Wednesday 18 January 2006 10.30am to 12.30pm

THEORETICAL PHYSICS I

Answer **three** questions only. The approximate number of marks allotted to each part of a question is indicated in the right margin where appropriate. The paper contains **3** sides and is accompanied by a book giving values of constants and containing mathematical formulae which you may quote without proof.

1 Two equal masses m joined by a length 2a of light thin wire are in orbit about a planet of mass M. The wire is in the plane of the orbit at all times. Write down a Lagrangian for the system in terms of the polar coordinates (r, ϕ) of its midpoint and the angle θ between the wire and the line to the planet centre, assuming $a \ll$ distance from the planet and the gravitational attraction between the two little masses is negligible.



Find the equation of motion for θ .	[7]
Determine the stability of equilibrium values of θ .	[8]
If the orbit is circular and the system is close to a stable value of θ , how	
many oscillations does it perform in one orbital period?	[12]
2 Describe the terms in the Bernoulli equation in fluid dynamics, and the	
conditions when it applies.	[10]
Consider a bath of cross-sectional area $0.6 \mathrm{m}^2$, filled to a depth of $20 \mathrm{cm}$.	
Estimate how long it takes to empty when the plug is removed, if the area of the	
plug hole is 10 cm^2 .	[12]
Two cylindrical jets of water, which have the same radius <i>a</i> and velocity	
components $(0,0,v)$ and $(0,0,-v)$ respectively, meet head-on at the origin and	
spread to form a sheet in the $z = 0$ plane. Show the thickness of this sheet at	
distance <i>r</i> from the origin is a^2/r .	[12]

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[7]

3 The Lagrangian *L* of a particle of mass *m* and charge *e* moving with velocity *v* in an electrostatic potential ϕ is $\frac{1}{2}mv^2 - e\phi$. Using the requirement that $\int L dt$ should be Lorentz invariant, or otherwise, explain why the generalisation of the potential energy term in the Lagrangian L = T - V to $V = e(\phi - v \cdot A)$ is required, where (ϕ, A) is the electromagnetic 4-potential.

Consider the non-relativistic motion of the charged particle in a constant magnetic field \boldsymbol{B} which is directed along *z*-axis. In cylindrical polar coordinates (r, θ, z) :

(a) By using the Stokes theorem for $\int \mathbf{A} \cdot d\mathbf{l}$, or otherwise, show that $\{A_{\theta} = \frac{1}{2}Br, A_r = 0, A_z = 0\}$ represents a magnetic field with a *z*-component *B*. [8]

(b) Write down the Lagrangian and derive the equations of motion: in particular, show that [8]

$$\dot{\theta} = -\frac{eB}{2m} - \frac{J}{r^2}$$

where J is a constant.

(c) Show that the radius of the helical orbit must be proportional to $B^{-1/2}$ and the angular frequency of particle on this orbit is $|\dot{\theta}| = eB/m$. [8]

(d) Further show that the helical pitch angle ψ (the angle between B and v) obeys $\tan \psi \propto B^{1/2}$. [4]

4 Describe how contour integration methods can be used to evaluate definite integrals of functions with no poles. Illustrate your answer by showing that:

$$\int_{0}^{\infty} \mathrm{d}x \cos(x^{2}) = \int_{0}^{\infty} \mathrm{d}x \sin(x^{2}) = \frac{\sqrt{2\pi}}{4}$$
[10]

(Hint: use a wedge-shaped contour with angle $\pi/4$, not forgetting that $e^{i\pi/4} = (1+i)/\sqrt{2}$.)

How are poles on the real axis treated in contour integration? Illustrate your answer by showing that:

$$\int_{-\infty}^{\infty} \mathrm{d}x \, \frac{\sin(x)}{x} = \pi \tag{12}$$

[6]

(Hint: $\sin z = (e^{iz} - e^{-iz})/2i$)

How are branch cuts dealt with in contour integration? Illustrate your answer by showing that:

$$\int_{0}^{\infty} \mathrm{d}x \, \frac{x^{\alpha}}{1 + \sqrt{2}x + x^{2}} = \sqrt{2}\pi \frac{\sin(\alpha \pi/4)}{\sin(\alpha \pi)} \,. \tag{12}$$

for $-1 < \alpha < 1$.

5 Derive the Kramers-Kronig relations between the real and imaginary parts of the generalized susceptibility for a perturbation potential of the form V = -x.f where x(t) is a position coordinate and f(t) a force. [8] The equation of motion for a damped harmonic oscillator has the form:

$$\ddot{x} + \gamma \dot{x} + \omega^2 x = f(t)$$

Derive the relationship between the Green's function for this differential equation G(t - t'), the position coordinate x and the force f.

Derive an expression for the Fourier Transform of the Green's function $G(\omega)$ and write down the Kramers-Kronig relations for its real and imaginary parts.

Convert the principal-value integrals to contour integrals and find their poles. [7]

By evaluating one of these integrals show that the corresponding Kramers-Kronig relation is obeyed by this Green's function. [8] (Hint: $P \int_{-\infty}^{\infty} f(x) dx = \lim_{\epsilon \to 0} \left[\int_{-\infty}^{-\epsilon} + \int_{\epsilon}^{\infty} f(x) dx \right] f(x) dx$ for a pole at x = 0.)

6 Thermal (Johnson) noise in a resistor is modelled as a white-noise voltage source, with a mean-square amplitude $\langle V^2 \rangle = 2Rk_BT$, connected in series with an ideal resistor R. Consider a circuit consisting of such a resistor connected across a capacitor C. Show that the equation of motion for the charge on the capacitor obeys a Langevin equation.

Describe the standard form for the Langevin (stochastic) equation,

$$\dot{q}_{\alpha} = F_{\alpha}(\boldsymbol{q}) + G_{\alpha}^{k}(\boldsymbol{q})A_{k}(t)$$

and its relation to classical Brownian motion.

Outline principles leading to the Fokker-Planck (kinetic) equation for the probability density f(q, t): [8]

$$\frac{\partial f(\boldsymbol{q},t)}{\partial t} = \left\{ -\frac{\partial K_{\alpha}(\boldsymbol{q})}{\partial q_{\alpha}} + \frac{1}{2} \frac{\partial^2}{\partial q_{\alpha} \partial q_{\beta}} Q_{\alpha\beta}(\boldsymbol{q}) \right\} f(\boldsymbol{q},t) ,$$

with $K_{\alpha} = F_{\alpha} + \frac{1}{2} \frac{\partial G_{\alpha}^i}{\partial q_{\beta}} G_{\beta}^k \delta_{ik}; \quad Q_{\alpha\beta} = G_{\alpha}^i G_{\beta}^k \delta_{ik} .$

Obtain the Fokker-Planck equation for the probability density f(q, t) of the charge on the capacitor.

Show that the equilibrium probability density f(q) of the charge on the capacitor is proportional to $\exp[-q^2/2Ck_BT]$ and comment on its form. [10]

[6]

[6]

[4]

[4]

[7]