Stable Skyrmions in Two-Component Bose-Einstein Condensates

Nigel Cooper T.C.M. Group, Cavendish Laboratory, University of Cambridge

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Richard Battye (Cambridge/Manchester), Paul Sutcliffe (Kent) [PRL 88, 080401 (2002)]

Overview

- Atomic Bose-Einstein Condensates
- Multi-component condensates
- Topological Solitons
- Stable "Skyrmions"
- Summary

Bose-Einstein Condensation



$$\frac{\hbar^2}{m\lambda_T^2} \sim k_B T \qquad \bar{a} \sim \bar{n}^{-1/3}$$
$$\lambda_T \gtrsim \bar{a} \quad \Rightarrow \quad k_B T \lesssim \frac{\hbar^2}{m} \bar{n}^{2/3}$$

Bosons, $T < T_c \Rightarrow$ B.E.C.

$$T = 0 \quad \rightarrow \quad \psi_N = \prod_{i=1}^N \psi(\vec{r_i})$$

Atomic Bose-Einstein Condensates

 $k_B T_{\rm c} \simeq \frac{\hbar^2}{m} \bar{n}^{2/3} \sim 100 {\rm nK}$

⁸⁷Rb (JILA, 1995)
⁷Li (Rice, 1995)
²³Na (MIT, 1995)
¹H (MIT, 1998)
⁸⁵Rb (JILA, 2000)
⁴He* (Orsay, 2001)



[Anderson et. al. [JILA], Science 269, 198 (1995).]

What's special about Atomic BECs?

Fantastic experimental control



[Abo-Shaeer et al.[MIT], Science 476, 476 (2001)]

Tunable effective interactions

s-wave scattering length, a.

atom	a
⁸⁷ Rb	5.77nm
⁷ Li	-1.45nm
⁸⁵ Rb	$-? \rightarrow \sim 500$ nm

Novel uncondensed states

High vortex density [Wilkin & Gunn,...]

Superfluid/(Mott) Insulator transition

[Greiner et al.[Munich], Nature 415, 39 (2002)]

Multi-component systems

Simultaneous trapping and cooling of atoms of different species.

Multicomponent Condensates

Hyperfine interaction $\Rightarrow \vec{F} = \vec{I} + \vec{J}$

Typically:
$$egin{array}{lll} J=S=1/2 \ ({
m Alkali\ gases}) \ I=3/2 \ ({}^{87}{
m Rb},\ {}^{23}{
m Na},\ {}^{7}{
m Li}) \ |F=2,m=-2,-1,0,1,2
angle \ |F=1,m=-1,0,1
angle \ |F=1,m=-1,0,1\endel{}$$

Magnetic trap

 $^{87} {\sf Rb}~|F=2,m=2\rangle$ and $|F=1,m=-1\rangle$ [Myatt $\it et.~al.[{\sf JILA}],~{\sf PRL}{\bf 78},~{\sf 586}~(1997)]$

 $^{87}\mathsf{Rb}\;|F=2,m=1\rangle$ and $|F=1,m=-1\rangle$ [Hall $\mathit{et.}\;\mathit{al.[JILA]},\,\mathsf{PRL}\textbf{81},\,\mathsf{1539}\;(\mathsf{1998})]$

Optical trap

 $^{23}\mathsf{Na}\ |F=1,m=-1,0,1\rangle$ [Stenger $\mathit{et.}\ \mathit{al.}[\mathsf{MIT}]$, Nature **396**, 345 (1998)]

Different atoms / Isotopes ⁸⁵Rb and ⁸⁷Rb [Bloch *et. al.*[Munich], PRA **64**, 021402 (2001)]

Vortex in a two-component BEC

[Matthews et. al. [JILA], PRL 83, 2498 (1998).]



Gross-Pitaevskii Mean-Field Theory

$$egin{array}{rl} \Psi_N & \propto & \prod_{i=1}^N \psi(ec{r_i}) \ N & = & \int d^3ec{r} \, |\psi(ec{r})|^2 \end{array}$$

Minimise the expectation value of the energy

$$E = \int d^{3}\vec{r} \left[\frac{\hbar^{2}}{2m} |\nabla\psi|^{2} + V(\vec{r})|\psi|^{2} + \frac{1}{2}U|\psi|^{4} \right]$$

at fixed N (chemical potential) $U \propto a$

Multi-component case

$$\Psi_N \propto \prod_{i=1}^N \left[\sum_{\alpha} \psi_{\alpha}(\vec{r_i}) |\alpha_i\rangle \right]$$
$$N_{\alpha} = \int d^3 \vec{r} |\psi_{\alpha}(\vec{r})|^2$$

Energy (density)

$$\sum_{\alpha} \frac{\hbar^2}{2m} |\nabla \psi_{\alpha}|^2 + V_{\alpha}(\vec{r}) |\psi_{\alpha}|^2 + \frac{1}{2} \sum_{\alpha,\beta} U_{\alpha\beta} |\psi_{\alpha}|^2 |\psi_{\beta}|^2$$

- $U_{\alpha\beta} \Rightarrow$ all mutual two-body scattering lengths
- N_{α} conserved \Rightarrow separate chemical potentials

•
$$\psi = \begin{pmatrix} \psi_1 \\ \psi_2 \\ \vdots \end{pmatrix}$$
 \Rightarrow topological solitons ?

Topological Textures

One-component condensate in one dimension.



"Topological invariant"

$$Q = \frac{1}{2\pi} \int_0^L \frac{d\theta}{dx} dx$$

Two-component condensate in three dimensions.

$$\begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix} = \sqrt{\rho} \begin{pmatrix} \cos(\theta/2)e^{i\phi_1} \\ \sin(\theta/2)e^{i\phi_2} \end{pmatrix}$$

$$Q = \frac{1}{8\pi^2} \int \sin\theta \nabla_i \theta \nabla_j \phi_1 \nabla_k \phi_2 \epsilon_{ijk} d^3 \vec{r}$$



[Al Kawaja and Stoof, Nature **411**, 918 (2001); Ruostekoski and Anglin, PRL **86**, 3934 (2001)]

[Three components, textures vortices, monopoles...]

How to obtain stable Skyrmions

- Large Trap $\psi(|\vec{r}| \to \infty) = \sqrt{\rho_0} \begin{pmatrix} 1 \\ 0 \end{pmatrix}$
- Constrain $N_2 = \int d^3 \vec{r} \, |\psi_2|^2$
- Regime of phase separation: $U_{12}^2 > U_{11}U_{22}$

We study $U_{11} \sim U_{12} \sim U_{22}$ s.t. $\rho(\vec{r}) = \rho_0$.

Find stable skyrmions of the form:



"Imprint" with lasers [Ruostekoski and Anglin, PRL 86, 3934 (2001)]

[cf. "Cosmic vortons"]

Mathematical Details

$$\begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix} = \sqrt{\rho_0} \begin{pmatrix} \cos(\theta/2)e^{i\phi_1} \\ \sin(\theta/2)e^{i\phi_2} \end{pmatrix}$$

$$N_{2} = \rho_{0} \int \sin^{2}(\theta/2) d^{3}\vec{r}$$
$$Q = \frac{1}{8\pi^{2}} \int \sin\theta \nabla_{i}\theta \nabla_{j}\phi_{1}\nabla_{k}\phi_{2} \epsilon_{ijk} d^{3}\vec{r}$$

Energy density

$$\frac{\hbar^2 \rho_0}{2m} \left[\frac{1}{4} |\nabla \theta|^2 + \cos^2(\theta/2) |\nabla \phi_1|^2 + \sin^2(\theta/2) |\nabla \phi_2|^2 \right] \\ + \Delta \sin^2 \theta$$

 $\left[\Delta \equiv \frac{1}{8}\rho_0^2 \left(2U_{12} - U_{11} - U_{22}\right)\right]$ Lengthscales: $\xi_{\Delta} \equiv \sqrt{\frac{\hbar^2 \rho}{2m\Delta}} \quad R_2 = \left(\frac{N_2}{\rho_0}\right)^{1/3}$

$$E(\Delta, N_2) = \left(\frac{\hbar^2 \rho_0 R_2}{m}\right) \mathcal{E}_Q(\eta)$$
$$\overbrace{\eta \equiv \frac{R_2}{\xi_\Delta}}$$

Q=1







$$Q=2$$









Axisymmetric Ansatz

 $egin{aligned} & heta(r,z), \phi_1(r,z), \phi_2 = m\chi \ & [(r,\chi,z) \text{ are cylindrical polar co-ordinates}] \end{aligned}$



Moving Vortex Rings

Constrain the impulse (momentum)

$$P_i = \frac{\hbar}{2i} \int d^3 \vec{r} [r_j \nabla_i \psi^*_\alpha \nabla_j \psi_\alpha - r_j \nabla_j \psi^*_\alpha \nabla_i \psi_\alpha]$$

$$\vec{v} = \frac{\partial E}{\partial \vec{P}}$$



Summary

• Atomic Bose-Einstein condensates offer the possibility of studying interacting Bose gases with a high level of control (interaction strength, confinement, numbers of components).

• In the regime of *phase separation*, two-component BECs have stable textures with the topology of Skyrmions (Q = 1, 2).

• This regime is relevant for 2-component ⁸⁷Rb systems. We expect that textures imprinted by lasers will relax to these stable Skyrmion configurations.