



Pulay Nodal Terms in Accurate Diffusion Monte Carlo Forces

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Motivation



Why forces in Diffusion Monte Carlo?

- Equilibrium geometries
- Energy derivatives are very useful!
- Molecular dynamics

Two main problems of forces in DMC are

- Infinite variance of force estimator
(addressed e.g. with pseudopotentials)¹
- Discontinuous 1st (higher) derivative
in DMC wavefunction at nodal surface

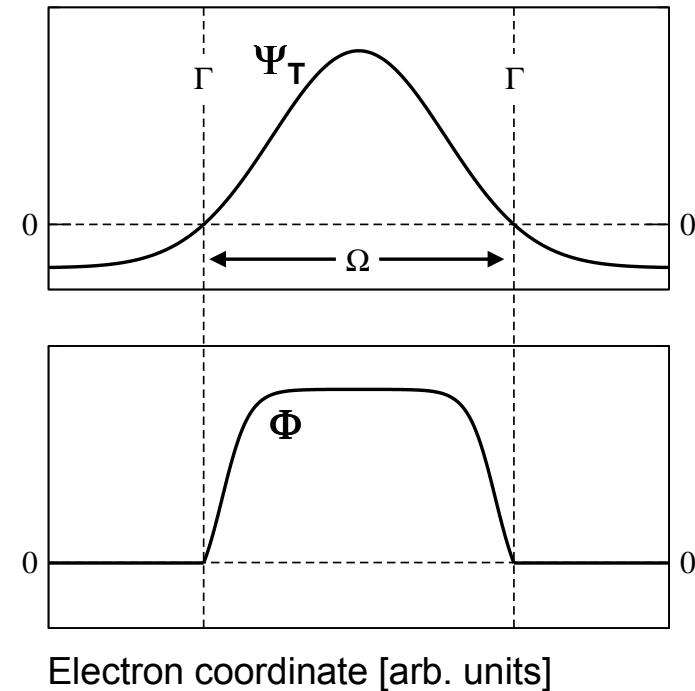
¹ESDG April 2007; A. Badinski, R.J. Needs PRE **76** 036707 (2007)

Diffusion Monte Carlo



Basics of DMC

- **Project out ground state**
 $|\Phi\rangle = e^{-\tau\hat{H}}|\Psi_T(\tau=0)\rangle$ for $\tau \rightarrow \infty$, $\tau = it$
using stochastic algorithm
- Φ is DMC wavefunction represented by ensemble of walkers
- Use **fixed node approx.** to eliminate fermionic sign problem
- Ψ_T is a given trial wavefunction
 Γ is nodal surface defined by $\Psi_T = 0$
Simulate nodal pockets individually



Diffusion Monte Carlo

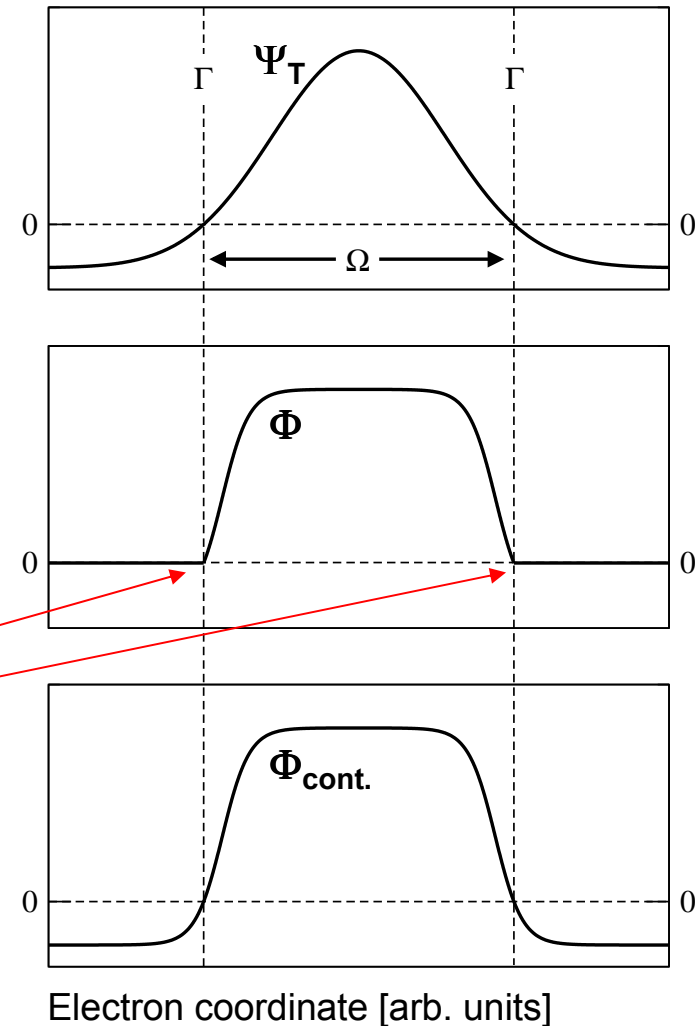


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Problem

Φ has discontinuous derivatives at Γ
Define $\Phi_{\text{cont.}}$ so it has no discontinuities



Energy in DMC



Effective Hamiltonian in DMC

After an involved derivation, we obtain

$$\hat{H}\Phi = \Theta(\Psi_T)\hat{H}\Phi_{cont.} - \frac{1}{2}\delta(\Psi_T)\frac{|\nabla\Psi_T|^2}{\Psi_T}\Phi_{cont.}$$

DMC energy

$$E_D = \frac{\int\Psi\hat{H}\Phi dV}{\int\Psi\Phi dV}$$

for **mixed DMC** ($\Psi=\Psi_T$) and **pure DMC** ($\Psi=\Phi$).

The δ function term (discontinuity in Φ) does not contribute to E_D .

But it may contribute when calculating derivatives of E_D

Forces in DMC



Differentiate E_D wrt nucleus coordinate λ

$$\frac{dE_D}{d\lambda} = \underbrace{\frac{\int \Psi \frac{d\hat{H}}{d\lambda} \Phi dV}{\int \Psi \Phi dV}} + \underbrace{\frac{\int \Psi (\hat{H} - E_D) \frac{d\Phi}{d\lambda} dV}{\int \Psi \Phi dV}} + \underbrace{\frac{\int \frac{d\Psi}{d\lambda} (\hat{H} - E_D) \Phi dV}{\int \Psi \Phi dV}}$$

$$\Psi = \Psi_T$$

Hellmann-Feynman force

use Reynolds' approx.¹

$$\frac{1}{\Phi} \frac{d\Phi}{d\lambda} \approx \frac{1}{\Psi_T} \frac{d\Psi_T}{d\lambda}$$

identify as nodal term
N(mixed)

$$\Psi = \Phi$$

Hellmann-Feynman force:
F(HFT, pure DMC)

nodal term²: **N(pure)**

¹ P. Reynolds, et al. Internat. J. Quant. Chem. **29** 589 (1986)

² F. Schautz and H.-J. Flad, J. Chem. Phys. **112**, 4421 (2000)

Nodal Term N



Volume integrals equal nodal term (steps omitted)

$$\mathbf{N(mixed)} = \frac{\int \frac{d\Psi_T}{d\lambda} (\hat{H} - E_D) \Phi dV}{\int \Psi_T \Phi dV} = -\frac{1}{2} \frac{\int_{\Gamma} \Psi_T \Phi \frac{|\nabla \Psi_T|}{\Psi_T} \frac{1}{\Psi_T} \frac{d\Psi_T}{d\lambda} dS}{\int \Psi_T \Phi dV}$$

$$\mathbf{N(pure)} = \text{Volume Terms} = -\frac{1}{2} \frac{\int_{\Gamma} \Phi \Phi \frac{|\nabla \Psi_T|}{\Psi_T} \frac{1}{\Psi_T} \frac{d\Psi_T}{d\lambda} dS}{\int_{\Gamma} \Phi \Phi dV}$$

1. these are exact expressions!
2. the averaged quantity only depends on Ψ_T
3. the averaged quantities are same in mixed and pure DMC

Using the extrapolation formula

$$\langle Q \rangle_{\text{pure}} \approx 2\langle Q \rangle_{\text{mixed}} - \langle Q \rangle_{\text{VMC}} \quad \text{with} \quad Q = \frac{|\nabla \Psi_T|}{\Psi_T} \frac{1}{\Psi_T} \frac{d\Psi_T}{d\lambda}$$

and $\langle Q \rangle_{\text{VMC}} = 0$ (proof omitted)

we find **N (pure) \approx 2 N (mixed)**

Summary



$$\text{mixed DMC } \frac{dE_D}{d\lambda} = \frac{\int \Psi_T \frac{d\hat{H}}{d\lambda} \Phi dV}{\int \Psi_T \Phi dV} + \frac{\int \Psi_T \Phi \left[\frac{1}{\Psi_T} \frac{d\Psi_T}{d\lambda} \frac{1}{\Psi_T} (\hat{H} - E_D) \Psi_T \right] dV}{\int \Psi_T \Phi dV} \\ + \frac{\int \Psi_T \Phi \frac{1}{\Psi_T} (\hat{H} - E_D) \frac{d\Psi_T}{d\lambda} dV}{\int \Psi_T \Phi dV} + O(\Delta\Psi_T)$$

$$\text{pure DMC } \frac{dE_D}{d\lambda} = \frac{\int \Phi \frac{d\hat{H}}{d\lambda} \Phi dV}{\int \Phi \Phi dV} + 2 \frac{\int \Psi_T \Phi \frac{1}{\Psi_T} (\hat{H} - E_D) \frac{d\Psi_T}{d\lambda} dV}{\int \Psi_T \Phi dV} + O(\Delta\Psi_T^2)$$

Computational Details



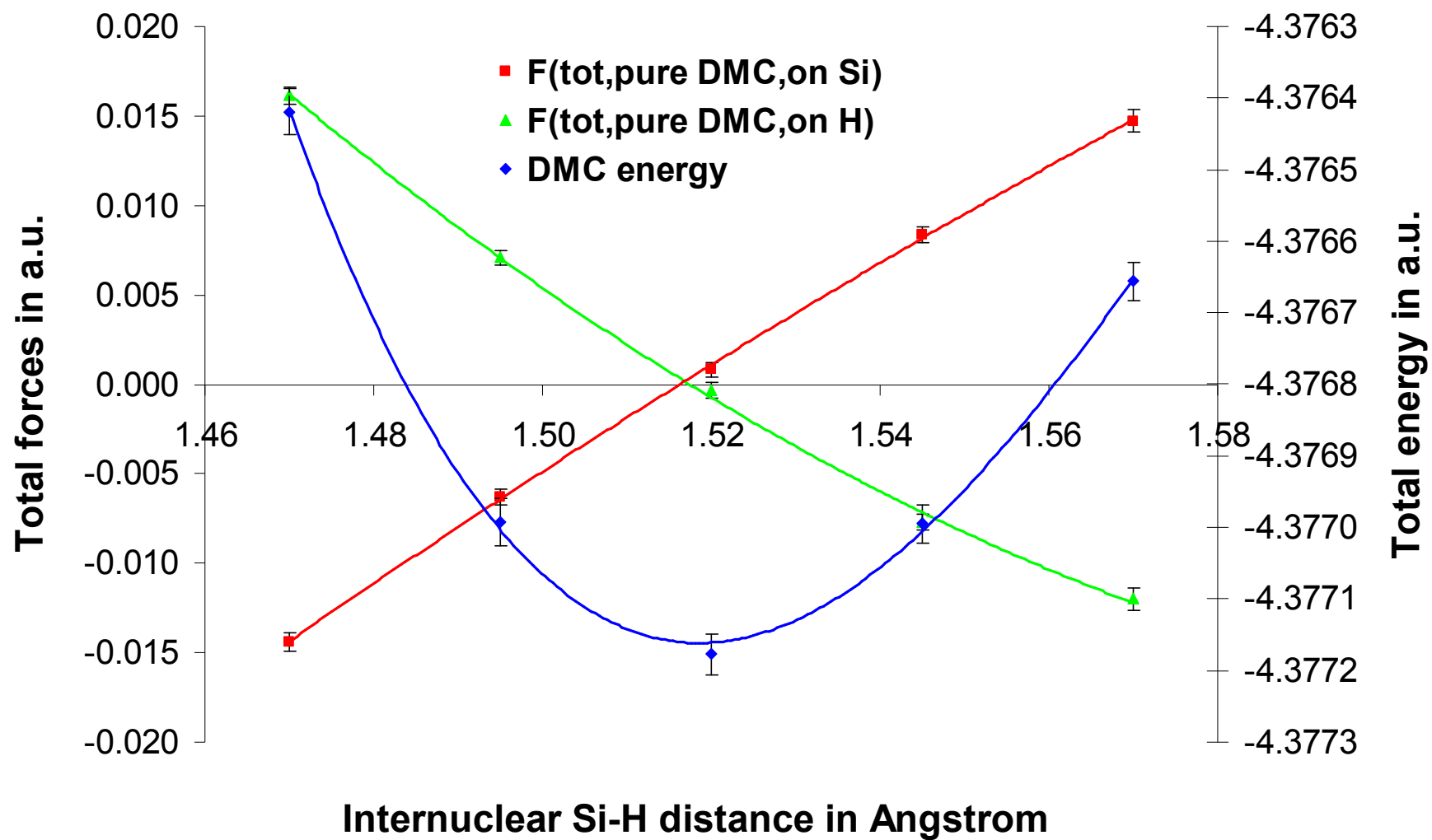
- GeH:
- **no electron-electron interaction** (nodal terms from kinetic energy!)
 - trial wavefunction: **single determinant** with 4 basis sets
 - **local pseudopotentials** (to avoid infinite variance!)

GeH, SiH, SiH₄:

- **full electron-electron interaction**
 - trial wavefunction: **single determinant** x Correlation function
 - **nonlocal pseudopotentials**¹
- calculate $\frac{\partial \Psi_T}{\partial \lambda}$ rather than $\frac{d\Psi_T}{d\lambda}$
 - use **future walking** method to calculate pure estimates
 - for reference also calculate energy gradient from potential energy curves

¹ A. Badinski, R.J. Needs PRE **76** 036707 (2007)

How to get geometries? (e.g. SiH)

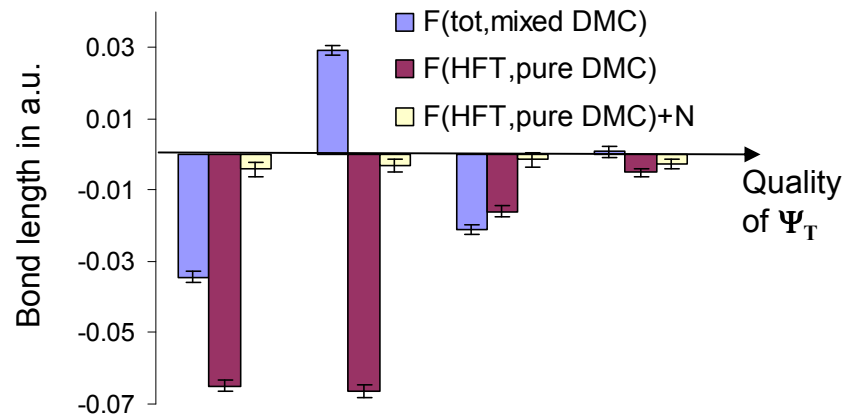


GeH (no e-e interaction)

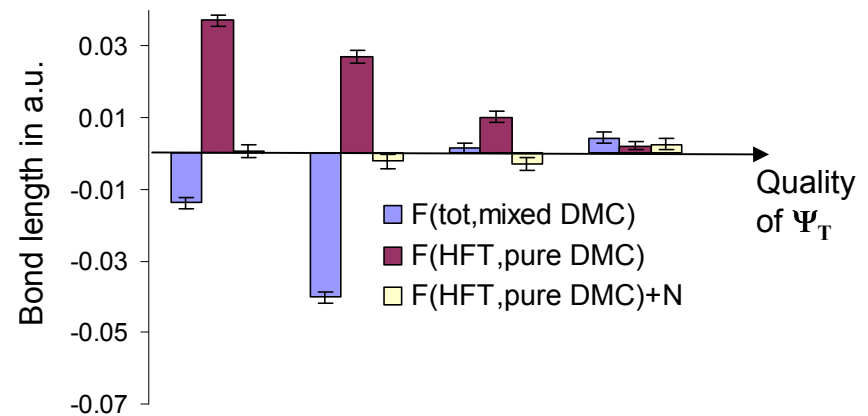


Difference between force & exact energy gradient (within basis set)

Forces on Ge



Forces on H



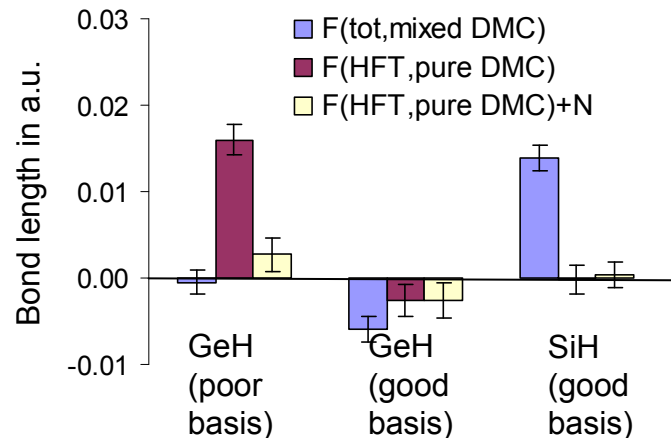
- F(tot,mixed DMC) is slightly better than F(HFT,pure DMC)
- Adding nodal term N to F(HFT,pure DMC) improves forces significantly
- F(tot,pure DMC) always better than F(tot,mixed DMC)

Nodal terms may be significant, should be included!

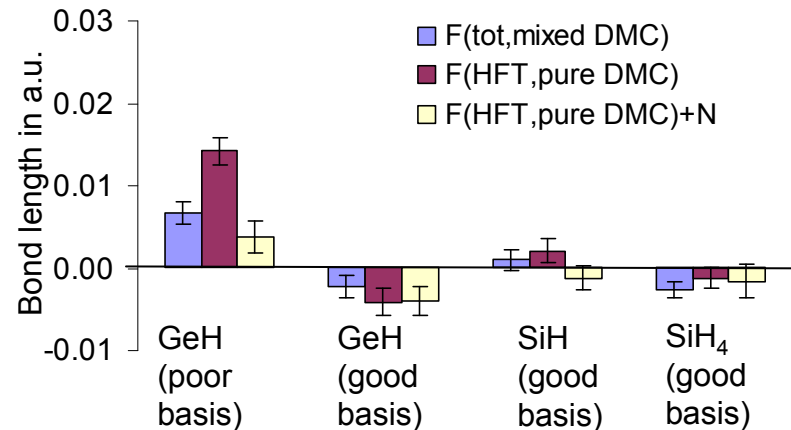
GeH, SiH, SiH₄ (with e-e interaction)

Difference between forces & exact energy gradient

Forces on Ge or Si



Forces on H



For poor basis:

- F(HFT,pure DMC) worse than F(tot,mixed DMC)
- Adding nodal term N to F(HFT,pure DMC) significantly improves forces

For good basis:

- F(HFT,pure DMC) equal or better than F(tot,mixed DMC)
- Adding nodal term N to F(HFT,pure DMC) has no significant effect

Nodal terms seem less important if basis set is good!

Conclusions



- We derived **exact expressions for forces** within mixed and pure DMC
- The nodal term in mixed DMC can be calculated straightforwardly
In pure DMC, it may be approximated as twice the mixed nodal term
- Tests for small molecules indicate that **nodal terms may be significant** and including them can significantly improve forces!
- Pure DMC forces including nodal terms seem more accurate than mixed DMC forces

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