

# Random Estimates in QMC

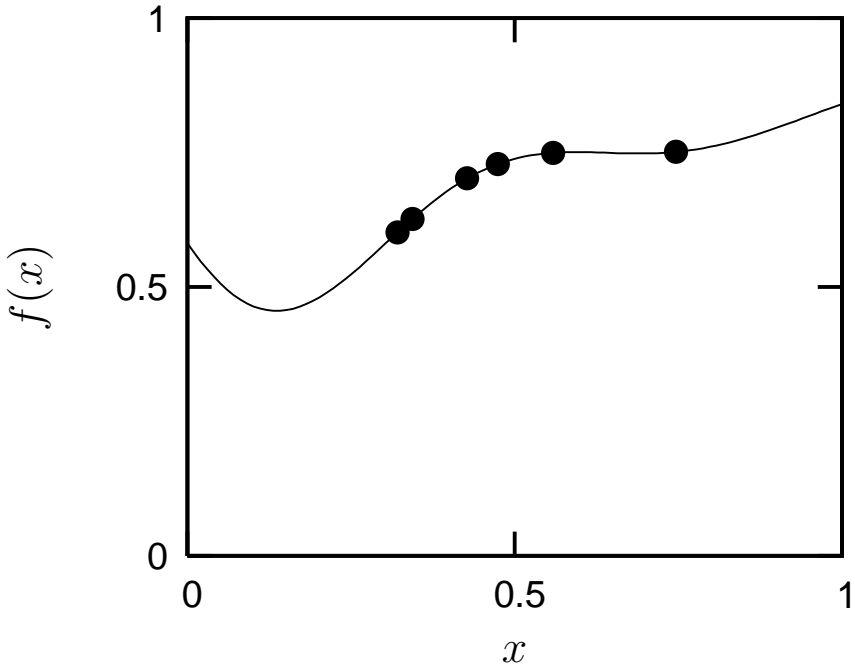
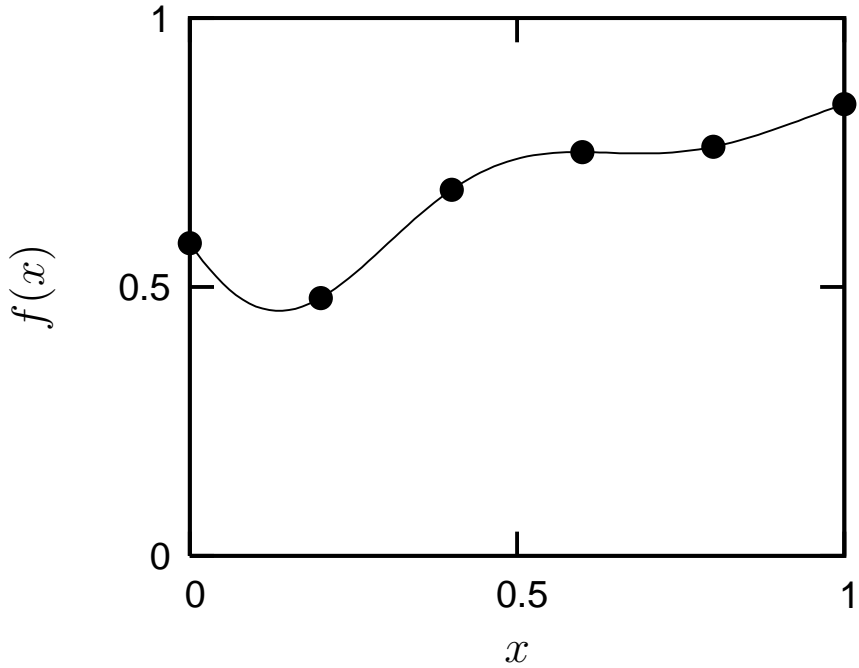
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*ESDG February 2007*

### Numerical Integration

Unkown function,  $r$  sample points



Evenly sampled grid:

$$\int_0^1 f(x)dx \approx \sum a_n f(x_n) \quad , \quad \epsilon \propto \Delta^p \approx \left(\frac{1}{r}\right)^{p/D}$$

Points sampled randomly, with PDF  $P(x)$ :

$$\int_0^1 f(x)dx \approx \frac{1}{r} \sum f(x_n)/P(x_n) \quad , \quad \epsilon \propto \frac{1}{\sqrt{r}}$$

## Quantum Monte Carlo (VMC)

Sample  $3N$  dimensional space with PDF  $P(\mathbf{R})$

$$\text{Est} [E_{tot}] = \frac{\sum \psi^2 E_L / P}{\sum \psi^2 / P} = \frac{\langle \psi | \hat{H} | \psi \rangle + Y}{\langle \psi | \psi \rangle + X}$$

where  $E_L = \psi^{-1} \hat{H} \psi$ .

Simplest case is 'Standard Sampling': Choose  $P(\mathbf{R}) = A\psi(\mathbf{R})^2$ , then

$$\text{Est} [E_{tot}] = \frac{1}{r} \sum E_L = \frac{\langle \psi | \hat{H} | \psi \rangle}{\langle \psi | \psi \rangle} + W$$

- $W$  is the random error in a sum of random variables, so what is its distribution?
- **IF** the CLT is valid then it is Gaussian with mean 0, and variance  $\sigma/r^{1/2}$ .

**3N → 1 dimension**

Why?: Easier to deal with the general case analytically

A change of the random variable from spatial to energy:

$$\begin{aligned} E_{tot} &= \int_V \psi^2 E_L d^{3N} \mathbf{R} / \int_V \psi^2 d^{3N} \mathbf{R} \\ &= \int_{-\infty}^{\infty} P_{\psi^2}(E) E dE \end{aligned}$$

with

$$P_{\psi^2}(E) = \int_{E=E_L} \frac{P(\mathbf{R})}{|\nabla_{\mathbf{R}} E_L|} d^{3N-1} \mathbf{R}$$

- A histogram of  $E_L$  approximates the 'seed' PDF  $P_{\psi^2}$
- $|\nabla_{\mathbf{R}} E_L|$  results from curvilinear co-ordinates and change of variables.
- Useless numerically, but useful analytically.

### What can we say about $P_{\psi^2}$ ?

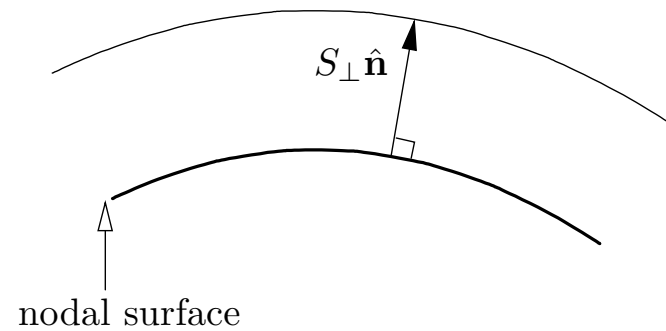
Singularities in the local energy:

$$E_L(\mathbf{R}) = -\frac{1}{2} \frac{\nabla_{\mathbf{R}}^2 \psi}{\psi} + \sum_{i < j} \frac{1}{r_{ij}} - \sum_i \frac{Z}{r_i}$$

- $E_L(\mathbf{R}) = E_{tot}$  if the trial wavefunction,  $\psi$ , is exact
- Enforce Kato cusp conditions  $\rightarrow$  no Coulomb singularities
- Nodal surface is  $\psi = 0$ , and is  $3N - 1$  dimensional
- Kinetic energy part gives singularity on a  $3N - 1$  dimensional surface
- Singularities provide information about  $P_{\psi^2}$  for large  $|E|$

What can we say about  $P_{\psi^2}$ ?

$$P_{\psi^2}(E) = \int_{E=E_L} \frac{P(\mathbf{R})}{|\nabla_{\mathbf{R}} E_L|} d^{3N-1} \mathbf{R}$$



$$\psi = a_1 S_{\perp} + \dots$$

$$E_L = b_{-1} S_{\perp}^{-1} + \dots$$

$$P(\mathbf{R})/|\nabla E_L| = c_4 S_{\perp}^4 + \dots$$

$$P_{\psi^2}(E) = d_{-4} E^{-4} + \dots$$

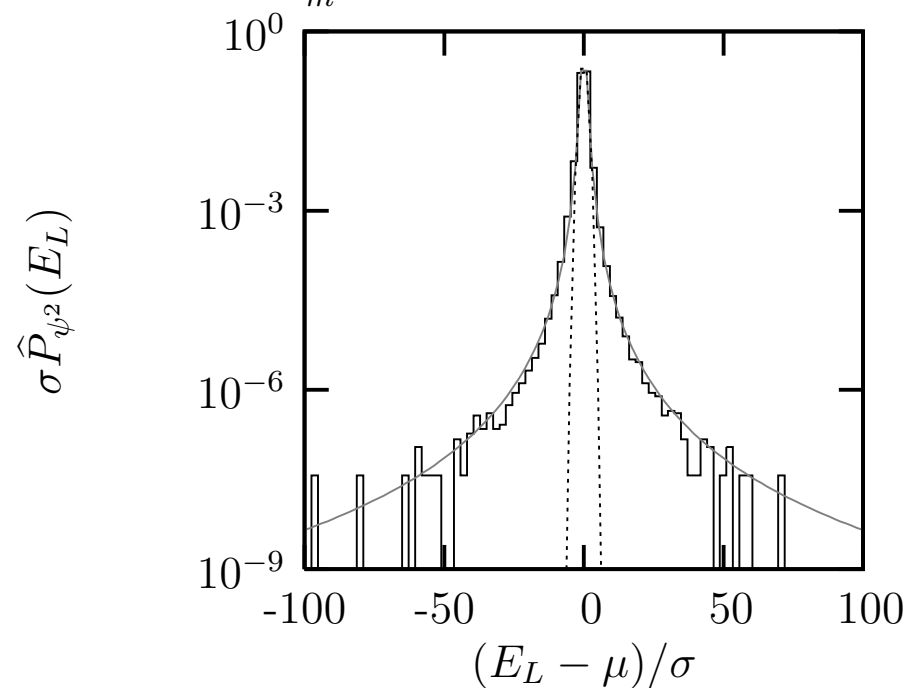
or more completely

$$P_{\psi^2}(E) = (E - E_0)^{-4} \left( e_0 + \frac{e_1}{(E - E_0)} + \dots \right) \quad |E| \gg E_0$$

### Example: All-electron isolated Carbon atom

- Jastrow + 48 determinants + backflow:

$$\psi = e^{J(\mathbf{R})} \sum_m a_m D_m^\uparrow(\mathbf{R}') D_m^\downarrow(\mathbf{R}') \quad \text{with} \quad \mathbf{R}' = \mathbf{R}'(\mathbf{R})$$



Estimated seed probability density function

- 93% correlation energy at VMC level

Also shown is  $\frac{\sqrt{2}}{\pi} \frac{\sigma^3}{\sigma^4 + (E - \mu)^4}$ , and a Normal distribution

### Random error in total energy estimate

$$\text{Est} [E_{tot}] = \frac{1}{r} (E_1 + \dots + E_r)$$

Product of probability of  $r$  samples energies that add up to  $rE_{tot} \rightarrow$  convolution integrals

$$P_{r=2}(2E_{tot}) = \int P_{\psi^2}(E_1)P_{\psi^2}(E_2)\delta(E_1 + E_2 - 2E_{tot})dE_1dE_2 = P_{\psi^2} \star P_{\psi^2}$$

$$P_r(rE_{tot}) = P_{\psi^2} \star P_{\psi^2} \star \dots \star P_{\psi^2}$$

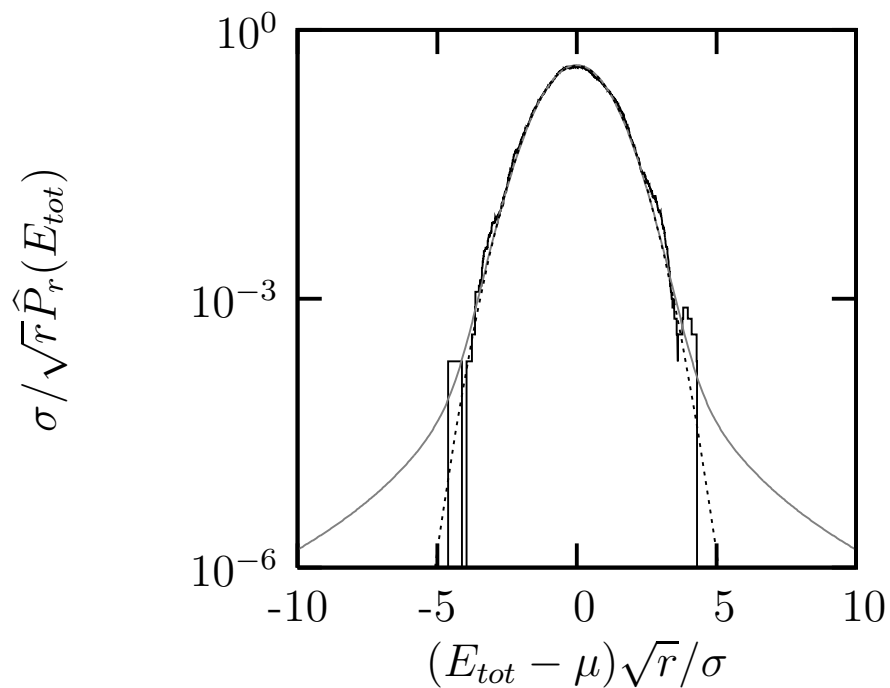
- Take Fourier transform of  $P_{\psi^2}(E)$
- Take the  $r^{th}$  power
- Take the inverse Fourier transform
- Rescale some variables to get the PDF of averages instead of sum

$$P_r(y) = \frac{1}{\sqrt{2\pi}} \left[ 1 + \frac{\eta}{\sqrt{r}} \frac{d^3}{dy^3} + \mathcal{O}\left(\frac{1}{r}\right) \right] e^{-y^2/2} + \left[ \frac{\lambda}{3\pi} \frac{1}{\sqrt{r}} \frac{d^3}{dy^3} D\left(\frac{y}{\sqrt{2}}\right) + \mathcal{O}\left(\frac{1}{r}\right) \right]$$

$$y = (E_{tot} - \mu)/\sigma$$



## PDF of estimate of Total energy



- Approximate PDF from  $10^4$  estimates of total energy, with  $r = 10^3$
- For small  $|E|$ , PDF is dominated by  $\frac{1}{\sqrt{2\pi}} e^{-x^2/2}$
- For large  $|E|$ , PDF is dominated by  $\frac{\sqrt{2}}{\pi} \frac{\lambda}{\sqrt{r}} 1/x^4$  ( $\lambda \approx 1$  for Carbon trial wavefunction)
- CLT is true in its weakest form

### PDF of estimate of the 'Residual Variance', $v$

Optimisation of wavefunctions using the 'residual variance',  $v$

$$\left(\hat{H} - E_{tot}\right) \psi = \delta = (E_L - E_{tot}) \psi$$

$$v = \int \delta^2 d\mathbf{R} \geq 0, \text{ and zero for exact } \psi$$

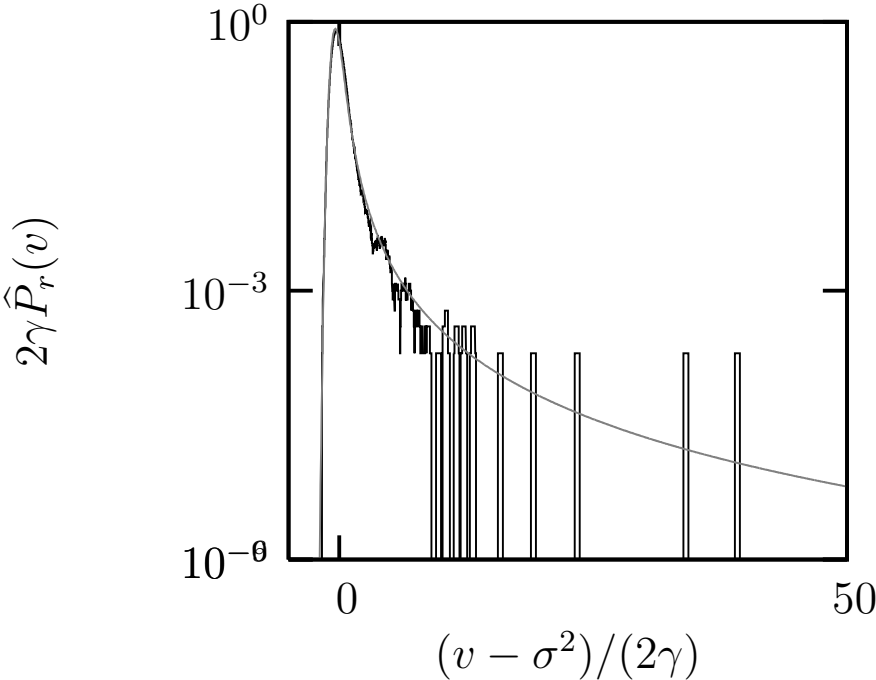
- To optimise the wavefunction  $v$  is usually minimised
- Analyse effect of tails, as before:

$$P_r(\bar{v}) = \frac{\sqrt{3}}{\pi} \frac{1}{2\gamma} \left[\frac{\bar{v} - \sigma^2}{2\gamma}\right]^2 \exp\left(\left[\frac{\bar{v} - \sigma^2}{2\gamma}\right]^3\right) \\ \times \left[ -\text{sgn}[\bar{v} - \sigma^2] K_{1/3}\left(\left|\frac{\bar{v} - \sigma^2}{2\gamma}\right|^3\right) + K_{2/3}\left(\left|\frac{\bar{v} - \sigma^2}{2\gamma}\right|^3\right) \right]$$

with the 'width' of the PDF decided by the magnitude of the tails

$$\gamma = \left[\frac{6\lambda^2}{\pi r}\right]^{1/3} \sigma^2 \quad (1)$$

PDF of estimate of the 'Residual Variance',  $v$



- Approximate PDF from  $10^4$  estimates of residual variance, with  $r = 10^3$
- Small  $v$ ,  $\propto e^{x^3}$ . Large  $v \propto 1/x^{5/2}$
- PDF has no variance,  $\gamma$  has no vigorous statistical estimate and is  $\propto r^{-1/3}$

- CLT is valid in its weakest form for the total energy
- CLT not valid for residual variance
- CLT is likely to be invalid for estimates of other physical quantities
- **Because:**  $\psi^2$  samples  $E_L$  rarely where it is largest, at the nodal surface

**Can the CLT be reinstated ?**

### Residual sampling

Instead of sampling with  $P = A\psi^2$ , sample with  $P = A\psi^2/w$ , then

$$\text{Est} [E_{tot}] = \frac{\sum w E_L}{\sum w} = \frac{\langle \psi | \hat{H} | \psi \rangle + \mathbf{Y}}{\langle \psi | \psi \rangle + \mathbf{X}}$$

and the residual variance,

$$\text{Est} \left[ \int \delta^2 d\mathbf{R} \right] = \frac{\sum w (E_L - E_{tot})^2}{\sum w} = \frac{\int \psi^2 (E_L - E_{tot})^2 d\mathbf{R} + \mathbf{Y}}{\int \psi^2 d\mathbf{R} + \mathbf{X}}$$

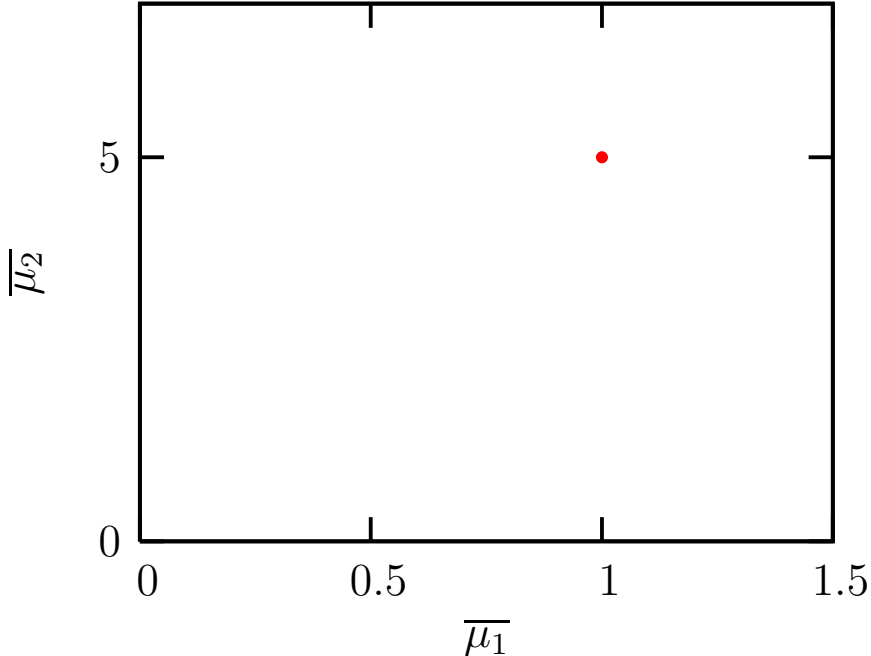
Choose the weighting function

$$w(E_L) = \frac{\epsilon^2}{(E_L - E_0)^2 + \epsilon^2}$$

to 'interpolate' between sampling the numerator and denominator perfectly.

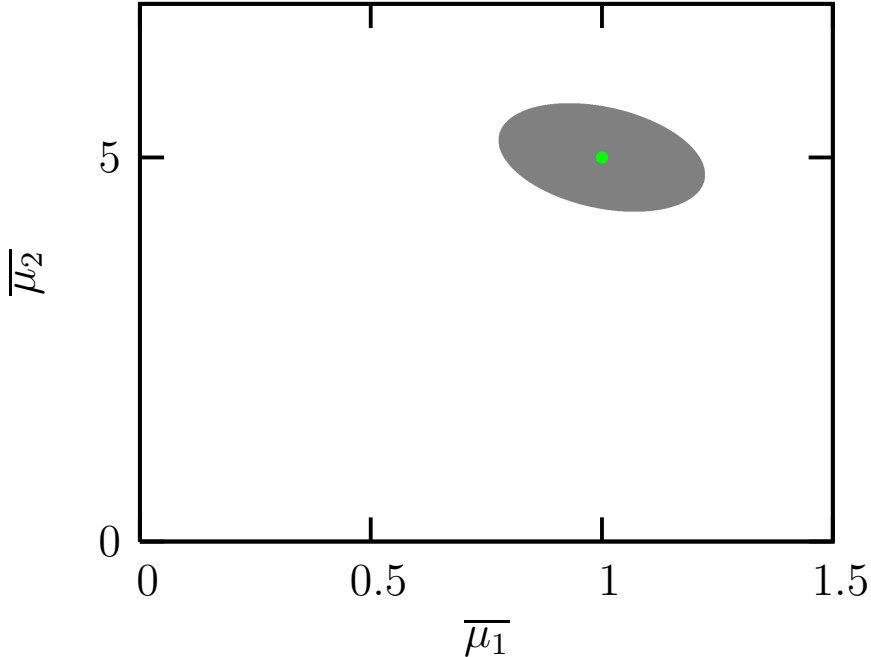
- No singularities, and no power law tails
- Quotient of two correlated random variables, each a sum of random variables

**Fieller's Theorem**



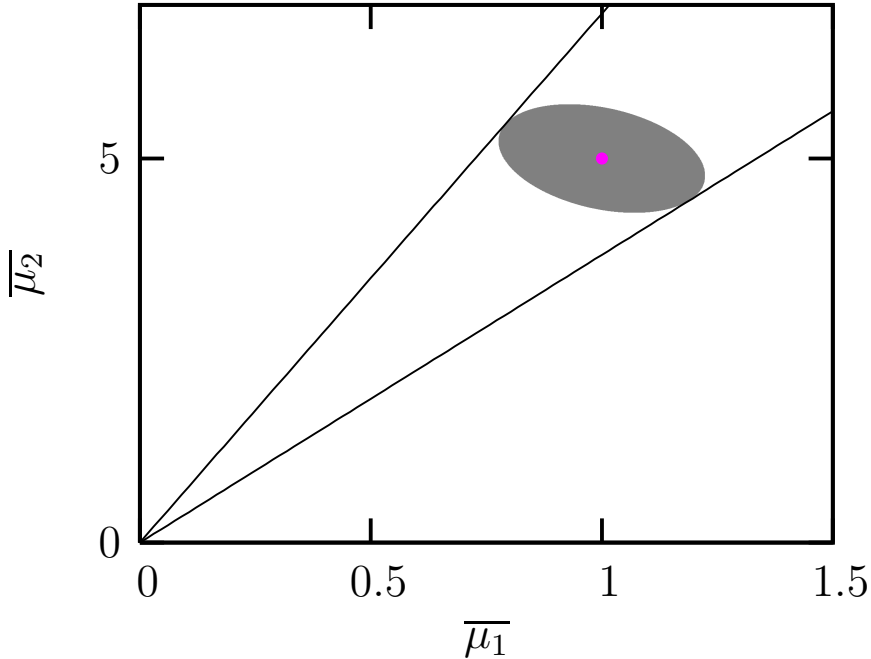
- $(\mu_2, \mu_1)$  that give  $\text{Est} = \mu_2/\mu_1$

### Fieller's Theorem



- $(\mu_2, \mu_1)$  that give  $\text{Est} = \mu_2/\mu_1$
- Ellipse containing 39% of probability from covariance matrix and bivariate CLT

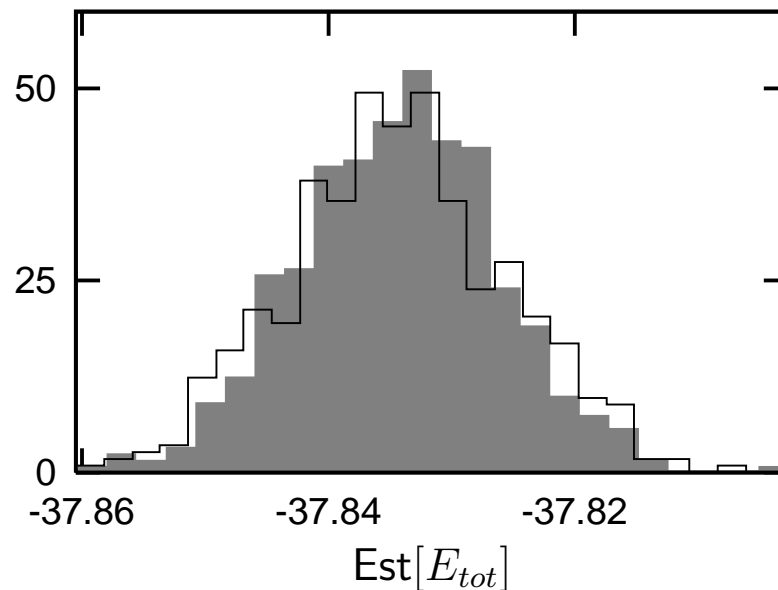
### Fieller's Theorem



- $(\mu_2, \mu_1)$  that give  $\text{Est} = \mu_2/\mu_1$
  - Ellipse containing 39% of probability from covariance matrix
  - Wedge that contains 68.3% of probability
- $\Rightarrow m_1 < \mu_2/\mu_1 < m_2$  with confidence 68.3%



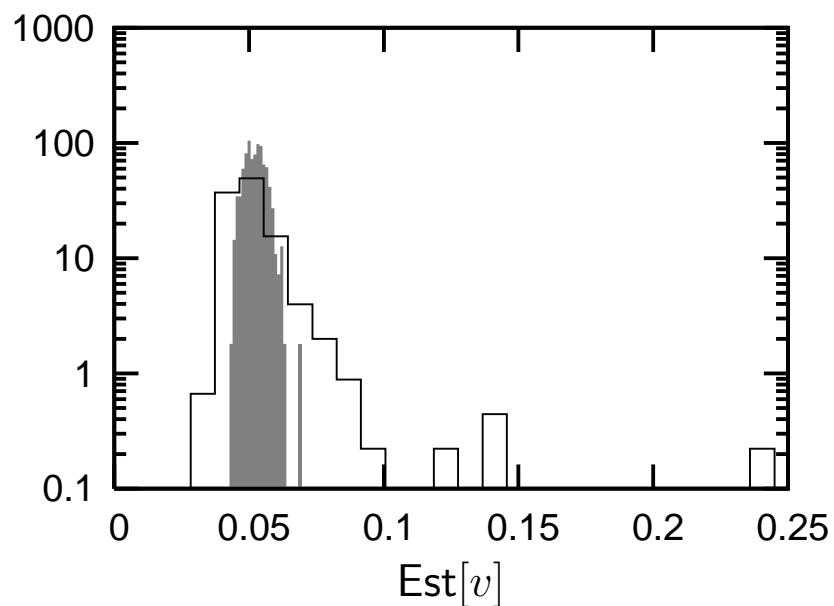
### Estimate of total energy



Histogram of  $10^3$  total energy estimates, each total energy estimate from  $10^3$  configurations.

- Residual sampling (filled) and standard sampling (unfilled) are not significantly different
- Residual sampling reduces error by  $\sim 30\%$
- For other systems standard sampling may give 'power law' outliers (depending on  $\lambda$ )
- For all systems residual sampling does not give 'power law' outliers

### Estimate of residual variance



Histogram of  $10^3$  residual variance estimates, each estimated from  $10^3$  configurations.

- Residual sampling and standard sampling are very different
- Standard sampling shows the  $v^{-5/2}$  tails and outliers expected
- Residual sampling gives well defined confidence limits for estimate via the bivariate CLT
- Standard sampling does not

## Optimisation

- 1) Take samples using wavefunction with parameter  $\alpha_0, \{\mathbf{R}\}_r$
- 2) These define random sample from a distribution of Optimisation functions,  $O(\alpha)$
- 3) Find the minimum of  $O(\alpha)$ , at  $\alpha = \alpha_{min}$
- 4) set  $\alpha_0 = \alpha_{min}$ , and return to 1)

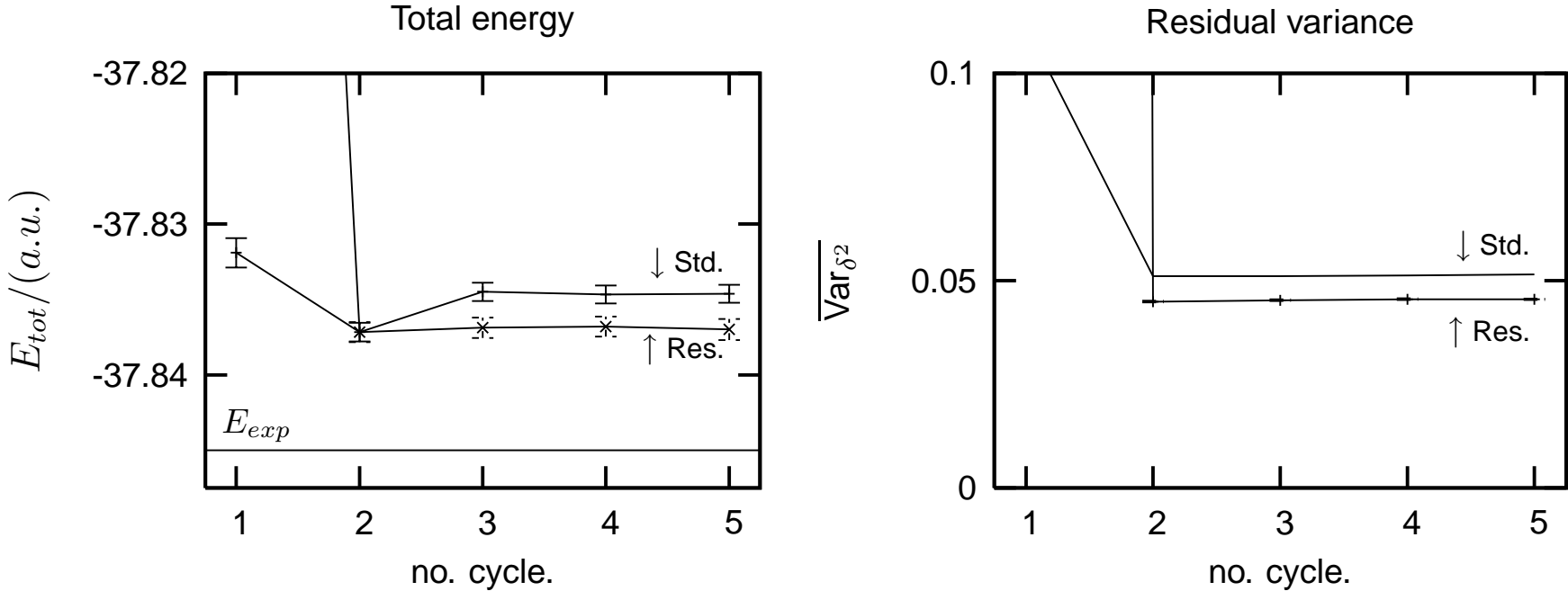
What is the distribution of  $O(\alpha)$ ?

$$O(\alpha) = \frac{\mathbf{a}_0 + \mathbf{a}_1(\alpha - \alpha_0) + \dots}{\mathbf{b}_0 + \mathbf{b}_1(\alpha - \alpha_0) + \dots}$$

- For each type of sampling/optimisation this expansion gives statistics of random error
- For 'Standard Sampling' error is not normal unless the nodes are suppressed by introducing a weight function into  $O(\alpha)$  solely for this purpose (eg  $\int g(E_L) \delta^2 d^{3N} \mathbf{R}$ ).
- For Residual sampling the error is always normal

### Optimisation

$r = 10^5$  configurations.



Std. - standard sampling with nodal surface 'suppressed' (93%  $E_{corr}$ )

Res. - residual sampling of residual variance (95%  $E_{corr}$ )

- New sampling provides lower total energy and a lower residual variance than standard sampling with nodal surface 'suppressed'

## Conclusions

- We cannot assume the CLT is true for estimates in ‘standard sampling QMC’
- ‘ $r$  large’ enough must be shown to be true for each estimate in ‘standard sampling QMC’
- The CLT can be reinstated by using an alternative sampling strategy
- Random functions whose minimum gives ‘optimum’ wavefunctions are not generally normally distributed. Some do not converge as  $r \rightarrow \infty$
- The residual sampling strategy can guarantee that the CLT is valid for estimates and optimisation functions, as long as they exist
- With residual sampling optimisation functions can be chosen on physical grounds - to give a good wavefunction at the nodal surface and small *fixed node error* in DMC

## Acknowledgements

*Financial support was provided by the Engineering and Physical Sciences Research Council (EPSRC), UK, and computational resources were provided by the Cambridge-Cranfield High Performance Computing Facility.*