

# Functional Integrals - a.k.a. field theory

$$Z = \int d^m x e^{-\beta H[m]}$$

cf. path integral saddle pt  $\leftrightarrow$  classical path  
 fluctuations  $\leftrightarrow$  quantum corrections

Gaussian integral in 1 real variable

$$Z_1 = \int_{-\infty}^{\infty} d\varphi e^{-\frac{1}{2G} \varphi^2 + h\varphi}$$

$$= \sqrt{2\pi G} e^{\frac{1}{2} Gh^2}$$

$$\langle \varphi \rangle = \frac{\partial}{\partial h} \log Z_1 = hG$$

$$\frac{1}{Z_1} \frac{\partial Z_1}{\partial h}$$

Define the constant

$$\langle \varphi^r \rangle_c = \left. \frac{\partial^r}{\partial k^r} \right|_{k=0} \log \langle e^{k\varphi} \rangle$$

A student exercise

$$\langle e^{k\varphi} \rangle = \frac{\sqrt{2\pi G} e^{\frac{1}{2}G(h+k)^2}}{\sqrt{2\pi h} e^{\frac{G}{2}h^2}} = e^{\frac{G}{2}(k^2+2hk)}$$

$$\begin{aligned}\langle \varphi^2 \rangle_c &= \partial_k^2 \Big|_{k=0} \text{ by } \langle e^{k\varphi} \rangle \\ &= \langle \varphi^2 \rangle - \langle \varphi \rangle^2 \\ &= \partial_k^2 \Big|_{k=0} \frac{G}{2}(k^2+2hk) \xrightarrow{\text{Var}} \text{2nd cumulant} \\ &= G\end{aligned}$$

N.B. Gaussian ensemble

$$\langle \varphi^r \rangle_c = 0 \quad r > 2.$$

Many real variables

$$Z_N = \int_{-\infty}^{\infty} \frac{N}{\prod_{i=1}^N} d\varphi_i e^{-\frac{1}{2} \varphi^T \tilde{G}^{-1} \varphi + h \cdot \varphi}$$

real symmetry  
⇒ Possible to diagonalise  
the unitary matrix.

$$u \bar{G}^{-1} \bar{u}^T = \bar{\hat{G}}^{-1}$$

↑  
eigenvectors

diagonal matrix of  
eigenvalues

$$\frac{1}{2} \varphi^T \bar{G}^{-1} \varphi - \underline{h} \cdot \varphi = \frac{1}{2} \underbrace{\varphi^T}_{\underline{x}^T} \underbrace{\bar{u}^T u \bar{G}^{-1} \bar{u}^T u}_{\bar{\hat{G}}^{-1}} \underbrace{\varphi}_{\underline{x}} - \underline{h}^T \bar{u}^T u \varphi$$

$$= \frac{1}{2} \underline{x}^T \bar{\hat{G}}^{-1} \underline{x} - \frac{1}{2} \left[ \underline{h}^T \bar{u}^T \underline{x} + \underline{x}^T u \underline{h} \right]$$

$$\text{Let } \underline{\hat{x}}^T = \underline{\hat{x}}^T + \underline{h}^T \bar{u}^T \bar{\hat{G}}$$

$$\underline{x} = \underline{\hat{x}} + \bar{\hat{G}} u \underline{h}$$

$$\begin{aligned} \frac{1}{2} \varphi^T \bar{G}^{-1} \varphi - \underline{h} \cdot \varphi &= \frac{1}{2} \left[ \underline{\hat{x}}^T + \underline{h}^T \bar{u}^T \bar{\hat{G}} \right] \bar{\hat{G}}^{-1} \left[ \underline{\hat{x}} + \bar{\hat{G}} u \underline{h} \right] \\ &\quad - \frac{1}{2} \left[ \underline{h}^T \bar{u}^T (\underline{\hat{x}} + \bar{\hat{G}} u \underline{h}) \right. \\ &\quad \left. + (\underline{\hat{x}}^T + \underline{h}^T \bar{u}^T \bar{\hat{G}}) u \underline{h} \right] \end{aligned}$$

$$= \frac{1}{2} \underline{\hat{x}}^T \bar{\hat{G}}^{-1} \underline{\hat{x}} - \frac{1}{2} \underline{h}^T \bar{u}^T \bar{\hat{G}} u \underline{h}$$

$$= \frac{1}{2} \tilde{x}^T \tilde{G}^{-1} \tilde{x} - \frac{1}{2} \underline{h}^T G \underline{h}$$

$$Z_N = \int_{-\infty}^{\infty} \prod_i d\chi_i e^{-\frac{1}{2} \underline{x}^T \tilde{G}^{-1} \underline{x} + \frac{1}{2} \underline{h}^T G \underline{h}}$$

$$= \sqrt{\det(2\pi G)} e^{\frac{1}{2} \underline{h}^T G \underline{h}}$$

Set of decoupled Gaussians  $\det = \prod_i \lambda_i$

$$\langle \varphi_i \dots \varphi_j \rangle_c = \partial_{k_i} \dots \partial_{k_j} \Big|_{\underline{k}=\underline{0}} \log \langle e^{\underline{k} \cdot \underline{\varphi}} \rangle$$

If Gaussian

$$\langle e^{\underline{k} \cdot \underline{\varphi}} \rangle = \exp \left[ \frac{1}{2} \underline{k}^T G \underline{k} + \underline{h}^T G \underline{k} \right]$$

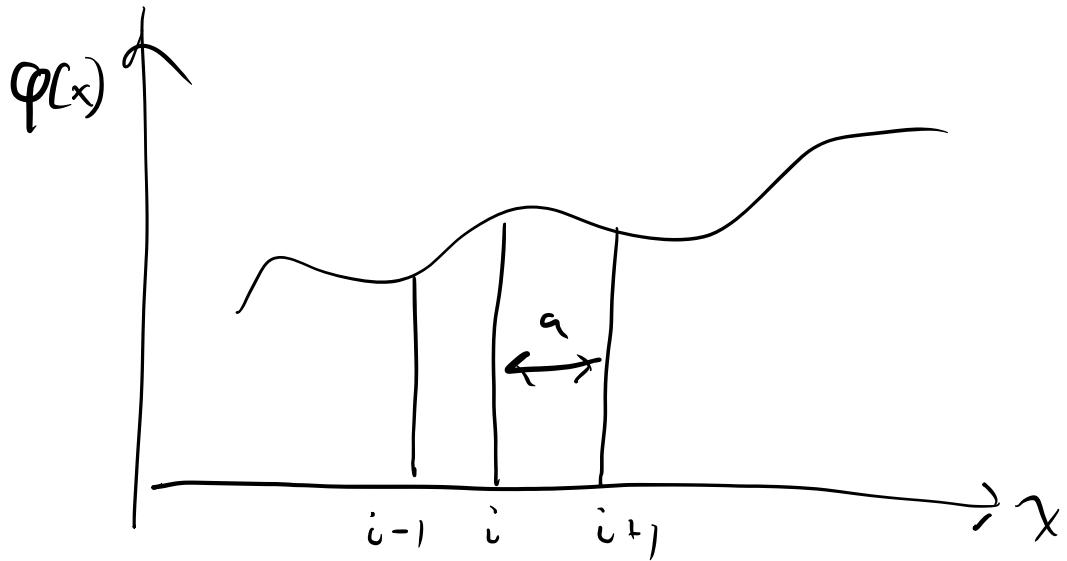
$$\langle \varphi_i \rangle_c = \sum_j G_{ij} h_j = G_{ij} h_j$$

$$\langle \varphi_i \varphi_j \rangle_c = G_{ij}$$

If  $A = g \cdot \varphi$

$$\langle e^A \rangle = e^{\langle A \rangle_c + \frac{1}{2} \langle A^2 \rangle_c}$$

Taking the continuum limit  $\rightsquigarrow$  gaussian function integral.



$$N \rightarrow \infty, a \rightarrow 0$$

$$\varphi_i \rightarrow \varphi(x)$$

↓  
label

$$\bar{G}_{ij}^{-1} \rightarrow \bar{G}^{-1}(\underline{x}, \underline{x}') \equiv \langle \underline{x} | \hat{G}^{-1} | \underline{x}' \rangle$$

↑  
operator

$$Z = \int d\underline{x} \varphi(\underline{x}) \exp \left[ - \int d\underline{x} \int d\underline{x}' \frac{1}{2} \varphi(\underline{x}) \bar{G}^{-1}(\underline{x}, \underline{x}') \varphi(\underline{x}') \right. \\ \left. + \int d\underline{x} h(\underline{x}) \varphi(\underline{x}) \right]$$

$$\alpha (\det G)^{\frac{1}{2}} \exp \left[ \frac{1}{2} \int d^d \underline{x} \int d^d \underline{x}' h(\underline{x}) G(\underline{x}, \underline{x}') h(\underline{x}') \right]$$

where  $\int d^d \underline{x}' G(\underline{x}, \underline{x}') G(\underline{x}', \underline{x}'') = \delta^d(\underline{x} - \underline{x}'')$

↑  
Green's function  
"propagator".

Averages

$$\langle \varphi(\underline{x}) \rangle_c = \int d^d \underline{x}' G(\underline{x}, \underline{x}') h(\underline{x}')$$

and  $\langle \varphi(\underline{x}) \varphi(\underline{x}') \rangle_c = G(\underline{x}, \underline{x}')$

↑  
"connected Green's function"

$$Z = \int \mathcal{D}\varphi(\underline{x}) \exp \left[ -\frac{1}{2} \int d^d \underline{x} \left( \frac{\varphi^2}{\xi^2} + (\nabla \varphi)^2 \right) \right]$$

↑  
 $-\varphi \nabla^2 \varphi$

$$= \int \mathcal{D}\varphi(\underline{x}) \exp \left[ -\frac{1}{2} \int d^d \underline{x} \int d^d \underline{x}' \varphi(\underline{x}) \delta^d(\underline{x} - \underline{x}') \right]$$

$$\left. \left( -\nabla'^2 + \xi^{-2} \right) \varphi(\underline{x}') \right] \downarrow G^{-1}(\underline{x}, \underline{x}')$$

$$\text{i.e. } (-\nabla'^2 + \xi^{-2}) G(\underline{x}', \underline{x}) = \delta^d(\underline{x}' - \underline{x})$$

Fourier Transform

$$G(f) = \frac{1}{f^2 + \xi^{-2}}$$

$$\langle \varphi(f) \varphi(f') \rangle_c = (2\pi)^d \delta^d(f + f') G(f)$$