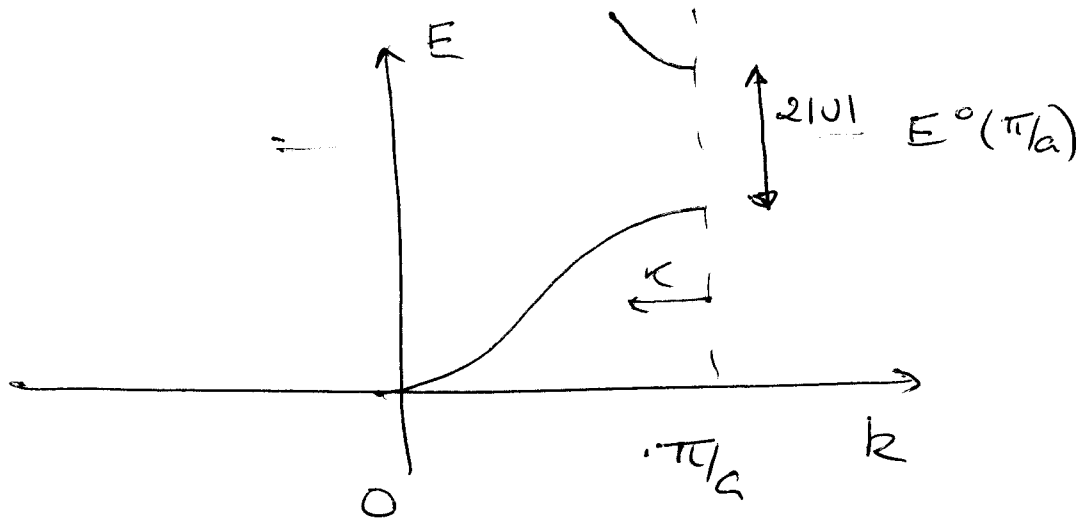


From (6.14)

$$E^{\pm}(k) = \frac{1}{2} \left[ E^0(k) + E^0(k - \frac{2\pi}{a}) \right] \pm \frac{1}{2} \sqrt{\left( E^0(k) - E^0(k - \frac{2\pi}{a}) \right)^2 + 4U^2}$$

where  $E^0(k) = \frac{\hbar^2 k^2}{2m}$



Write  $k = \pi/a + \kappa$ ,  $|\kappa| \ll \pi/a$

$$\begin{aligned} E^0(k) - E^0(k - \frac{2\pi}{a}) &= \frac{\hbar^2}{2m} \left\{ \left( \frac{\pi}{a} + \kappa \right)^2 - \left( \frac{\pi}{a} - \kappa \right)^2 \right\} \\ &= \frac{\hbar^2}{2m} \cdot \frac{4\pi}{a} \cdot \kappa \end{aligned}$$

Then

$$\begin{aligned} E^{\pm} &\approx E^0\left(\frac{\pi}{a}\right) \pm |U| \sqrt{\left[ 1 + \left( \frac{\hbar^2}{2m} \cdot \frac{4\pi}{a} \cdot \frac{\kappa}{2U} \right)^2 \right]} \\ &\approx E^0\left(\frac{\pi}{a}\right) \pm |U| \left[ 1 + \frac{1}{2} \left( \frac{\hbar^2}{2m} \cdot \frac{4\pi}{a} \cdot \frac{\kappa}{2U} \right)^2 + \dots \right] \\ &= E^0\left(\frac{\pi}{a}\right) \pm |U| \pm \frac{1}{2} \frac{\hbar^2 \kappa^2}{m^*} \end{aligned}$$

$$\therefore \frac{1}{m^*} = \frac{1}{m} \cdot \frac{\hbar^2 \pi^2}{m a^2} \cdot \frac{1}{|U|}, \text{ or since } E_{\text{gap}} = 2U$$

$$\frac{m^*}{m} = \frac{E_{\text{gap}}}{4 E^0(\pi/a)}$$

	$E_{\text{gap}} / (m^*/m_e)$
$\ln \text{ Sb}$	15
$\ln \text{ As}$	16.5
$\ln \text{ P}$	19.45

$$4E^0(\pi/a) \sim 4 \cdot \frac{\hbar^2 \pi^2}{2ma^2}$$

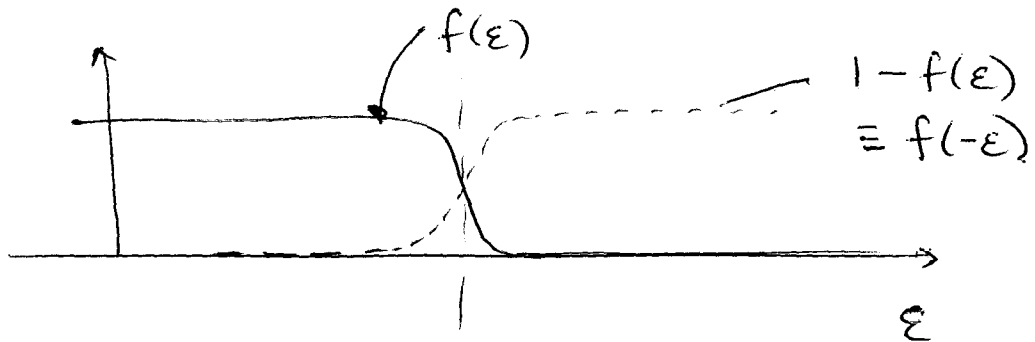
$$a \sim 2.5 \times 10^{-10} \text{ m. (not much variation between materials)}$$

$$\text{So } 4E^0(\pi/a) \sim 20 \text{ eV.}$$

Conclusion: Rough consistency,  
most of the trend of  $m^*$  variation  
accounted for by change of band gap

Simple algebra.

$$\frac{1}{e^x + 1} + \frac{1}{e^{-x} + 1} = \frac{e^{-x} + 1 + e^x + 1}{1 + 1 + e^x + e^{-x}} = 1$$



ionized donors | acceptors at RT.

$$n - p = N_d - N_a \approx n \quad (p \text{ very small})$$

$$n = 2 \times 10^{22} \text{ m}^{-3}$$

$$np = n_i^2 \Rightarrow p = 2.9 \times 10^{16} \text{ m}^{-3}.$$

$$\sigma = ne\mu \quad , \quad R_H = -\frac{1}{n|e|}$$

Effective Hamiltonian for a donor impurity is

$$H = -\frac{\hbar^2}{2m^*} \nabla^2 + \frac{e^2}{4\pi\epsilon\epsilon_0 r}$$

$\Rightarrow$  Hydrogenic, but with.

$$\text{Rydberg} = \frac{m^*}{2\hbar^2} \cdot \left( \frac{e^2}{4\pi\epsilon\epsilon_0} \right)^2 = \left( \frac{m^*}{m} \right) \frac{1}{\epsilon^2} \times 13.6 \text{ eV}$$

$$\text{Bohr radius} = \frac{\epsilon \hbar^2 4\pi\epsilon_0}{m^* e^2} = \frac{\epsilon}{(m^*/m)} 0.529 \times 10^{-10} \text{ m}$$

$$\text{For InSb, Rydberg}^* = \frac{0.015}{18^2} \cdot 13.6 \text{ eV} \\ = 6 \times 10^{-4} \text{ eV}$$

— very small ( $1 \text{ meV} \approx 10 \text{ Kelvin}$ )

$$a_{\text{Bohr}}^* = 60 \text{ nm}$$

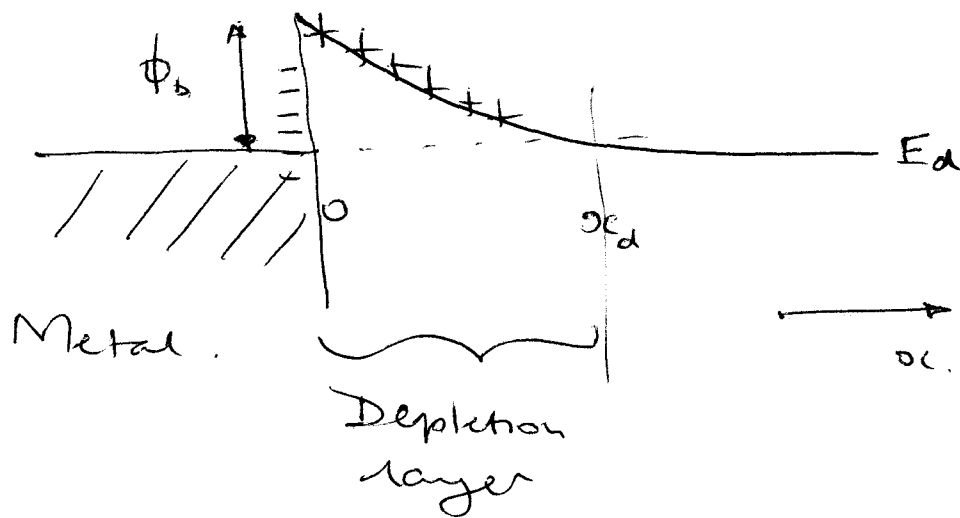
Significant overlap when

$$n(a_B^*)^3 \sim 1$$

$$\text{i.e. } n \sim \left( \frac{1}{a_B^*} \right)^3 \sim \text{few} \times 10^{21} \text{ m}^{-3}$$

Schottky barrier

See figure 13.1 (c)



In the depletion layer, charge density is  $N_d e$  where  $N_d$  is donor density.

Poisson  $\Rightarrow \nabla \cdot D = \frac{N_d e}{\epsilon_0}$

So the electrostatic potential ~~is given by~~

$$\frac{d^2 \bar{\Phi}}{dx^2} = - \frac{N_d e}{\epsilon \epsilon_0}$$

Solution is

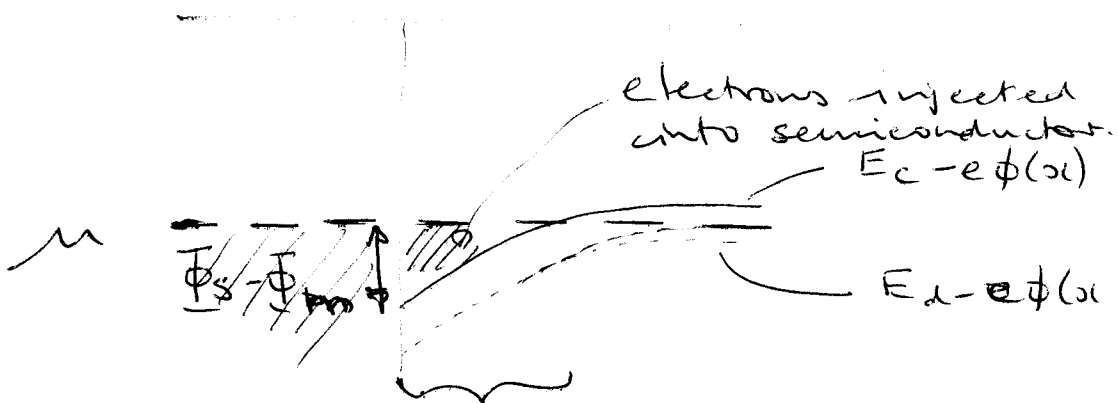
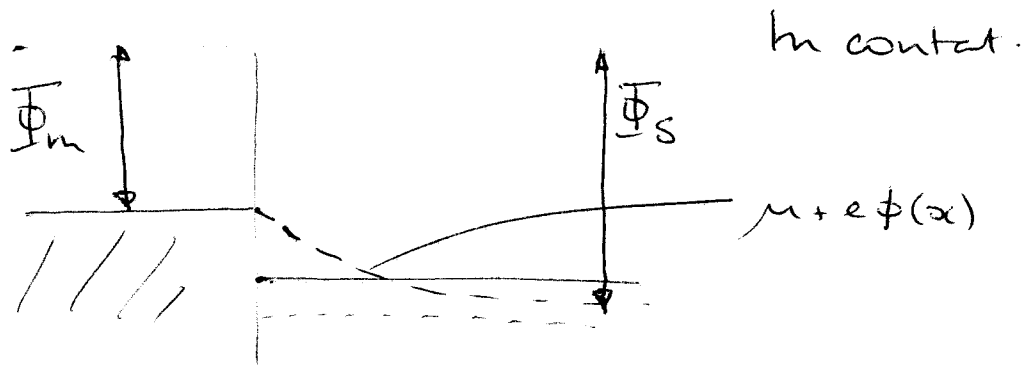
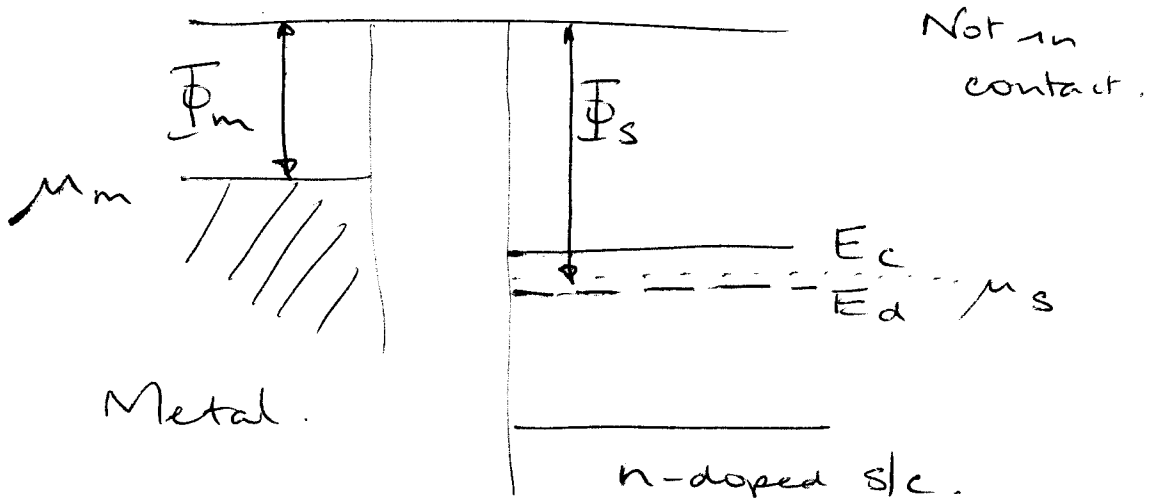
$$\bar{\Phi} = \bar{\Phi}_b - \frac{N_d e x^2}{2 \epsilon \epsilon_0}$$

→ Width of depletion layer

$$x_d = \left( \frac{2 \epsilon \epsilon_0 \bar{\Phi}_b}{N_d e} \right)^{1/2}$$

$$N_d = 10^{22} \text{ m}^{-3}, \bar{\Phi}_b = 0.5 \text{ V}, \epsilon = 12$$

$$\Rightarrow \boxed{x_d \sim 0.2 \mu\text{m}}$$



this length scale is short,  
because the n-doped s/c is  
metallized & screens  $\phi$  well

— no barrier

— good metallic contact

— "ohmic" contact