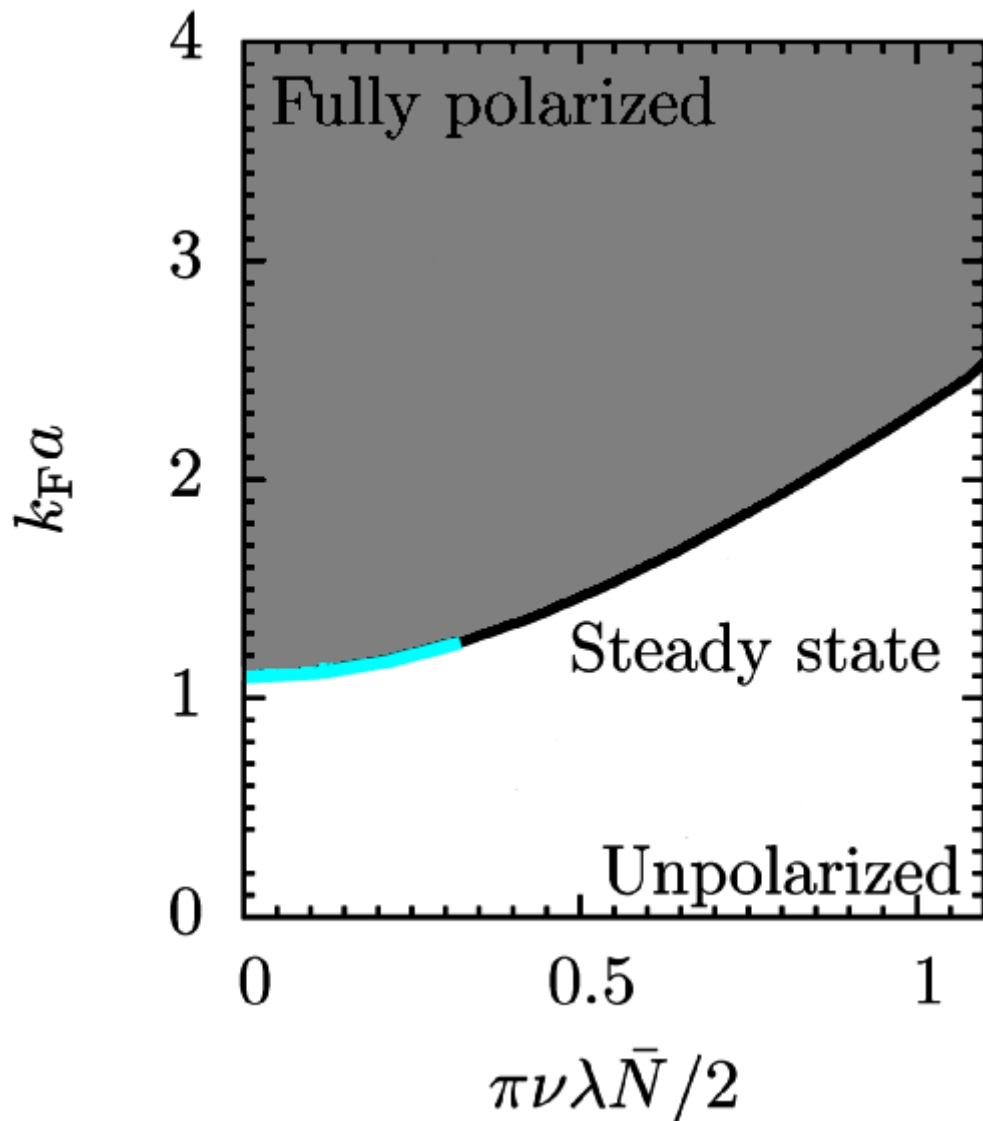


Two-body loss



Phase boundary with atom loss

- Atom loss raises the interaction strength required for ferromagnetism



Outstanding questions

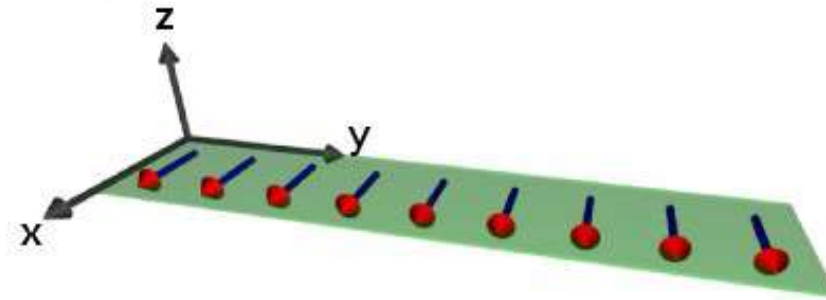
- Two-body competing instability
[Pekker et al., Phys. Rev. Lett. **106**, 050402 (2011)]
- Three-body loss interaction renormalization
[GJC & Altman, Phys. Rev. A **83**, 043618 (2011)]
- Instability of fully polarized gas with single spin impurity
[Zhai, Phys. Rev. A **80**, 051605(R) (2009)]
- Instability of gas with contact interactions (Tan relations)
[Barth & Zwirger, arXiv:1101.5594]

Outline: protocols to avoid loss

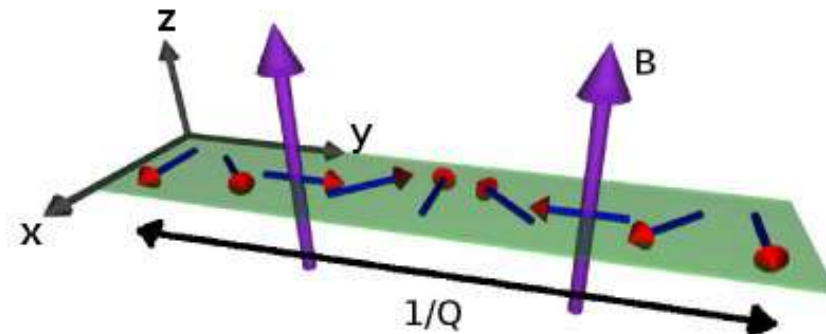
- Experimental results and mean-field analysis
- Competing many-body instabilities
- Experimental protocols that circumvent atom loss
 - Collective modes within a spin spiral
 - Ferromagnetism with mass imbalance

Alternative strategy: spin spiral

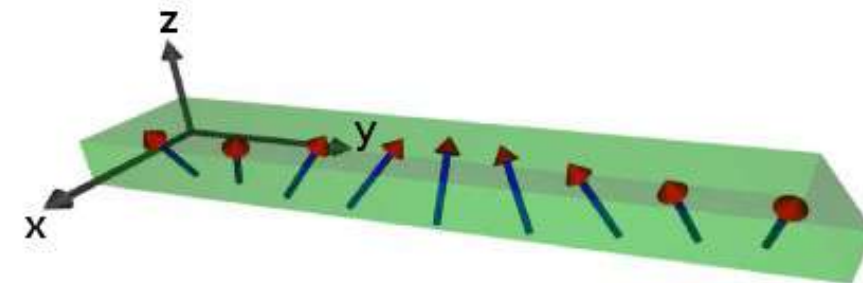
(a) Fully polarized state



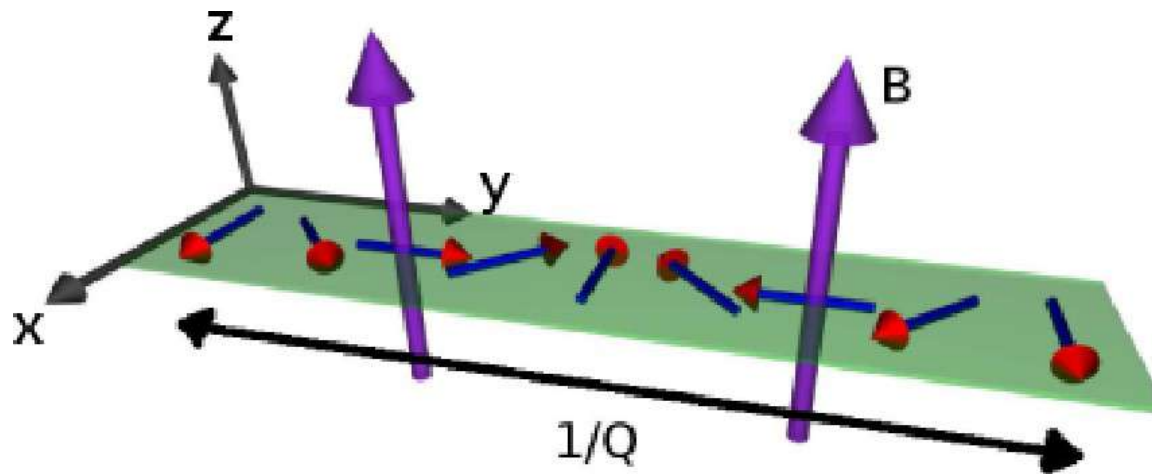
(b) Magnetic field gradient forms spin spiral



(c) Interactions cant the spiral



Heisenberg model



- Describe with Heisenberg Hamiltonian

$$\hat{H} = -J \sum_{i \neq j} \hat{S}_i \cdot \hat{S}_j$$

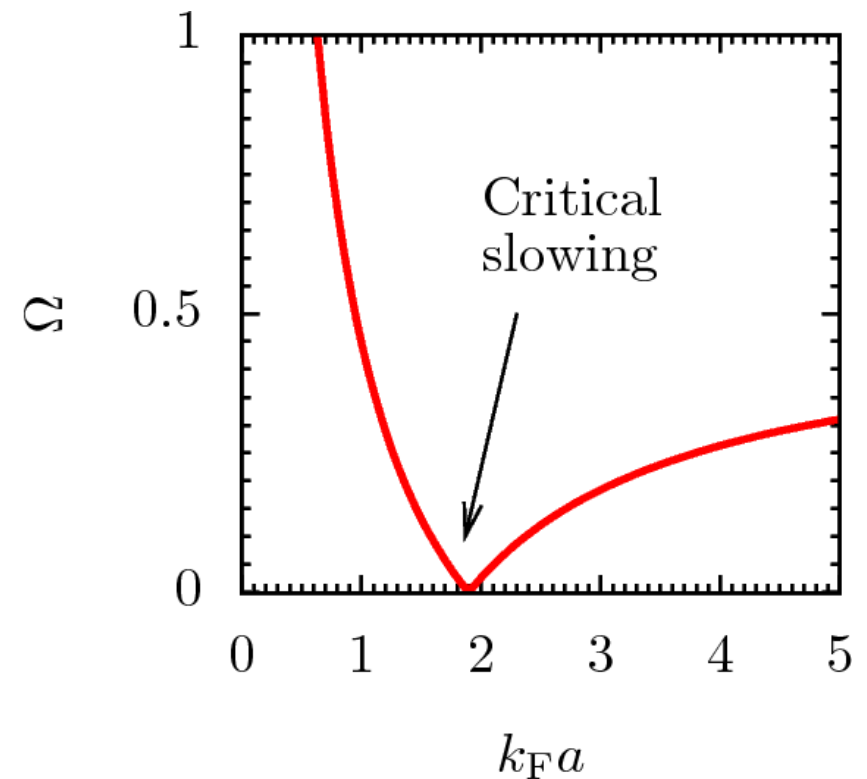
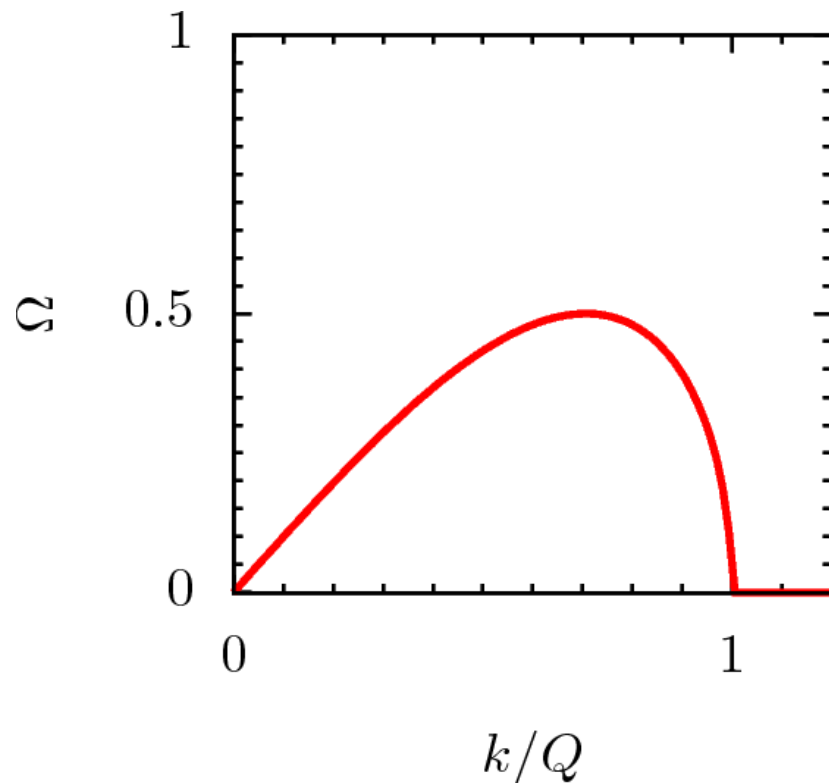
- Dispersion gives mode growth for $k < Q$ independent of sign of J

$$\Omega = \pm J S a^2 k \sqrt{Q^2 - k^2}$$

Spin spiral collective modes

- Exponentially growing collective modes if $k < Q$
[GJC & Altman, PRA **82**, 043603 (2010)]

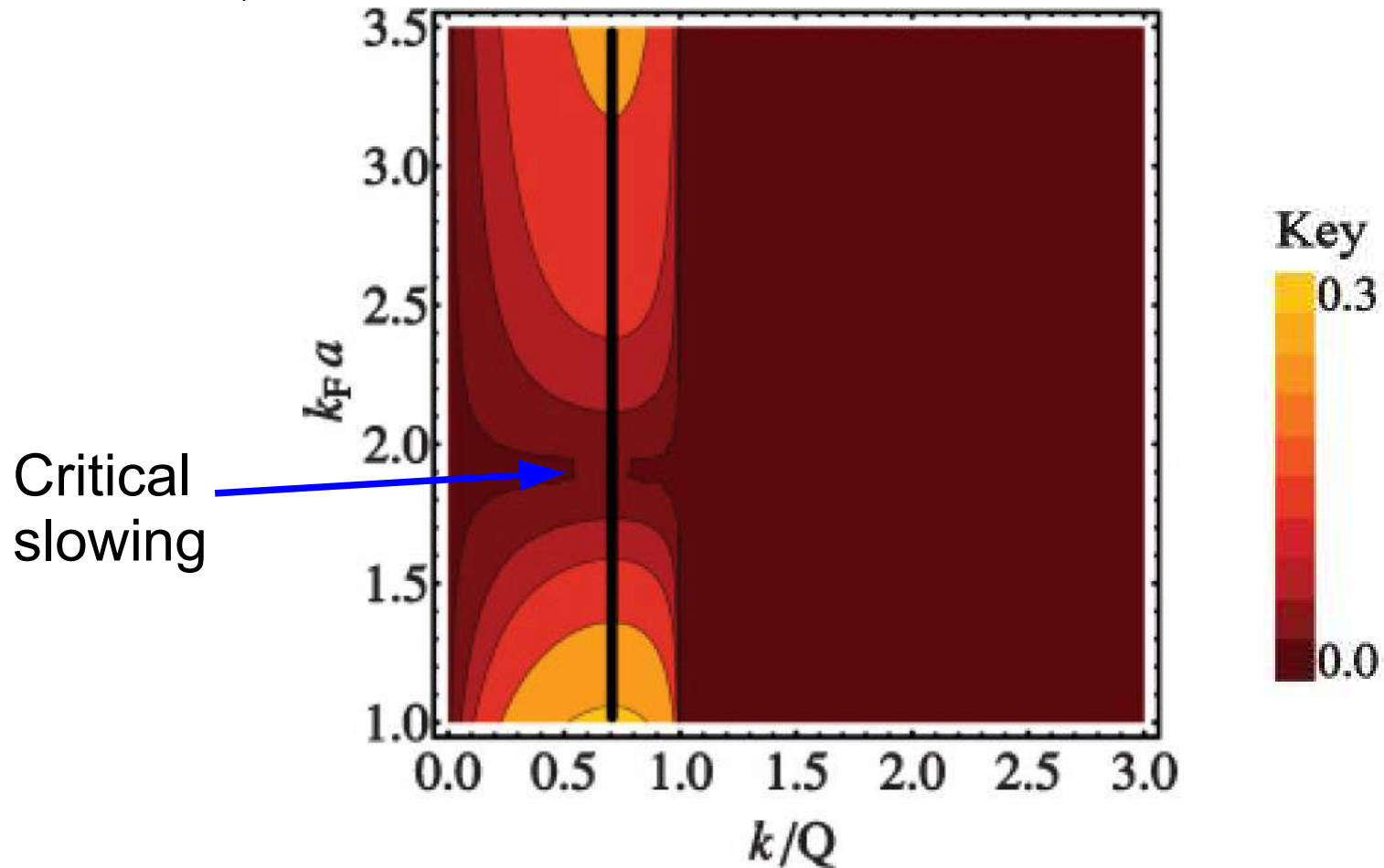
$$\Omega = \pm \left(\frac{1}{2} - \frac{2^{2/3} 3}{5k_F a} \right) k \sqrt{Q^2 - k^2}$$



Spin spiral collective modes

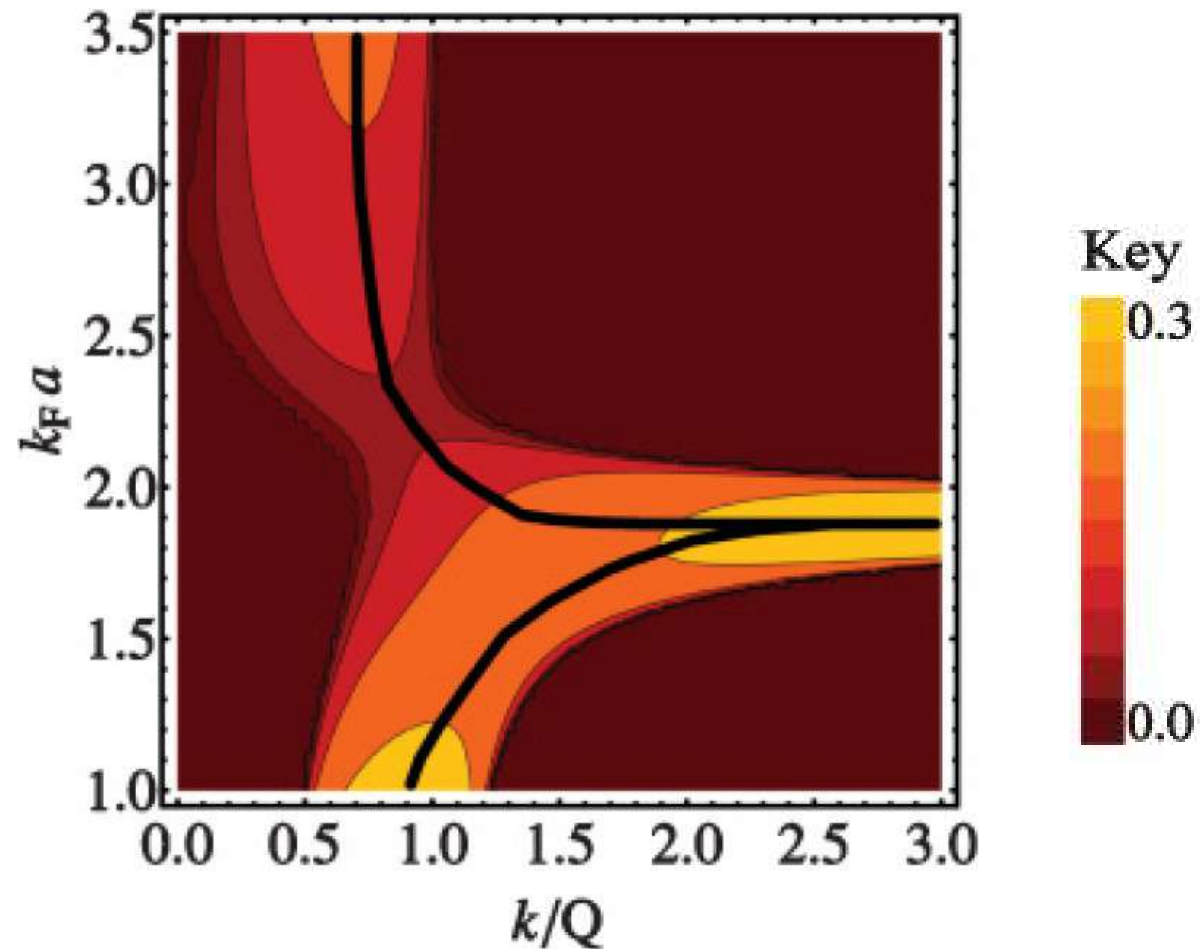
- Exponentially growing collective modes if $k < Q$
[GJC & Altman, PRA **82**, 043603 (2010)]

$$\Omega = \pm \left(\frac{1}{2} - \frac{2^{2/3} 3}{5k_F a} \right) k \sqrt{Q^2 - k^2}$$



Spin spiral collective modes

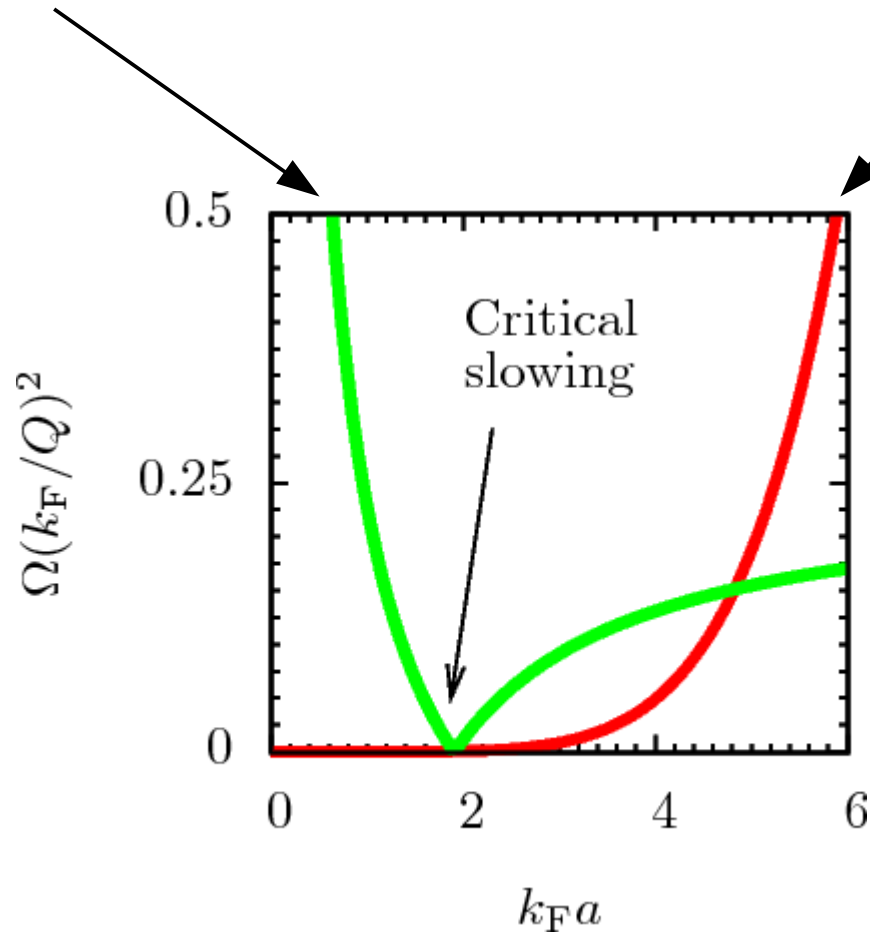
- Critical slowing down is eliminated and a new unstable branch develops
- Long-wavelength instabilities are quenched on the paramagnetic side of the transition



Overcoming loss

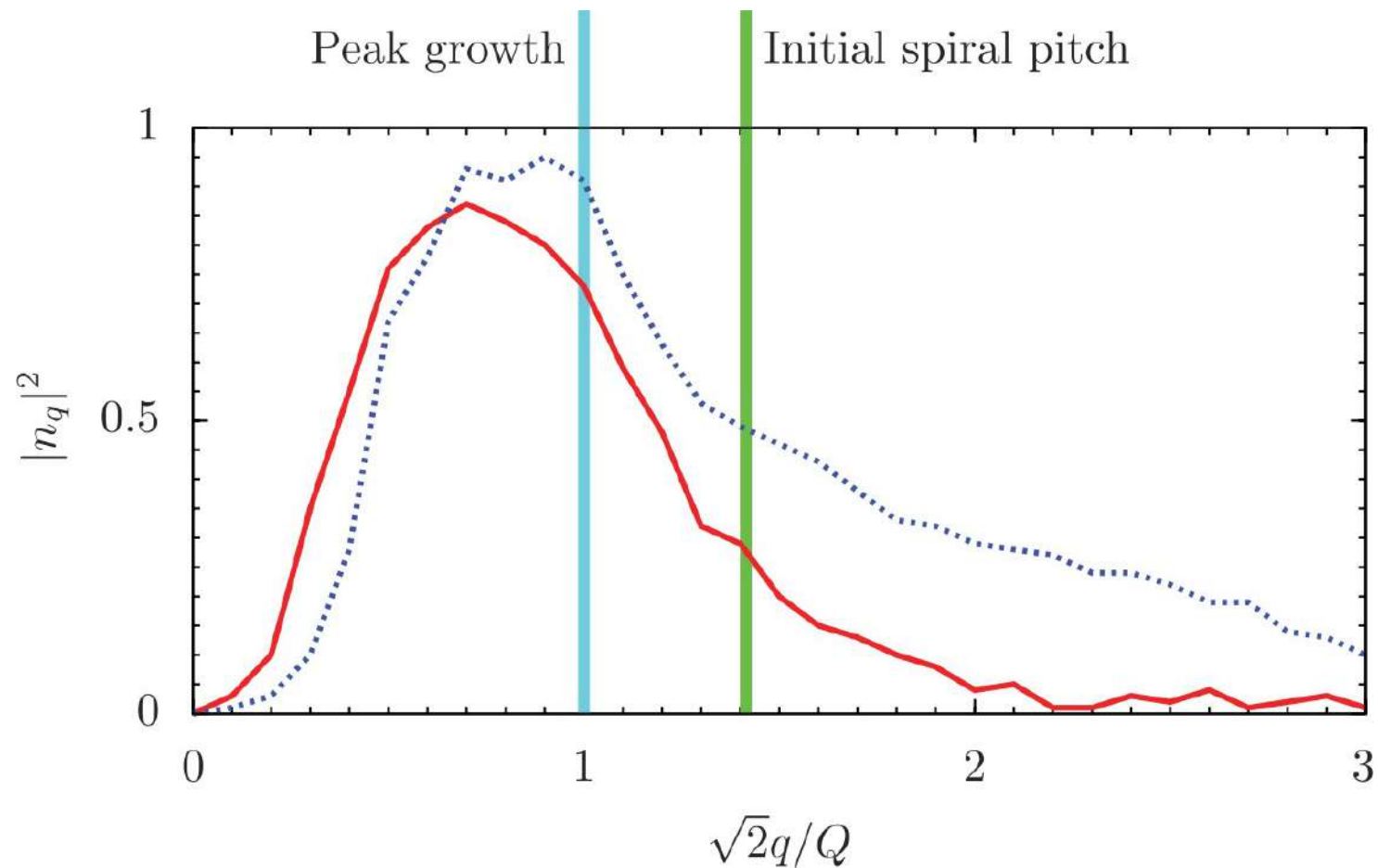
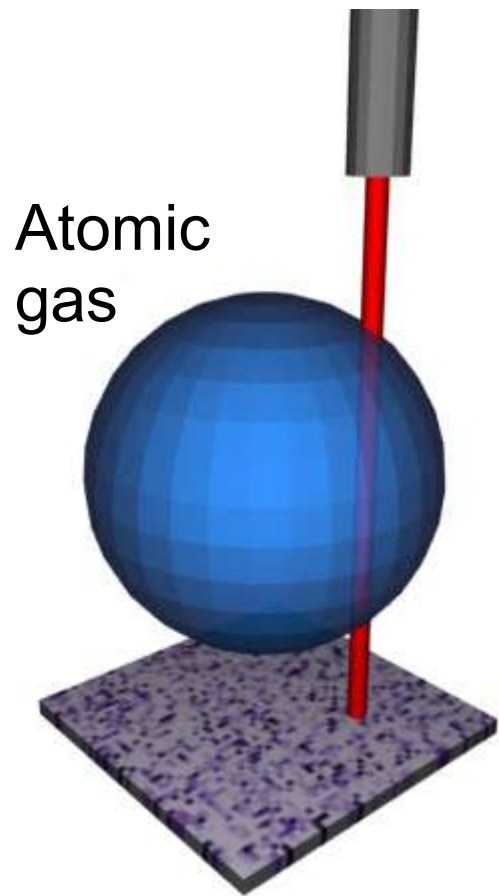
Growth rate
of domains

Three-body
loss rate

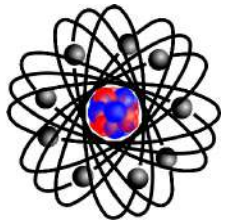


Phase-contrast imaging

- Phase-contrast imaging displays signatures of domain growth
- Domain size fixed across the sample



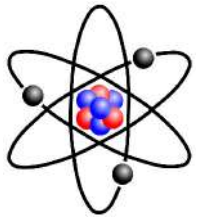
Mass imbalance ferromagnetism



^{40}K atom, $m=40m_0$



Up spin electron



^6Li atom, $m=6m_0$



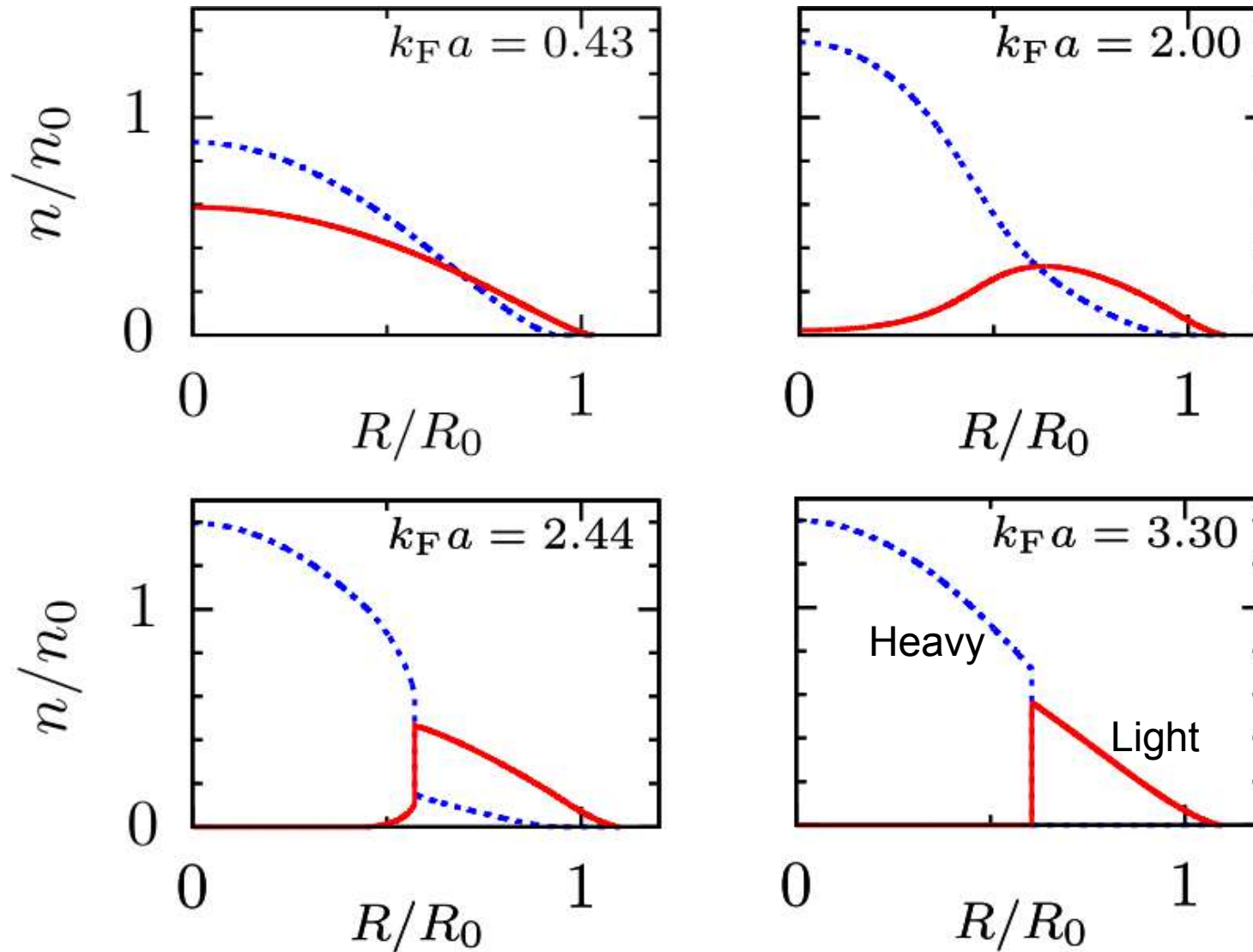
Down spin electron

$$\hat{H} = \sum_k \frac{k^2}{2m_\uparrow} c_{k\uparrow}^\dagger c_{k\uparrow} + \sum_k \frac{k^2}{2m_\downarrow} c_{k\downarrow}^\dagger c_{k\downarrow} + g \sum_{kk'q} c_{k\uparrow}^\dagger c_{k'+q\downarrow}^\dagger c_{k'+q\downarrow} c_{k'\uparrow}$$

- Magnetic moment formed along quantization axis

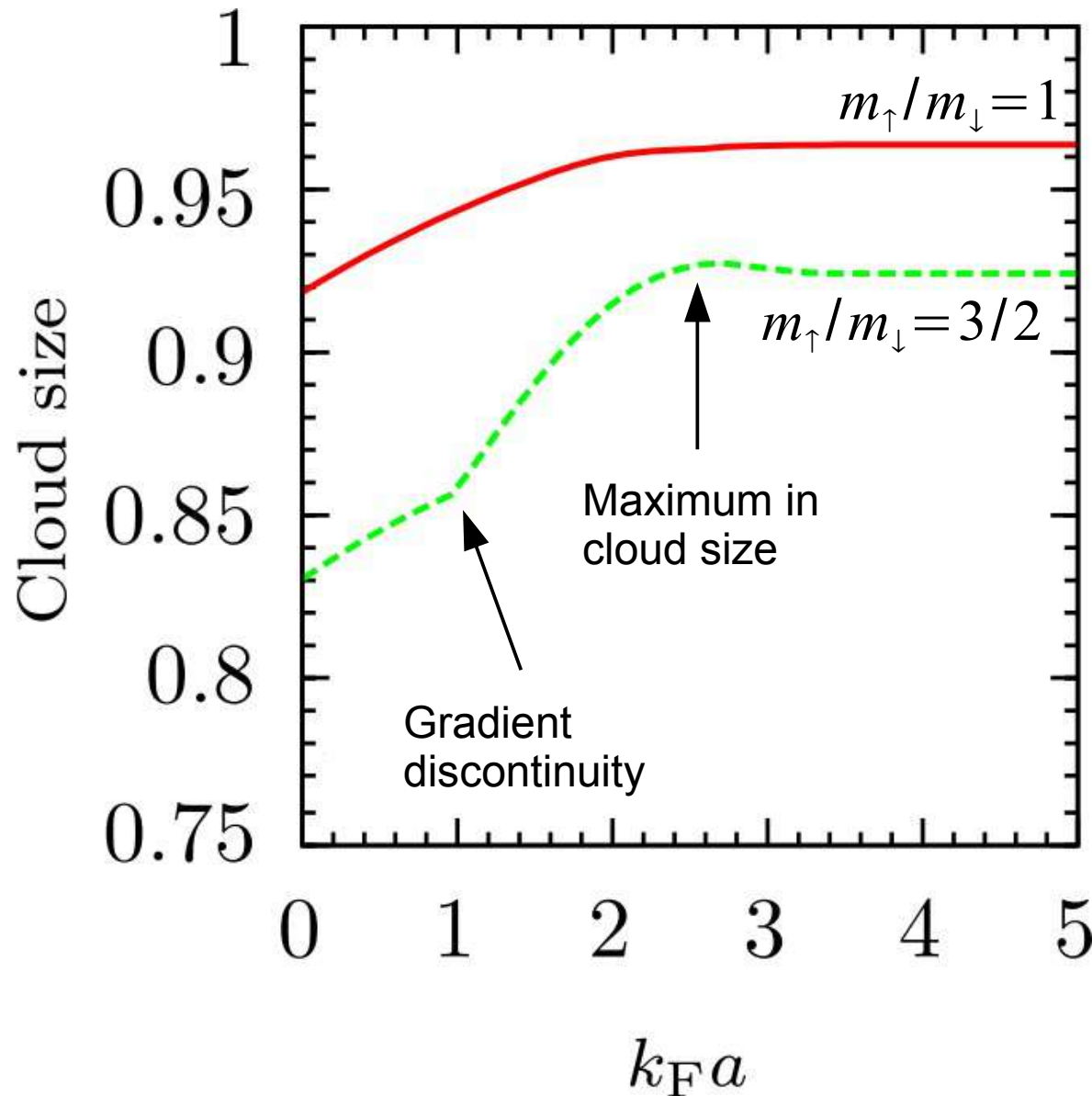
Behavior in a trap

- At zero interaction strength atoms spread all over trap, at high interaction strength light atoms forced to outside



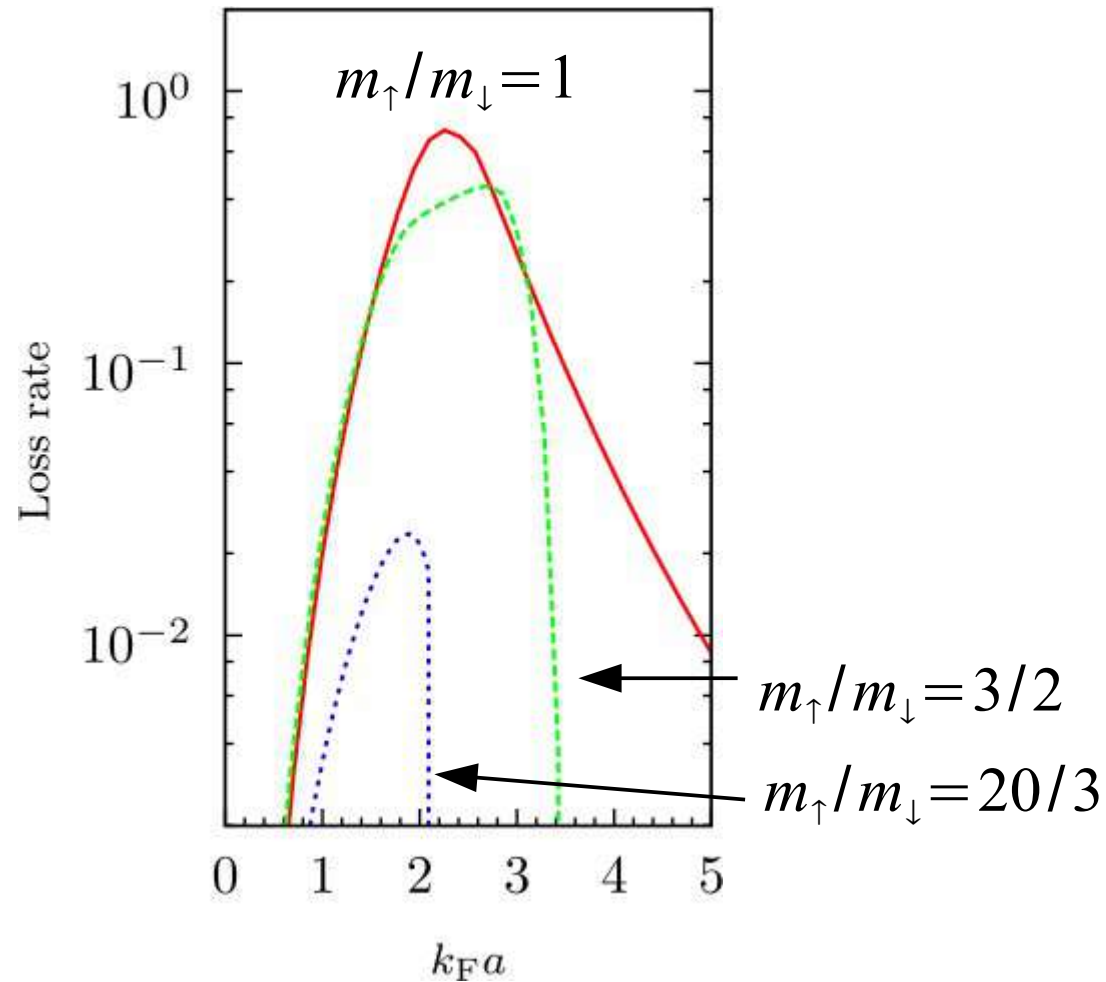
Unique signatures of ferromagnetism

- Expulsion creates unique signatures of ferromagnetism



Reduced three-body losses

- Dramatically reduced three-body loss



Summary

- Equilibrium theory provides a reasonable qualitative description of the transition
- Competing many-body instabilities provide alternative explanation
- Circumvent three-body loss by studying the evolution of a spin spiral
- Suppress losses and give stronger signatures of ferromagnetism by studying mass imbalance
- Answer long-standing questions about solid state ferromagnetism and motivate new research arenas