

Quantum Phase Transitions

Dualities

$$\underline{\nabla} \cdot \underline{E} = \frac{\rho}{\epsilon_0}$$

$$\underline{\nabla} \cdot \underline{B} = 0$$

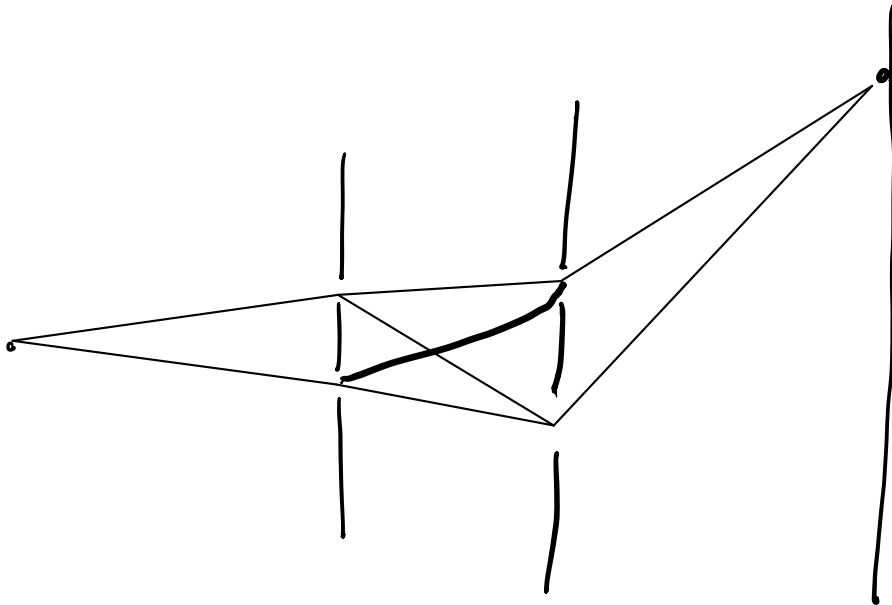
$$\underline{\nabla} \times \underline{E} = -\dot{\underline{B}}$$

$$\underline{\nabla} \times \underline{B} = \mu_0 \underline{J} + \frac{1}{c^2} \dot{\underline{E}}$$

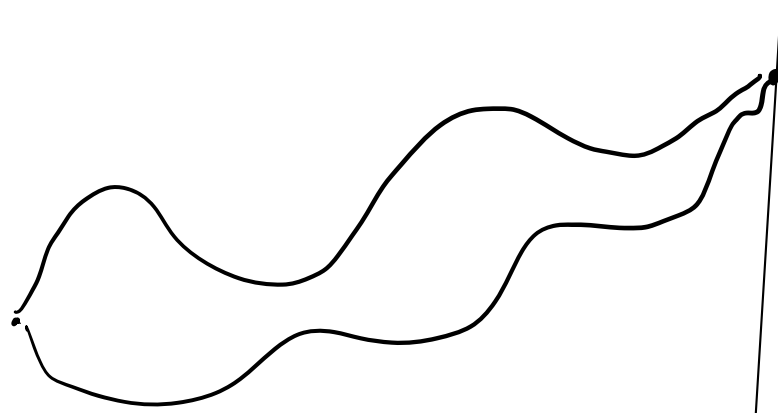
AdS/CFT - holographic

Quantum - Classical mapping

Intuitive Picture



$$e^{iS/\hbar}$$



"sum over histories"
or "path integral"

$$[\hat{x}, \hat{p}] = i\hbar \leftarrow \hbar = 1$$

$$x \delta'(x) = -\delta(x)$$

$$\langle x | [\hat{x}, \hat{p}] | x' \rangle = i\hbar \langle x | x' \rangle = i\hbar \delta(x-x')$$

$$\langle x | \hat{x} \hat{p} - \hat{p} \hat{x} | x' \rangle = (x-x') \langle x | \hat{p} | x' \rangle$$

$$\Rightarrow \langle x | \hat{p} | x' \rangle = -i\hbar \frac{\partial}{\partial x} \delta(x-x')$$

$$\psi_p(x) = \langle x | p \rangle$$

$$\hat{p} | p \rangle = p | p \rangle$$

$$\Rightarrow \langle x | \hat{p} | p \rangle = p \langle x | p \rangle$$

$$I = \int dx' |x'\rangle \langle x|$$

$$\int dx' \langle x | \hat{p} | x' \rangle \langle x' | p \rangle = p \langle x | p \rangle$$

$$\Rightarrow \int dx' -i\hbar \frac{\partial}{\partial x} \delta(x-x') \langle x' | p \rangle = p \langle x | p \rangle$$

$$\Rightarrow -i\hbar \frac{\partial}{\partial x} \int dx' \delta(x-x') \langle x' | p \rangle = p \langle x | p \rangle$$

$$-i\hbar \frac{\partial}{\partial x} \langle x | p \rangle = p \langle x | p \rangle$$

$$\psi_0(x) = \langle x | p \rangle = \frac{e^{\frac{i p x}{\hbar}}}{\sqrt{2\pi\hbar}}$$

Path integral for quantum partition function

$$\hat{H} = \sum_{i=1}^N \frac{p_i^2}{2m} + \hat{V}(\hat{x}_1, \hat{x}_2, \dots, \hat{x}_N)$$

↑
many body potential.

$$[\hat{x}_\mu, \hat{p}_\nu] = i\delta_{\mu\nu} \quad d\text{-dimension}$$

$$Z = \text{tr} e^{-\beta \hat{H}} = \int \prod_{i=1}^N d^d x_i \langle x_1, x_2, \dots, x_N | e^{-\beta \hat{H}} | x_1, \dots, x_N \rangle$$

\uparrow \uparrow
 dX X

$$= \int dX \langle X | e^{-\beta \hat{H}} | X \rangle$$

$X^2 = \sum x_i^2$
 $X \cdot P = \sum x_i \cdot p_i$

$$= \int dX \langle X | e^{-\frac{\beta}{M} \hat{H}} e^{-\frac{\beta}{M} \hat{H}} \dots e^{-\frac{\beta}{M} \hat{H}} | X \rangle$$

\uparrow \uparrow \uparrow \uparrow
 $\int dX_1 |X_1\rangle \langle X_1|$ "2" "3" "N" $\epsilon = \frac{\beta}{M}$

$$\langle X_1 | e^{-\epsilon \hat{H}} | X_2 \rangle \langle X_2 | e^{-\epsilon \hat{H}} | X_3 \rangle \dots \langle X_N | e^{-\epsilon \hat{H}} | X_1 \rangle$$

\uparrow \uparrow \uparrow
 $\int dP_1 |P_1\rangle \langle P_1|$ "2" "N"

key bracket:

$$\langle P_i | e^{-\epsilon \hat{H}} | X_{i+1} \rangle$$

$$= \langle P_i | (1 - \epsilon \hat{H}) | X_{i+1} \rangle$$

$$= \left(1 - \varepsilon \left(\frac{p_i^2}{2m} + V(x_{i+1}) \right) \right) \langle p_i | x_{i+1} \rangle$$

$$= e^{-\varepsilon \left(\frac{p_i^2}{2m} + V(x_{i+1}) \right)} e^{-i p_i \cdot x_{i+1}}$$

So putting back together

$$Z = \int \prod_{i=1}^m dx_i \int \prod_{i=1}^m dp_i \exp \left\{ -i \sum_{i=1}^m p_i \cdot (x_{i+1} - x_i) - \varepsilon \sum_{i=1}^m \left[\frac{p_i^2}{2m} + V(x_i) \right] \right\}$$

Note the periodic b.c.

$$x_{N+1} = x_1$$

Taking the continuum limit

$$x_i \rightarrow x(\tau) \quad \tau = (i-1)\varepsilon$$

and we find the functional integral

$$Z = \int_{x(\beta) = x(0)} \mathcal{D}x(\tau) \int \mathcal{D}p(\tau) e^{-\int_0^\beta d\tau \left(i p(\tau) \partial_\tau x + \frac{p^2}{2m} + V(x(\tau)) \right)}$$

Integrate out momenta fields

$$Z = \int \mathcal{D}\underline{x}_i(\tau) e^{-H[\underline{x}_i(\tau)]}$$

$$\underline{x}_i(\beta) = \underline{x}_i(0)$$

$$H[\underline{x}_i(\tau)] = \int_0^\beta d\tau \sum_{i=1}^N \frac{m |\partial_\tau x_i|^2}{2} + V[x_i(\tau)]$$

↓
classical Hamiltonian

Quantum - Classical Mapping

A d -dimensional quantum system can be mapped onto a $d+1$ dimensional classical system.

Extra dimension is the imaginary time τ .

Key correspondences of the Quantum-Classical mapping

Correspondence	Quantum	Classical
Hamiltonian $\hat{H} \leftrightarrow \beta H$	$\hat{H} = \sum_i \frac{\hat{p}_i^2}{2M} + V[\hat{x}_i]$, where $[\hat{x}_i, \hat{p}_j] = i\delta_{ij}$	$\beta H = \int_0^\beta d\tau \left[\sum_i \frac{M}{2} \left(\frac{\partial m_i(\tau)}{\partial \tau} \right)^2 + V[m_i(\tau)] \right]$
Order Parameter $\hat{x}_i \leftrightarrow m_i(\tau)$	\hat{x}_i	$m_i(\tau)$
Partition function $Z_Q = Z_C$	$Z_Q = \sum_n \langle n e^{-\beta \hat{H}} n \rangle$	$Z_C = \int_{m_i(0)=m_i(\beta)} \mathcal{D}m_i(\tau) e^{-\beta H[m_i(\tau)]}$
Expectation Value $\langle \hat{x}_i \rangle = \langle m_i(0) \rangle$	$\langle \hat{x}_i \rangle = \frac{1}{Z_Q} \sum_n \langle n e^{-\beta \hat{H}} \hat{x}_i n \rangle$	$\langle m_i(0) \rangle = \frac{1}{Z_C} \int_{m_i(0)=m_i(\beta)} \mathcal{D}m_i(\tau) m_i(0) e^{-\beta H[m_i(\tau)]}$
Correlator $\langle e^{\hat{H}\tau} \hat{x}_j e^{-\hat{H}\tau} \hat{x}_i \rangle = \langle m_j(\tau) m_i(0) \rangle$, where $\tau > 0$	$\langle e^{\hat{H}\tau} \hat{x}_j e^{-\hat{H}\tau} \hat{x}_i \rangle = \frac{1}{Z_Q} \sum_n \langle n e^{-\beta \hat{H}} e^{\hat{H}\tau} \hat{x}_j e^{-\hat{H}\tau} \hat{x}_i n \rangle$	$\langle m_j(\tau) m_i(0) \rangle = \frac{1}{Z_C} \int_{m_i(0)=m_i(\beta)} \mathcal{D}m_i(\tau) m_j(\tau) m_i(0) e^{-\beta H[m_i(\tau)]}$

Notes:

Example

S.H.O. \leftrightarrow 1d string model

$$\begin{aligned} \hat{H} &= \frac{\hat{p}^2}{2M} + \frac{1}{2} M \omega^2 \hat{x}^2 \leftrightarrow \beta H[m(x)] \\ &= \int_0^\beta dx \frac{1}{2} M \left(\frac{\partial m}{\partial \tau} \right)^2 + \frac{1}{2} M \omega^2 m^2 \end{aligned}$$

By the correspondence

$$K = M \quad \text{and} \quad t = M \omega^2$$

Recall $K \xi^{-2} = t$

$$\begin{aligned} \text{Correlation length } \xi &= \sqrt{\frac{K}{t}} = \sqrt{\frac{M}{M \omega^2}} \\ &= \frac{1}{\omega} = \frac{K}{\Delta E} \end{aligned}$$

General result: Correlation length (along imaginary time) is the inverse of the energy gap between the ground state and 1st excited state.