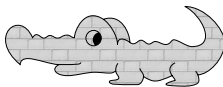


A general backflow transformation

Pablo López Ríos

TCM group. Cavendish Laboratory. University of Cambridge.

October 7, 2009



TCM

Quantum Monte Carlo

The Schrödinger equation for an electronic system

$$\left[\frac{1}{2} \sum_i \nabla_i^2 + \sum_{i < j} \frac{1}{|\mathbf{r}_i - \mathbf{r}_j|} - \sum_I \sum_i \frac{Z_I}{|\mathbf{r}_i - \mathbf{R}_I|} \right] \Phi_0(\mathbf{R}) = E_0 \Phi_0(\mathbf{R})$$

- Solvable exactly for very few systems.
- In general we don't know the exact $\Phi_0(\mathbf{R})$.
- We can give an approximate $\Psi(\mathbf{R})$ instead.

Quantum Monte Carlo

The Schrödinger equation for an electronic system

$$\left[\frac{1}{2} \sum_i \nabla_i^2 + \sum_{i < j} \frac{1}{|\mathbf{r}_i - \mathbf{r}_j|} - \sum_I \sum_i \frac{Z_I}{|\mathbf{r}_i - \mathbf{R}_I|} \right] \Phi_0(\mathbf{R}) = E_0 \Phi_0(\mathbf{R})$$

- Solvable exactly for very few systems.
- In general we don't know the exact $\Phi_0(\mathbf{R})$.
- We can give an approximate $\Psi(\mathbf{R})$ instead.

Variational Monte Carlo

Variational estimate of the energy

$$\begin{aligned} E_0 \leq E_V[\Psi] &= \frac{\int \Psi(\mathbf{R}) \hat{H}(\mathbf{R}) \Psi(\mathbf{R}) d\mathbf{R}}{\int \Psi^2(\mathbf{R}) d\mathbf{R}} \approx \\ &\approx \frac{1}{M} \sum_{m; \Psi^2} \frac{\hat{H}(\mathbf{R}_m) \Psi(\mathbf{R}_m)}{\Psi(\mathbf{R}_m)} = E_{\text{VMC}}[\Psi] \end{aligned}$$

The VMC method:

- is a very **simple** method
- along with the variational principle, allows for optimization of a parametrized $\Psi(\mathbf{R}; \alpha)$
- is entirely **dependent** on the quality of $\Psi(\mathbf{R})$

Variational Monte Carlo

Variational estimate of the energy

$$E_0 \leq E_V[\Psi] = \frac{\int \Psi(\mathbf{R}) \hat{H}(\mathbf{R}) \Psi(\mathbf{R}) d\mathbf{R}}{\int \Psi^2(\mathbf{R}) d\mathbf{R}} \approx$$
$$\approx \frac{1}{M} \sum_{m; \Psi^2} \frac{\hat{H}(\mathbf{R}_m) \Psi(\mathbf{R}_m)}{\Psi(\mathbf{R}_m)} = E_{\text{VMC}}[\Psi]$$

The VMC method:

- is a very **simple** method
- along with the variational principle, allows for **optimization** of a parametrized $\Psi(\mathbf{R}; \boldsymbol{\alpha})$
- is entirely **dependent** on the quality of $\Psi(\mathbf{R})$

Diffusion Monte Carlo

The time-dependent Schrödinger equation

$$\hat{H}\Phi(\mathbf{R},t) = i\frac{\partial\Phi(\mathbf{R},t)}{\partial t}$$

$$\Phi(\mathbf{R},t) = \sum_n c_n \Phi_n(\mathbf{R}) e^{-iE_n t}$$

In imaginary time $\tau = it$, with energy shift E_T

$$(\hat{H} - E_T)\Phi(\mathbf{R},\tau) = -\frac{\partial\Phi(\mathbf{R},\tau)}{\partial\tau}$$

$$\Phi(\mathbf{R},\tau) = \sum_n c_n \Phi_n(\mathbf{R}) e^{-(E_n - E_T)\tau}$$

Diffusion Monte Carlo

The time-dependent Schrödinger equation

$$\hat{H}\Phi(\mathbf{R}, t) = i \frac{\partial \Phi(\mathbf{R}, t)}{\partial t}$$

$$\Phi(\mathbf{R}, t) = \sum_n c_n \Phi_n(\mathbf{R}) e^{-iE_n t}$$

In imaginary time $\tau = it$, with energy shift E_T

$$(\hat{H} - E_T) \Phi(\mathbf{R}, \tau) = - \frac{\partial \Phi(\mathbf{R}, \tau)}{\partial \tau}$$

$$\Phi(\mathbf{R}, \tau) = \sum_n c_n \Phi_n(\mathbf{R}) e^{-(E_n - E_T)\tau}$$

Diffusion Monte Carlo

Application:

- Set $\Phi(\mathbf{R}, 0) = \Psi(\mathbf{R})$.
- $\lim_{\tau \rightarrow \infty} \Phi(\mathbf{R}, \tau) = \Phi_0(\mathbf{R})$.
- DMC projects out $\Phi_0(\mathbf{R})$ as $\tau \rightarrow \infty$.
- Implementation: ensemble of configurations subjected to drift+diffusion+branching in a discretized timeline.

The DMC method:

- is a more complicated, expensive method
- is a more accurate method
- (for fermions) only depends on the quality of the nodes of $\Psi(\mathbf{R})$ (nodes \equiv region where $\Psi(\mathbf{R}) = 0$).

Diffusion Monte Carlo

Application:

- Set $\Phi(\mathbf{R}, 0) = \Psi(\mathbf{R})$.
- $\lim_{\tau \rightarrow \infty} \Phi(\mathbf{R}, \tau) = \Phi_0(\mathbf{R})$.
- DMC projects out $\Phi_0(\mathbf{R})$ as $\tau \rightarrow \infty$.
- Implementation: ensemble of configurations subjected to drift+diffusion+branching in a discretized timeline.

The DMC method:

- is a more **complicated, expensive** method
- is a more **accurate** method
- (for fermions) only depends on the quality of the **nodes** of $\Psi(\mathbf{R})$ (nodes \equiv region where $\Psi(\mathbf{R}) = 0$).

The Jastrow factor

The Slater-Jastrow wave function

$$\Psi(\mathbf{R}) = \exp[J(\mathbf{R})] \Psi_S(\mathbf{R})$$

The Jastrow factor

$$J(\mathbf{R}) = J_{e-e}(\mathbf{R}) + J_{e-N}(\mathbf{R}) + J_{e-e-N}(\mathbf{R}) + \dots$$

The Jastrow factor:

- is a **compact** parametrization to describe electronic correlation
- does not change the **nodes** of the wave function, thus leaving the DMC energy unchanged

The Jastrow factor

The Slater-Jastrow wave function

$$\Psi(\mathbf{R}) = \exp[J(\mathbf{R})] \Psi_S(\mathbf{R})$$

The Jastrow factor

$$J(\mathbf{R}) = J_{e-e}(\mathbf{R}) + J_{e-N}(\mathbf{R}) + J_{e-e-N}(\mathbf{R}) + \dots$$

The Jastrow factor:

- is a compact parametrization to describe electronic correlation
- does not change the nodes of the wave function, thus leaving the DMC energy unchanged

The Jastrow factor

The Slater-Jastrow wave function

$$\Psi(\mathbf{R}) = \exp[J(\mathbf{R})] \Psi_S(\mathbf{R})$$

The Jastrow factor

$$J(\mathbf{R}) = J_{e-e}(\mathbf{R}) + J_{e-N}(\mathbf{R}) + J_{e-e-N}(\mathbf{R}) + \dots$$

The Jastrow factor:

- is a **compact** parametrization to describe electronic correlation
- does not change the **nodes** of the wave function, thus leaving the DMC energy unchanged

Backflow transformations

The Slater-Jastrow-backflow wave function

$$\Psi(\mathbf{R}) = \exp[J(\mathbf{R})] \Psi_S[\mathbf{X}(\mathbf{R})]$$

The backflow transformation

$$\mathbf{x}_i(\mathbf{R}) = \mathbf{r}_i + \xi_i(\mathbf{R})$$

$$\xi_i(\mathbf{R}) = \xi_{e-e}(\mathbf{R}) + \xi_{e-N}(\mathbf{R}) + \xi_{e-e-N}(\mathbf{R}) + \dots$$

The backflow transformation:

- is a compact parametrization to describe electronic correlation
- changes the nodes of the wave function
- is relatively expensive (but not necessarily affecting the scaling of QMC with system size)

Backflow transformations

The Slater-Jastrow-backflow wave function

$$\Psi(\mathbf{R}) = \exp[J(\mathbf{R})] \Psi_S[\mathbf{X}(\mathbf{R})]$$

The backflow transformation

$$\mathbf{x}_i(\mathbf{R}) = \mathbf{r}_i + \boldsymbol{\xi}_i(\mathbf{R})$$

$$\boldsymbol{\xi}_i(\mathbf{R}) = \boldsymbol{\xi}_{e-e}(\mathbf{R}) + \boldsymbol{\xi}_{e-N}(\mathbf{R}) + \boldsymbol{\xi}_{e-e-N}(\mathbf{R}) + \dots$$

The backflow transformation:

- is a compact parametrization to describe electronic correlation
- changes the nodes of the wave function
- is relatively expensive (but not necessarily affecting the scaling of QMC with system size)

Backflow transformations

The Slater-Jastrow-backflow wave function

$$\Psi(\mathbf{R}) = \exp[J(\mathbf{R})] \Psi_S[\mathbf{X}(\mathbf{R})]$$

The backflow transformation

$$\mathbf{x}_i(\mathbf{R}) = \mathbf{r}_i + \boldsymbol{\xi}_i(\mathbf{R})$$

$$\boldsymbol{\xi}_i(\mathbf{R}) = \boldsymbol{\xi}_{e-e}(\mathbf{R}) + \boldsymbol{\xi}_{e-N}(\mathbf{R}) + \boldsymbol{\xi}_{e-e-N}(\mathbf{R}) + \dots$$

The backflow transformation:

- is a **compact** parametrization to describe electronic correlation
- changes the **nodes** of the wave function
- is relatively **expensive** (but **not** necessarily affecting the scaling of QMC with system size)

Fixed-form Jastrow factor terms

Fixed-form Jastrow factor terms

$$J_{e-e}(\mathbf{R}) = \sum_{i < j}^M f(r_{ij}) \sum_{\nu}^p \lambda_{\nu}^{P_{ij}} r_{ij}^{\nu}$$

$$J_{e-N}(\mathbf{R}) = \sum_i^N \sum_I^M f(r_{iI}) \sum_{\mu}^q \lambda_{\mu}^{S_{iI}} r_{iI}^{\mu}$$

$$J_{e-e-N}(\mathbf{R}) = \sum_{i < j}^N \sum_I^M f(r_{iI}) f(r_{jI}) \sum_{\nu_{ij}}^p \sum_{\mu_{iI} \mu_{jI}}^q \lambda_{\nu_{ij} \mu_{iI} \mu_{jI}}^{P_{ij} S_{iI} S_{jI}} r_{ij}^{\nu_{ij}} r_{iI}^{\mu_{iI}} r_{jI}^{\mu_{jI}}$$

General Jastrow factor

General Jastrow factor

$$\begin{aligned}
 J_{n,m}(\mathbf{R}) = & \sum_{i_1 < \dots < i_n}^N \sum_{I_1 < \dots < I_m}^M \sum_{\alpha < \beta}^p \sum_{\alpha=1, \gamma=1}^q \lambda_{\{v\}\{\mu\}}^{\{P\}\{S\}} \times \\
 & \times \left(\prod_{\alpha < \beta}^n \Phi_{v_{i_\alpha i_\beta}}^{P_{\sigma(i_\alpha)\sigma(i_\beta)}}(\mathbf{r}_{i_\alpha i_\beta}) \right) \left(\prod_{\alpha}^n \prod_{\gamma}^m \Theta_{\mu_{i_\alpha I_\gamma}}^{S_{\sigma(i_\alpha)I_\gamma}}(\mathbf{r}_{i_\alpha I_\gamma}) \right)
 \end{aligned}$$

Fixed-form backflow terms

Fixed-form backflow terms

$$\xi_{e-e}^i(\mathbf{R}) = \sum_{j \neq i}^M f(r_{ij}) \sum_{\nu}^p \lambda_{\nu}^{P_{ij}} r_{ij}^{\nu} \mathbf{r}_{ij}$$

$$\xi_{e-N}^i(\mathbf{R}) = \sum_I^M f(r_{iI}) \sum_{\mu}^q \omega_{\mu}^{S_{iI}} r_{iI}^{\mu} \mathbf{r}_{iI}$$

$$\begin{aligned} \xi_{e-e-N}^i(\mathbf{R}) = & \sum_{j \neq i}^N \sum_I^M f(r_{iI}) f(r_{jI}) \sum_{\nu_{ij}}^p \sum_{\mu_{iI} \mu_{jI}}^q \left[\lambda_{\nu_{ij} \mu_{iI} \mu_{jI}}^{P_{ij} S_{iI} S_{jI}} r_{ij}^{\nu} r_{iI}^{\mu_{iI}} r_{jI}^{\mu_{jI}} \mathbf{r}_{ij} + \right. \\ & \left. + \omega_{\nu_{ij} \mu_{iI} \mu_{jI}}^{P_{ij} S_{iI} S_{jI}} r_{ij}^{\nu} r_{iI}^{\mu_{iI}} r_{jI}^{\mu_{jI}} \mathbf{r}_{iI} \right] \end{aligned}$$

General backflow transformation

General backflow transformation

$$\xi_{n,m}^i(\mathbf{R}) = \sum_{i_2 < \dots < i_n}^N \sum_{I_1 < \dots < I_m}^M \sum_{\{\nu_{\alpha i \beta}\}_{\alpha < \beta}}^p \sum_{\{\mu_{\alpha I \gamma}\}_{\alpha=1, \gamma=1}^q} \times$$

$$\left[\sum_{k=2}^N [ik] \{P\}\{S\} \lambda_{\{v\}\{\mu\}} \left(\prod_{\alpha < \beta}^n \Phi_{\nu_{\alpha i \beta}}^{P_{\sigma(i\alpha)\sigma(i\beta)}}(\mathbf{r}_{i\alpha i\beta}) \right) \left(\prod_{\alpha}^n \prod_{\gamma}^m \Theta_{\mu_{\alpha I \gamma}}^{S_{\sigma(i\alpha)I\gamma}}(\mathbf{r}_{i\alpha I\gamma}) \right) \mathbf{r}_{ik} + \right.$$

$$\left. \sum_{K=2}^N [iK] \{P\}\{S\} \omega_{\{v\}\{\mu\}} \left(\prod_{\alpha < \beta}^n \Delta_{\nu_{\alpha i \beta}}^{P_{\sigma(i\alpha)\sigma(i\beta)}}(\mathbf{r}_{i\alpha i\beta}) \right) \left(\prod_{\alpha}^n \prod_{\gamma}^m \Lambda_{\mu_{\alpha I \gamma}}^{S_{\sigma(i\alpha)I\gamma}}(\mathbf{r}_{i\alpha I\gamma}) \right) \mathbf{r}_{iK} \right]$$

Summary

- Some work to do...

Summary

- Some work to do...