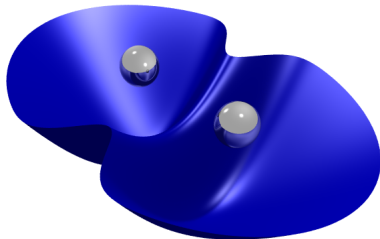


Cold Atoms: A quantum test bed for magnetism?

Pascal Bugnion and Gareth Conduit



Why cold atoms?

They are useful to **simulate** complex many-body systems.

Simulating electronic systems

Prepare a mixture of ${}^6\text{Li}$ atoms, such that the **total angular momentum** for each atom is $F = 1/2$.

This gives two hyperfine states: $m_F = +1/2$ and $m_F = -1/2$.

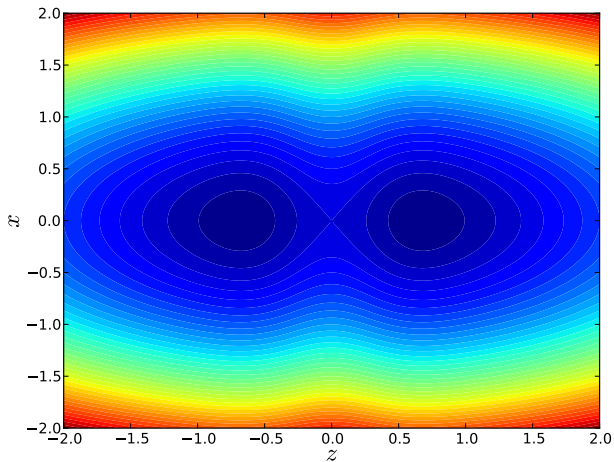
Simulating electronic systems

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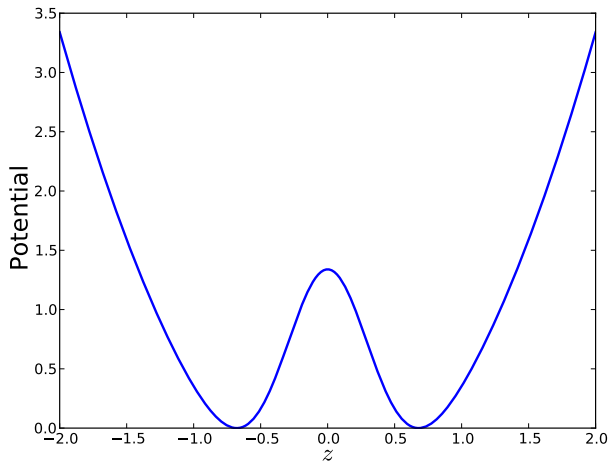
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The interaction between different atoms can be **tuned** by changing the underlying magnetic field.

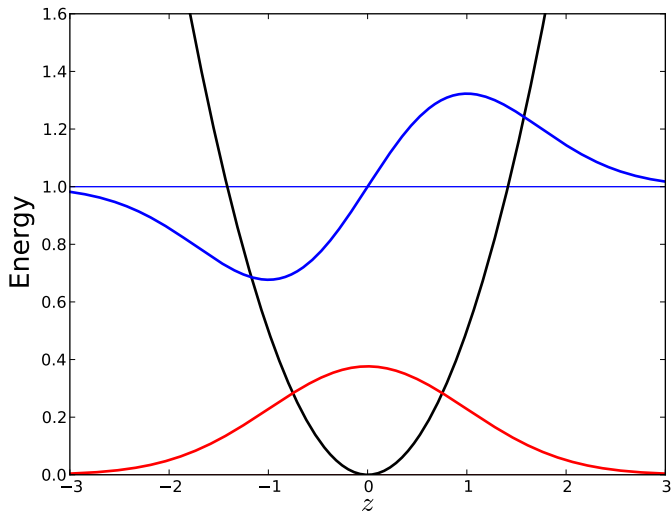
Experimental setup



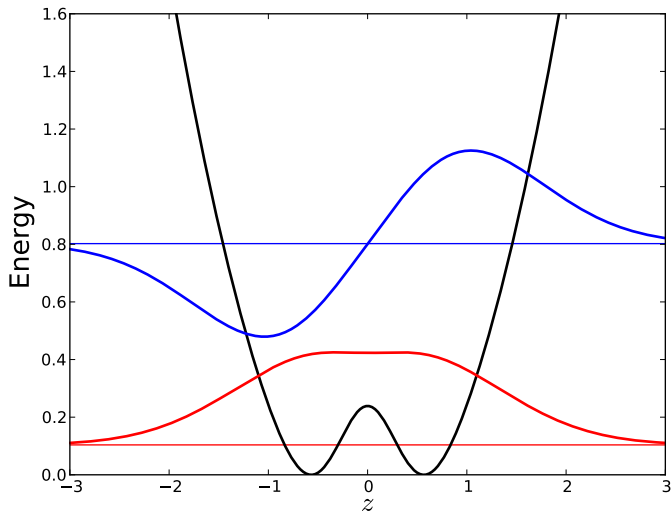
Experimental setup



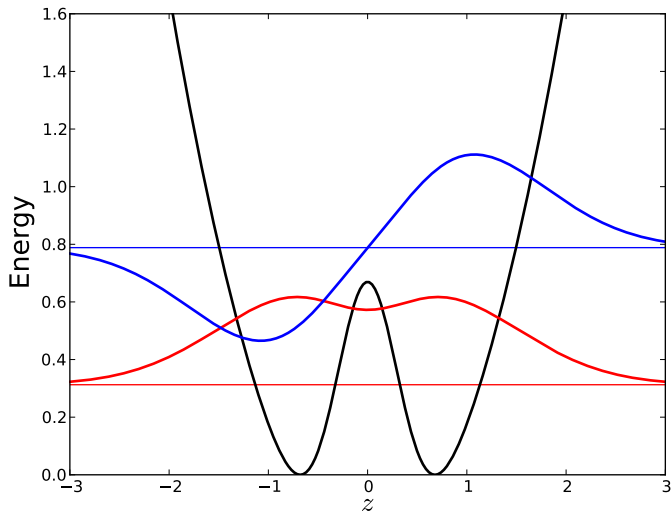
Low-barrier limit



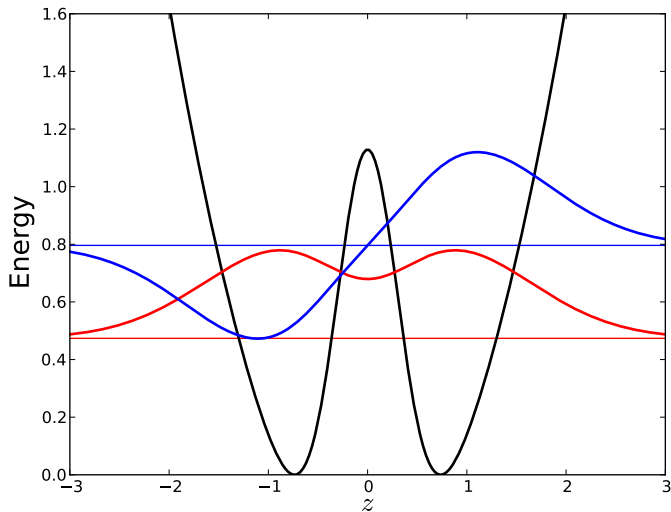
Increasing barrier height



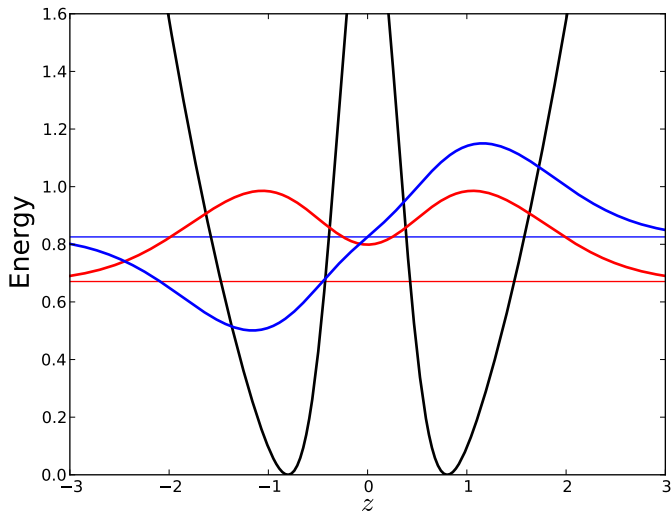
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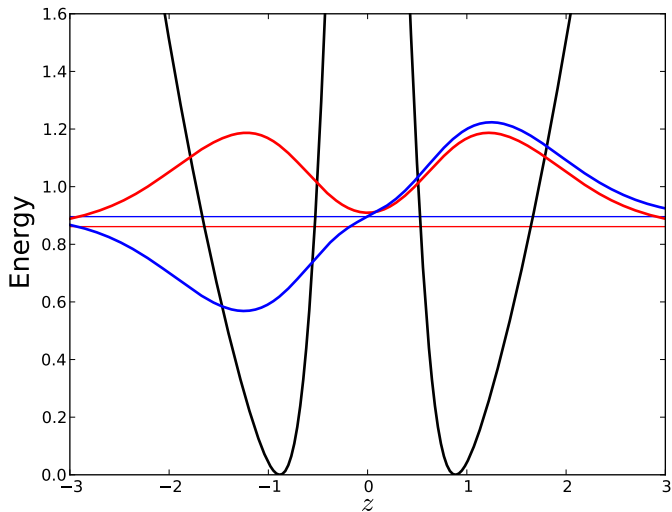
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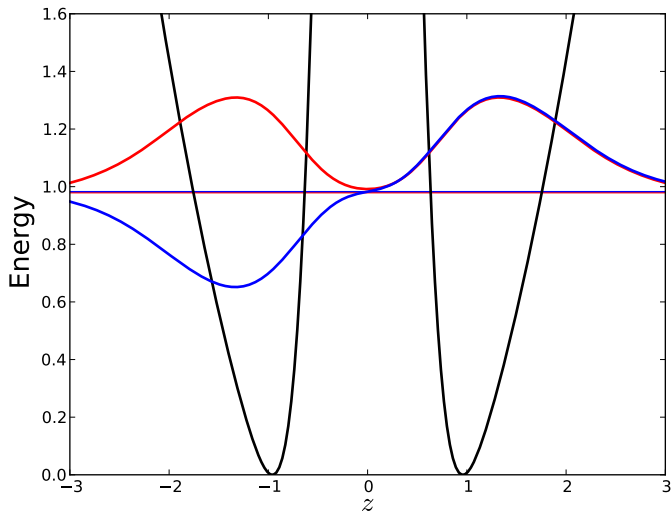
Increasing barrier height



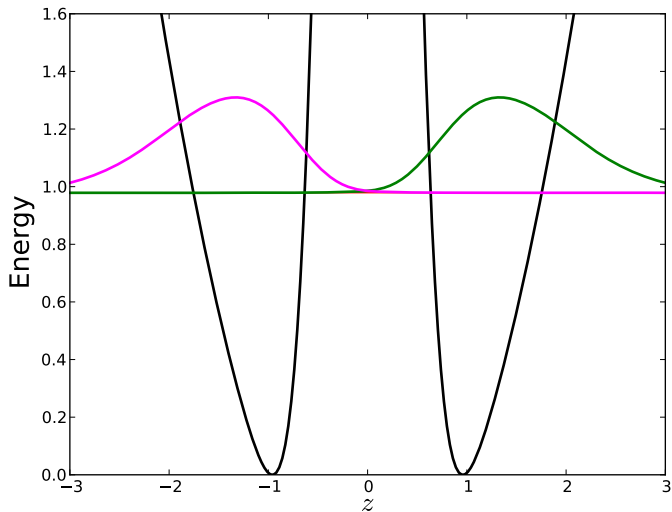
Increasing barrier height



High barrier



High barrier

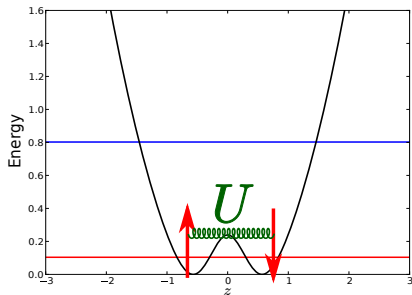


Simulating magnetism

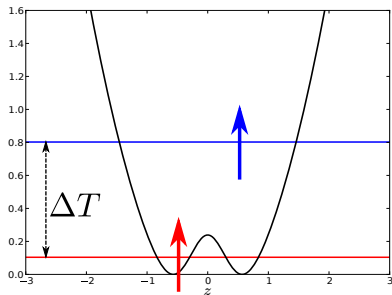
Ashcroft and Mermin on magnetism:

“ *The development of a tractable model of a magnetic material, capable of describing both the characteristic electron spin correlations as well as the electronic transport properties [. . .] remains one of the major unsolved problems of modern solid state theory.* ”

Diamagnetic-Ferromagnetic transition

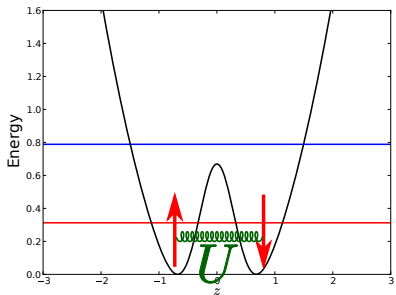


$$E_{\uparrow\downarrow} = U$$

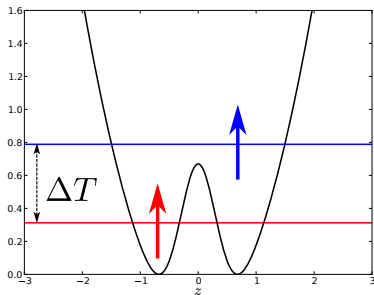


$$E_{\uparrow\uparrow} = \Delta T$$

Diamagnetic-Ferromagnetic transition

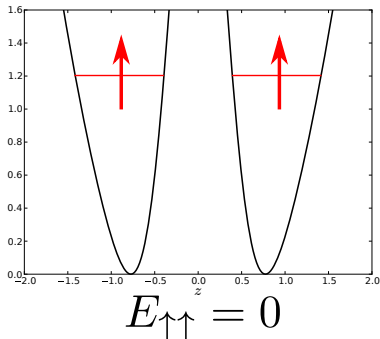
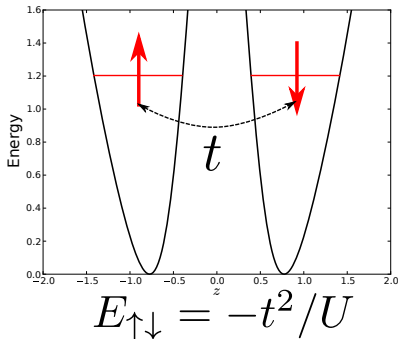


$$E_{\uparrow\downarrow} = U$$



$$E_{\uparrow\uparrow} = \Delta T$$

Ferromagnetic-Antiferromagnetic transition



Full Configuration Interaction

We want to find the stationary states of the **time-independent Schrödinger equation**, that is, all $E, |\Psi\rangle$ pairs such that

$$\hat{H} |\Psi\rangle = E |\Psi\rangle$$

Full Configuration Interaction

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$$\hat{H} |\Psi\rangle = E |\Psi\rangle$$

where

$$\hat{H} = \sum_i \hat{h}(\mathbf{r}_i) + \sum_{i>j} \hat{V}(\mathbf{r}_i, \mathbf{r}_j)$$

Full Configuration Interaction (2)

The Hamiltonian eigenvalue equation $\hat{H} |\Psi\rangle = E |\Psi\rangle$ can always be solved by expanding $|\Psi\rangle$ in a basis with the correct symmetry.

Let the basis functions be called $\{|i\rangle\}$. The method proceeds as follows:

- 1 Write the Hamiltonian matrix in the $\{|i\rangle\}$ basis: $\langle i | \hat{H} | j \rangle$.
- 2 Diagonalise this matrix. The eigenvalues are the energies and the eigenvectors are the stationary states $|\Psi\rangle$ expressed in the $\{|i\rangle\}$ basis.

Full Configuration Interaction (3)

- 1 Define a set of one-electron orbitals $\{\phi_i\}$.
- 2 Build all n -electron Slater determinants:

$$|D_\alpha\rangle = |\phi_{\alpha_1}\phi_{\alpha_2}\dots\phi_{\alpha_n}\rangle$$

that you can in this basis.

- 3 Construct all matrix elements $\langle D_\alpha | \hat{H} | D_\beta \rangle$.
- 4 Diagonalise the resultant matrix.

Full Configuration Interaction (4)

Pros

- Arbitrarily improvable. Gives the right answer in the limit of complete basis set.
- Variational.
- Gives excited states.
- Conceptually simple (and easy enough to implement).

Full Configuration Interaction (4)

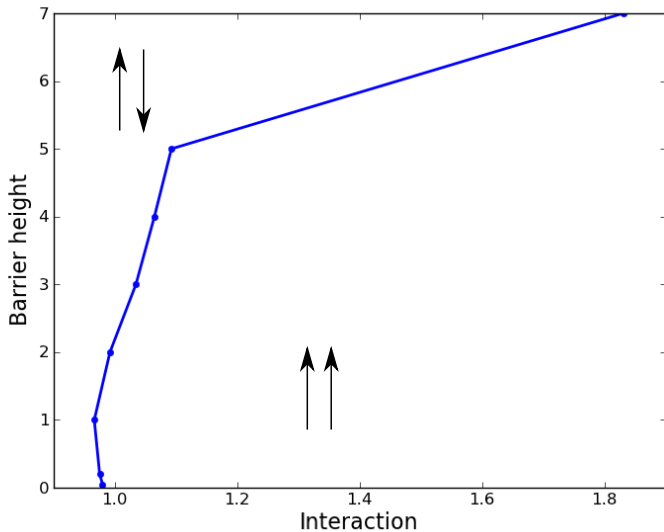
Pros

- Arbitrarily improvable. Gives the right answer in the limit of complete basis set.
- Variational.
- Gives excited states.
- Conceptually simple (and easy enough to implement).

Cons

Very expensive: for N orbitals and n electrons, there are ${}^N C_n$ Slater determinants. This gives a cost of $O(n!^3)$.

Phase diagram





Thank you for your attention!