

The out of equilibrium birth of a (holographic) superfluid

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MIT



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Dynamical phase transitions

$$\epsilon(t) = 1 - \frac{T(t)}{T_c} = t/\tau_Q$$

$$\tau_{eq} = \tau_0 |\epsilon|^{-\nu z}$$

$$\xi_{eq} = \xi_0 |\epsilon|^{-\nu}$$

T_c

Unbroken Phase

$$\langle \psi \rangle = 0$$

Broken phase

$$\langle \psi \rangle \neq 0$$

$$\langle \psi \rangle = \Delta(x, t) e^{i\theta(x, t)} ? \quad t$$

Kibble

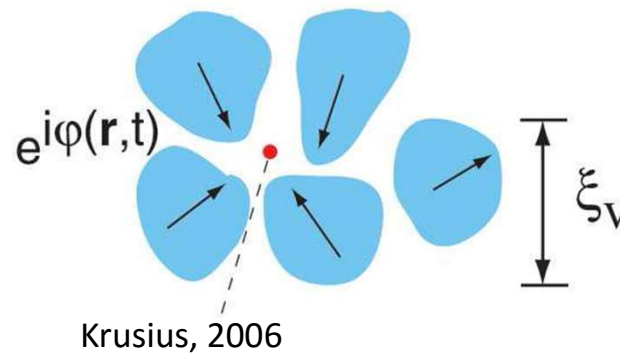
J.

Phys. A: Math. Gen. 9: 1387. (1976)

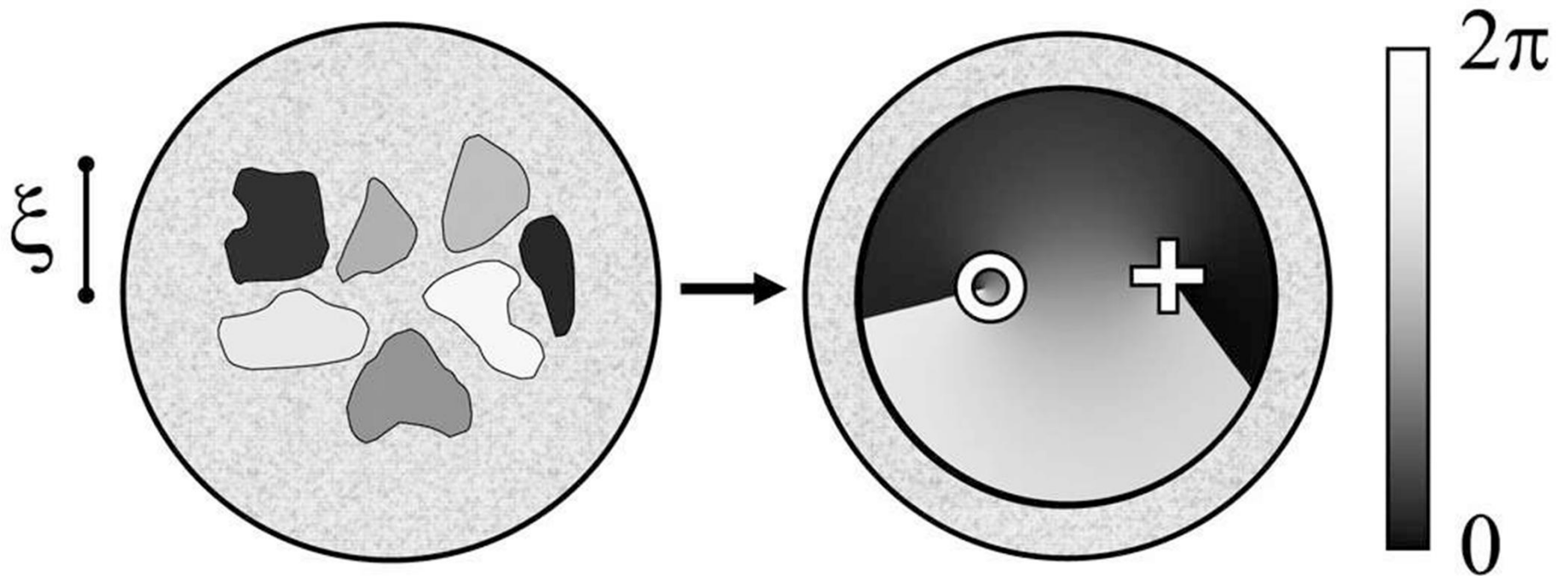
Causality

Vortices in the sky

Cosmic strings



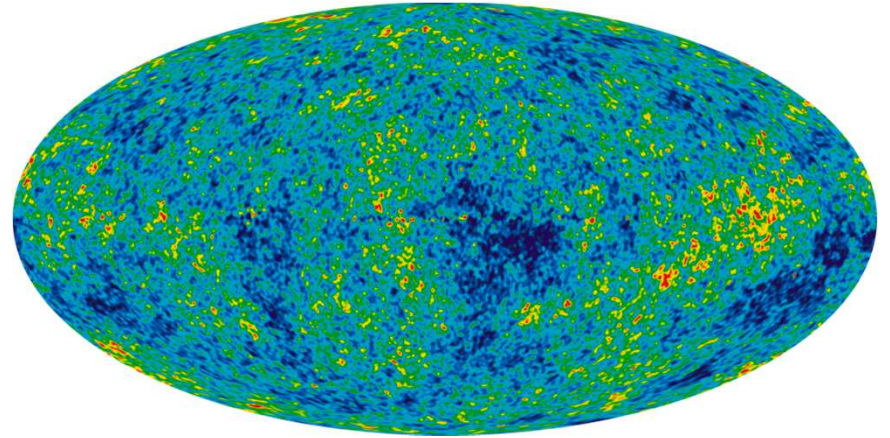
Generation of Structure



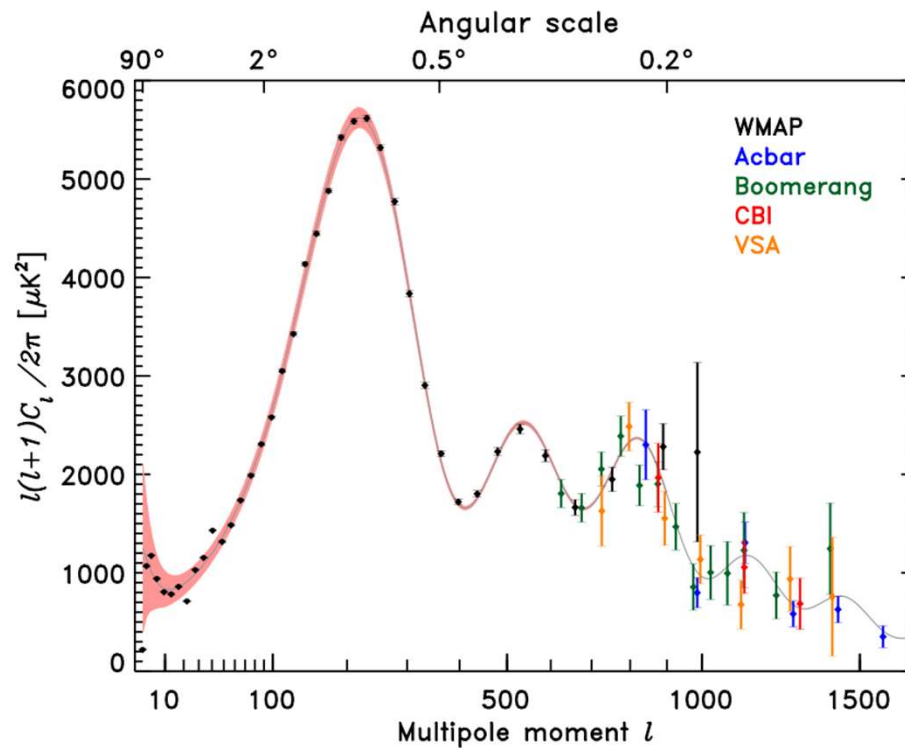
Weyler, Nature 2008

No evidence so far !

CMB, galaxy distributions...



NASA/WMAP



Cosmological experiments in superfluid helium?

Doable for ^4He !!



Zurek

W. H. Zurek

Theoretical Astrophysics, Los Alamos National Laboratory,
Los Alamos, New Mexico 87545, USA Nature 317, 505 (1985)

$$T \approx T_c$$

2nd order



Scaling

$$\tau(T_c) = \infty$$

$$\epsilon(t) = 1 - \frac{T(t)}{T_c} = t/\tau_Q$$

$$t = -\hat{t} \equiv -t_{freeze}$$

$$t = \hat{t} \equiv t_{freeze}$$

Non adiabatic evolution

Defect generation!

$$\epsilon(t) = t/\tau_Q$$

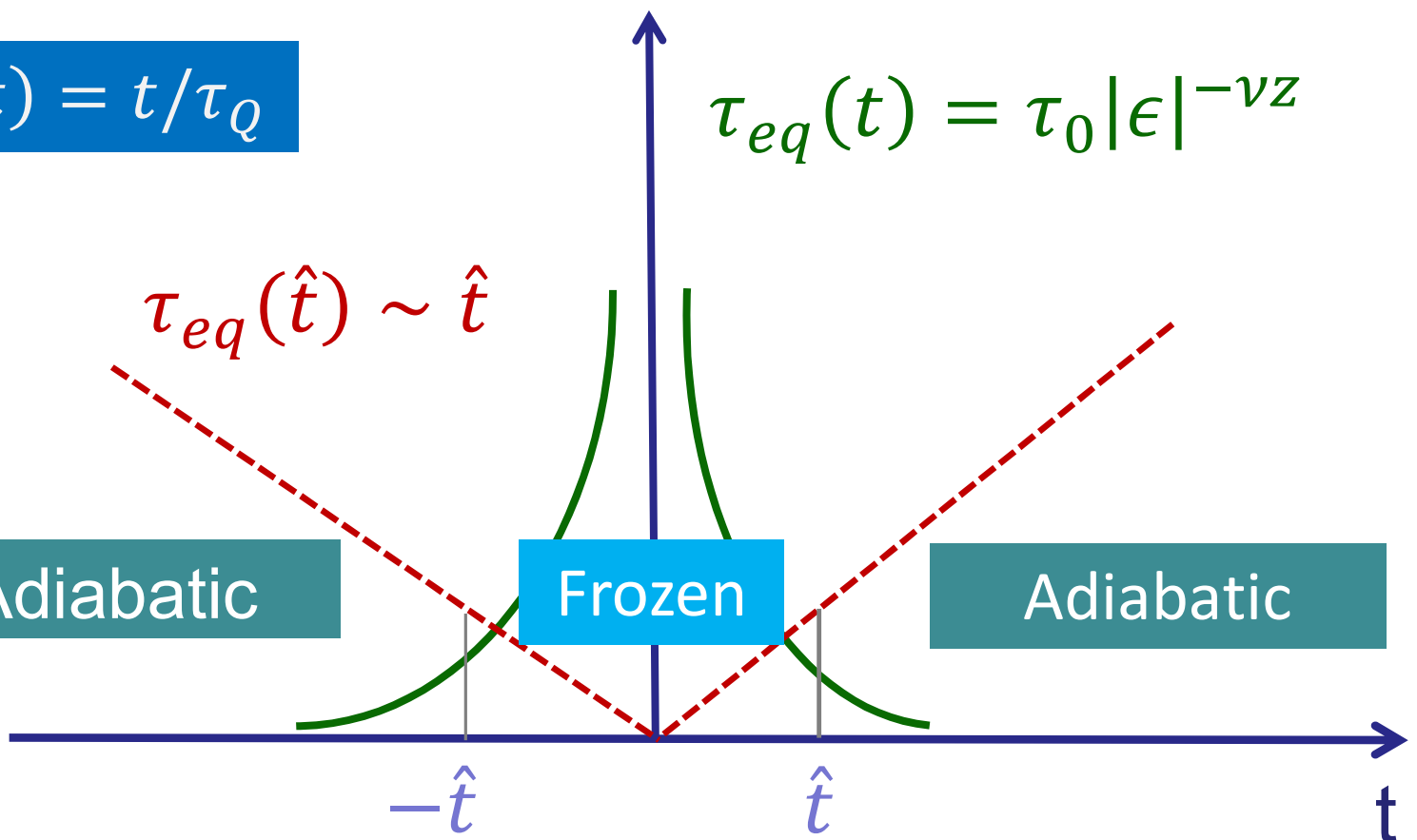
$$\tau_{eq}(t) = \tau_0 |\epsilon|^{-\nu z}$$

$$\tau_{eq}(\hat{t}) \sim \hat{t}$$

Adiabatic

Frozen

Adiabatic



$$\hat{\xi} = \xi_0 |\hat{\epsilon}|^{-\nu} = \xi_0 (\tau_Q/\tau_0)^{\nu/(1+\nu z)}$$

Kibble-Zurek mechanism

$$\rho \sim \hat{\xi}^{-d} \sim \tau_Q^{-d\nu/(1+\nu z)}$$

Generation of defects in superfluid ^4He as an analogue of the formation of cosmic strings

P. C. Hendry*, N. S. Lawson*, R. A. M. Lee*,
P. V. E. McClintock* & C. D. H. Williams†

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Transient
attenuation of
second sound
amplitude

But vortices
induced by stirring
up!

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PHYSICAL REVIEW LETTERS

26 OCTOBER 1998

Nonappearance of Vortices in Fast Mechanical Expansions of Liquid ^4He through the Lambda Transition

M. E. Dodd,¹ P. C. Hendry,¹ N. S. Lawson,¹ P. V. E. McClintock,¹ and C. D. H. Williams²

¹*Department of Physics, Lancaster University, Lancaster, LA1 4YB, United Kingdom*

²*Department of Physics, University of Exeter, Stocker Road, Exeter EX4 4QL, United Kingdom*

No vortices in ^4He !!

G. Karra, R. J. Rivers, PRL. 81, 3707 (1998)

Vortex formation in neutron-irradiated superfluid ^3He as an analogue of cosmological defect formation

Ruutu, Nature 382, 334-336 (1996)

Laboratory simulation of cosmic string formation in the early Universe using superfluid ^3He

C. Bäuerle et al. Nature 382, 332 (1996)

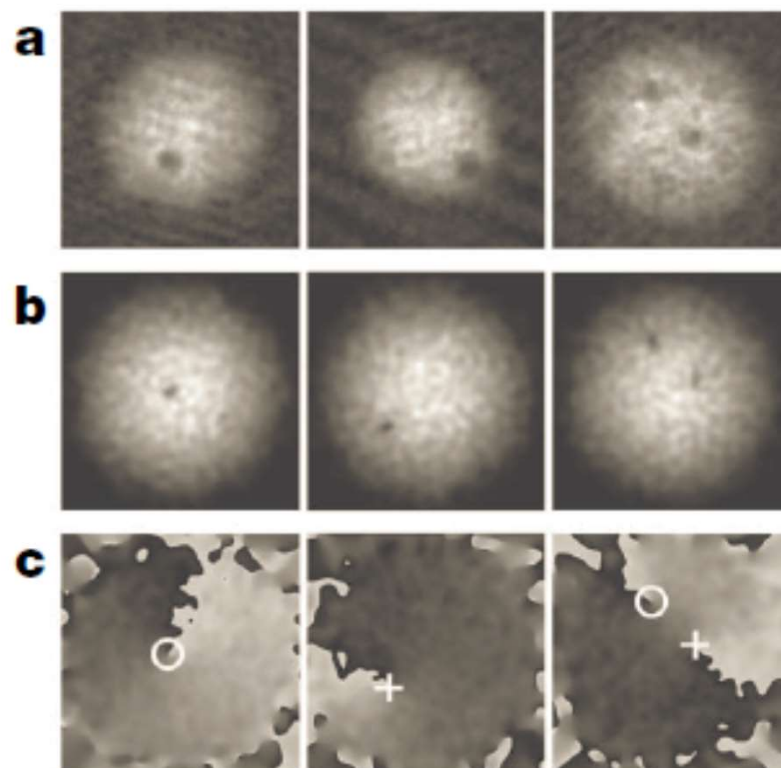
Thin SC films, nematic liquid crystal..



LETTERS

Spontaneous vortices in the formation of Bose–Einstein condensates

Chad N. Weiler¹, Tyler W. Neely¹, David R. Scherer¹, Ashton S. Bradley²†, Matthew J. Davis² & Brian P. Anderson¹



ARTICLE

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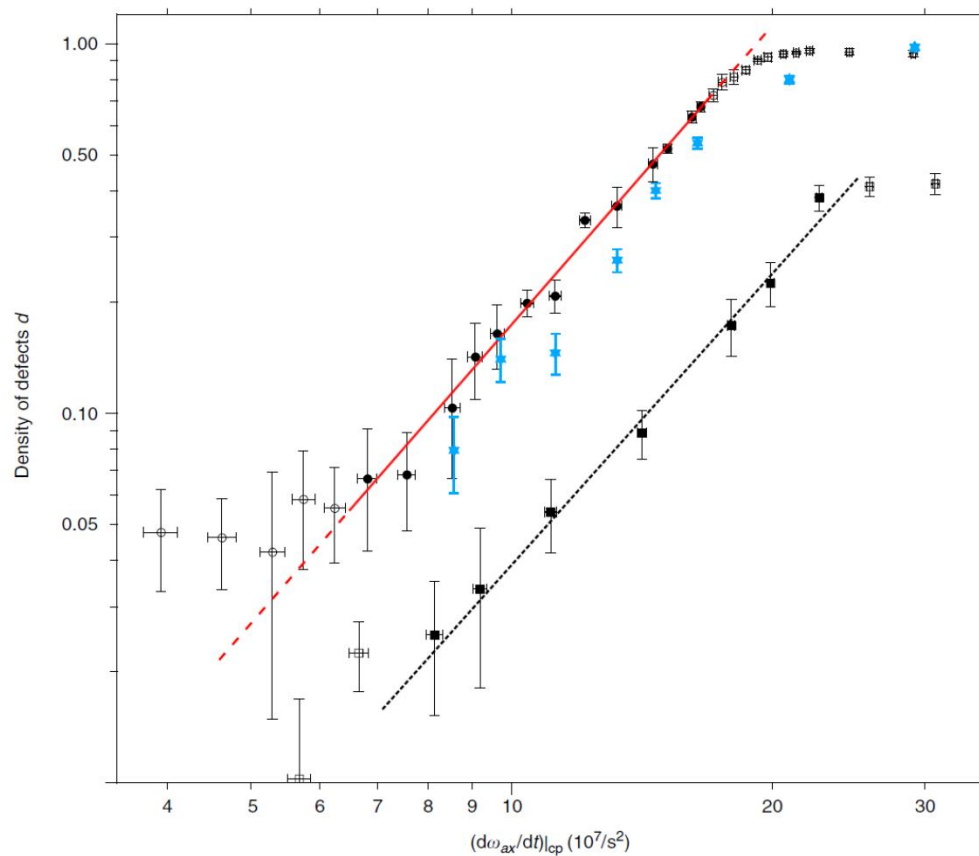
DOI: 10.1038/ncomms3290

Observation of the Kibble–Zurek scaling law for defect formation in ion crystals

S. Ulm¹, J. Roßnagel¹, G. Jacob¹, C. Degünther¹, S.T. Dawkins¹, U.G. Poschinger¹, R. Nigmatullin^{2,3}, A. Retzker⁴, M.B. Plenio^{2,3}, F. Schmidt-Kaler¹ & K. Singer¹

KZ scaling with the
quench speed

Too few defects



Extension to quantum phase transitions

Zurek, Zoller, et al, "Dynamics of a quantum phase transition." , PRL 95.10 (2005): 105701.

Demonstration of KZ scaling in a 1d Ising chain in transverse field

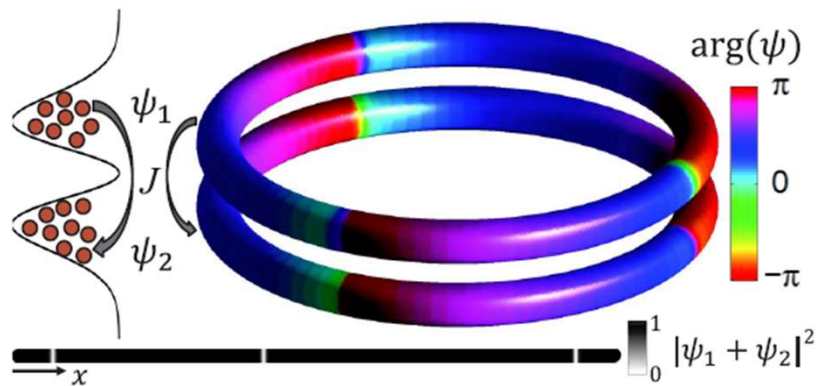
Dziarmaga "Dynamics of a quantum phase transition: Exact solution of the quantum Ising model." PRL 95.24 (2005): 245701.

Calculation of correlation functions

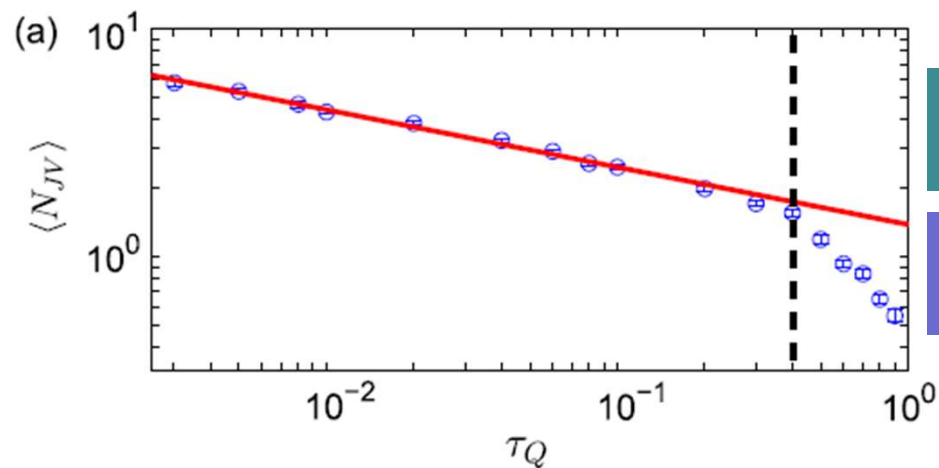
Kibble-Zurek problem: Universality and the scaling limit
PRB 86, 064304, (2012), Gubser, Sondhi et al.

Kibble-Zurek Scaling and its Breakdown for Spontaneous Generation of Josephson Vortices in Bose-Einstein Condensates

Shih-Wei Su,¹ Shih-Chuan Gou,² Ashton Bradley,³ Oleksandr Fialko,⁴ and Joachim Brand⁴



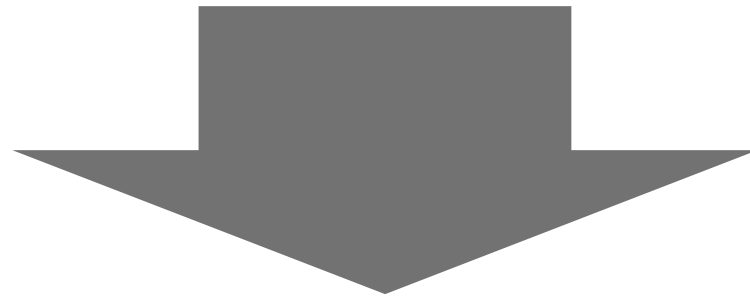
Stochastic
Gross-Pitaevskii



Breaking of KZ scaling

Too few vortices !

Adiabatic at t_{freeze} ?
Defects without a
condensate?



$t_{eq} > t > t_{freeze}$ is relevant

PRX 5, 021015 (2015)

Slow Quenches

Linear response

$t > t_{\text{freeze}}$

Scaling

KZ

Frozen

Adiabatic

US

Frozen

Coarsening

Adiabatic

t_{freeze}

t_{eq}

$$\frac{t_{\text{eq}}}{t_{\text{freeze}}} \sim (\log R)^{\frac{1}{1+\nu z}}$$

$$|\psi|^2(t) \propto e^{a_2 \bar{t}^{1+z\nu}}$$

$$|\psi|^2(\epsilon) \propto \epsilon^{2\beta}$$

$$\Lambda = (d - z)\nu - 2\beta$$

$$R \sim \xi^{-1} \tau_Q^{\Lambda/1+\nu z}$$

$$\gamma = \frac{1 + (z - 2)\nu}{2(1 + z\nu)}$$

$$\rho(t_{\text{eq}}) \sim [\log R]^\gamma \rho_{\text{KZ}}$$

Non adiabatic growth after t_{freeze}

$$C(t, \mathbf{r}) \equiv \langle \psi^*(t, \mathbf{x} + \mathbf{r}) \psi(t, \mathbf{x}) \rangle$$

$$\psi(t, \mathbf{q}) = \int dt' G_R(t, t', q) \varphi(t, \mathbf{q})$$

$$\langle \varphi^*(t, \mathbf{x}) \varphi(t', \mathbf{x}') \rangle = \zeta \delta(t - t') \delta(\mathbf{x} - \mathbf{x}')$$

$$G_R(t, t', q) = \theta(t - t') H(q) e^{-i \int_{t'}^t dt'' \mathbf{w}_0(\epsilon(t''), q)}$$

$$C(t, q) = \int dt' \zeta |G_R(t, t', q)|^2$$

Linear response

$t > t_{freeze}$

$|\partial_t \log \mathbf{w}_0| < |\mathbf{w}_0|$

$$C(t, q) = \int_{t_{freeze}}^t dt' \zeta |H(q)|^2 e^{2 \int_{t'}^t dt'' \text{Im } \mathbf{w}_0(\epsilon(t''), q)} + \dots$$

$$\mathbf{w}_0(\epsilon, q) = \epsilon^{z\nu} h(q\epsilon^{-\nu})$$

$$\text{Im } \mathbf{w}_0 = -a\epsilon^{(z-2)\nu} q^2 + b\epsilon^{z\nu} + \dots, \quad q_{max} \sim \epsilon(t)^\nu$$

$$\text{Im } \mathbf{w}_0 > 0$$

Unstable Modes



Growth

$\langle \psi(t) \rangle t > t_{freeze}$

Protocol

$$\epsilon(t) = t/\tau_Q$$

$$t \in (t_i, t_f)$$

$$t_i = (1 - T_i/T_c)\tau_Q < 0$$

$$t_f = (1 - T_f/T_c)\tau_Q > 0$$

Slow quenches

$$t_f \geq t_{eq}$$

$$t > t_{freeze}$$

Correlation length
increases

$$C(t, r) \sim |\psi|^2(t) e^{-\frac{r^2}{2\ell_{co}^2(t)}}, \quad \bar{t} \equiv \frac{t}{t_{freeze}}$$

$$\ell_{co}(\bar{t}) = a_3 \zeta_{freeze} \bar{t}^{\frac{1+(z-2)\nu}{2}}$$

Condensate growth

$$|\psi|^2(t) \sim \tilde{\varepsilon}(t) e^{a_2 \bar{t}^{1+z\nu}}$$

$$\tilde{\varepsilon}(t) \equiv \zeta t_{freeze} \ell_{co}^{-d}(t)$$

Adiabatic evolution

$$t = t_{eq} \gg t_{freeze}$$

$$|\psi|^2(t = t_{eq}) \sim |\psi|_{eq}^2(\epsilon(t_{eq}))$$

Defects

$$\rho(t_{eq}) \sim 1/\ell_{co}^{d-D}(t_{eq}) \sim [\log(\zeta^{-1} \tau_Q \Lambda)]^{-\frac{(d-D)(1+(z-2)\nu)}{2(1+z\nu)}} \rho_{KZ}$$

Fast quenches

$$t_f \ll t_{eq}$$

$$q_{max}(T_f) = \epsilon(t_f)^{d\nu}$$

Breaking of τ_Q scaling

$$KZ \quad t_f < t_{freeze}$$

$$US \quad t_{freeze} \ll t_f \ll t_{eq}$$

Exponential growth

$$|\psi|^2(t) \sim \epsilon_f^{(d-z)\nu} \zeta \exp [2b(t - t_{freeze})\epsilon_f^{\nu z}]$$

Number of defects

$$\rho \sim \begin{cases} \epsilon_f^{(d-D)\nu} & R_f \lesssim O(1) \\ \epsilon_f^{(d-D)\nu} \log^{-\frac{d-D}{2}} R_f & R_f \gg 1 \end{cases}$$

Independent of τ_Q

$$R_f \equiv \frac{\epsilon_f^{2\beta}}{\zeta \epsilon_f^{(d-z)\nu}} \quad \epsilon_f \equiv \frac{T_c - T_f}{T_c}$$

Holography?

Defects survive large
N limit

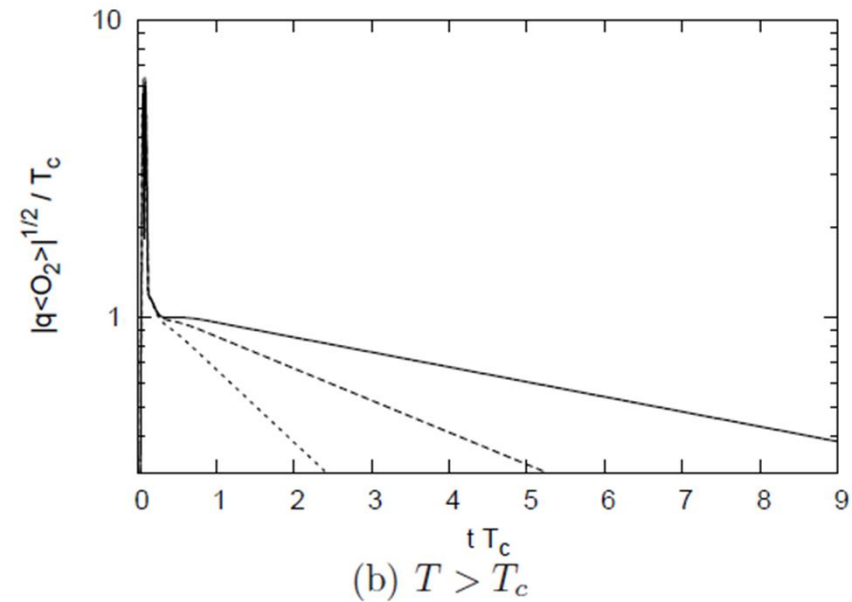
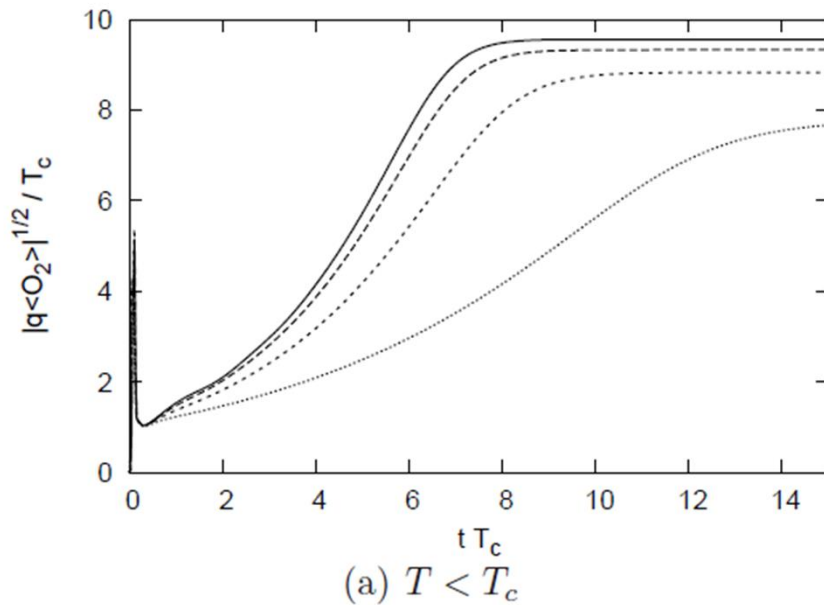
Universality

Real time

Holographic SC out of equilibrium

$$S_{\text{grav}} = \frac{1}{16\pi G_{\text{Newton}}} \int d^4x \sqrt{-G} \left[R + \Lambda + \frac{1}{e^2} \left(-\frac{1}{4} F_{MN} F^{MN} - |D\Phi|^2 - m^2 |\Phi|^2 \right) \right]$$

Gubser, Herzog, Horowitz, Hartnoll

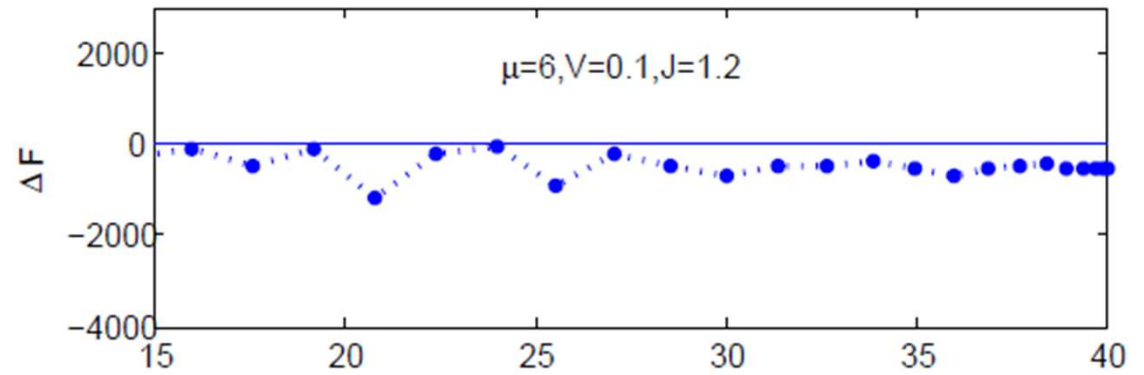
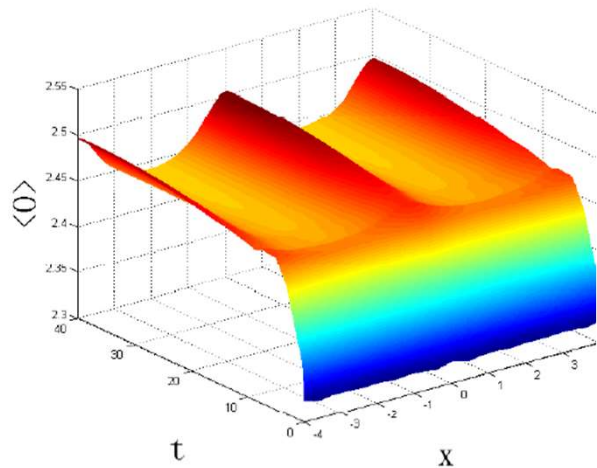


$$|\langle \mathcal{O}_2(t) \rangle| = C_1 \exp(-t/t_{\text{relax}}) + C_2 \quad |\langle \mathcal{O}_2(t) \rangle| = C \exp(-t/t_{\text{relax}})$$

Exponential growth/decay

Murata, et al.,
arXiv:1005.0633

Oscillations in space

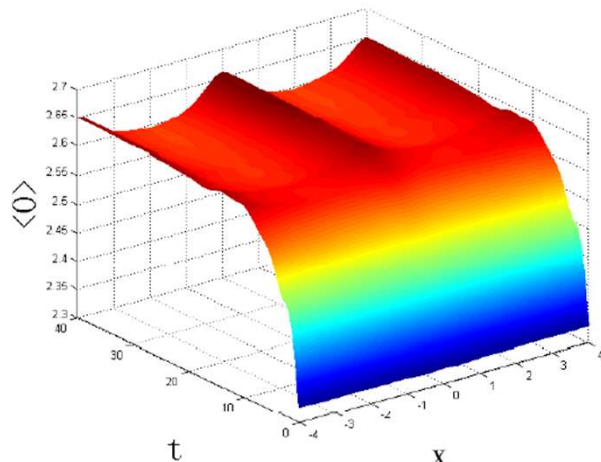


$$\Psi \approx z\psi_1 + \psi(x, t)z^2$$
$$\psi_1(t) = J(1 - \tanh vt)$$

$$\langle O \rangle \sim \psi(x, t)$$

Probe limit

Conservation laws!



Basu et al., arXiv:1308.4061

AGG, Zhang, Bi, arXiv:1308.5398

Oscillations in space: BdG

$$\hat{\xi} = -\vec{\nabla}^2/2m - \mu$$

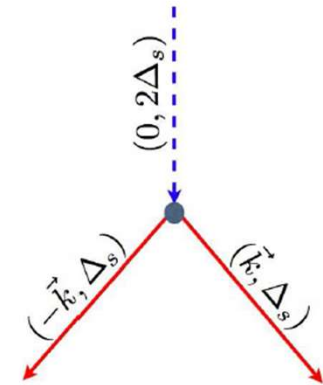
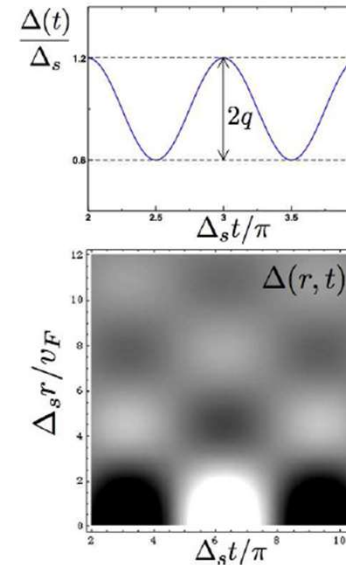
$$i\dot{u}_{\mathbf{p}}(\mathbf{r}, t) = \hat{\xi}u_{\mathbf{p}}(\mathbf{r}, t) + \Delta(\mathbf{r}, t)v_{\mathbf{p}}(\mathbf{r}, t),$$

$$i\dot{v}_{\mathbf{p}}(\mathbf{r}, t) = -\hat{\xi}v_{\mathbf{p}}(\mathbf{r}, t) + \bar{\Delta}(\mathbf{r}, t)u_{\mathbf{p}}(\mathbf{r}, t)$$

$$\Delta(\mathbf{r}, t) = \bar{\Delta}(t) + \delta\Delta(\mathbf{r}, t)$$

$$\delta\Delta(\vec{r}, t) \approx \frac{C e^{\nu_m t} \cos[\Delta_s(t - \tau)] \sin(k_m R) e^{-R^2/l^2(t)}}{\sqrt{\Delta_s t} k_m R}$$

$$l(t) \approx \xi \sqrt{\Delta_s t} \quad \nu_m \approx 2q\Delta_s$$



Conservation
laws

Instability to spatial inhomogeneity

KZ scaling in holographic superconductors

Quantum Quench Across a Zero Temperature
Holographic Superfluid Transition,

P. Basu, D. Das, S. R. Das,
T. Nishioka

arXiv:1211.7076

Kibble-Zurek Scaling in Holographic Quantum
Quench : Backreaction

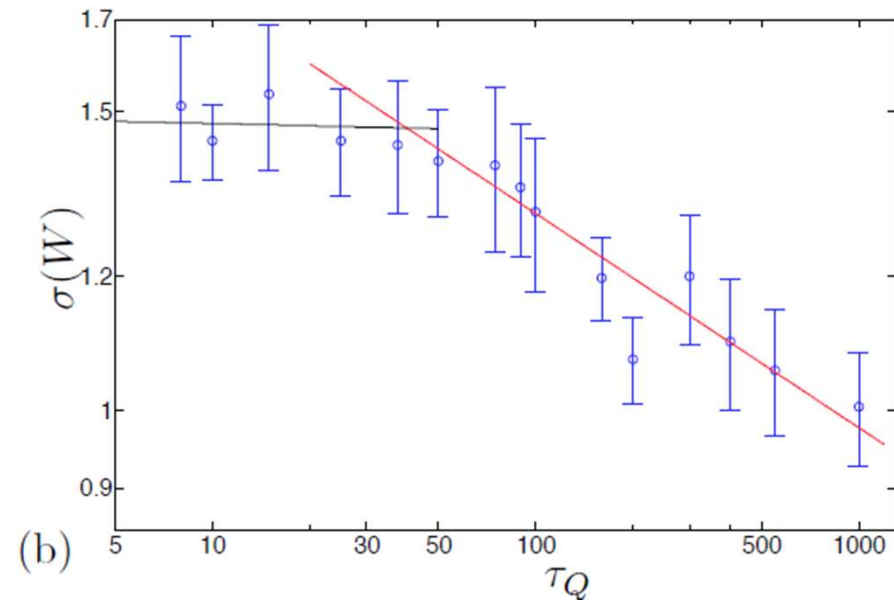
S. R. Das, T. Morita

arXiv:1409.7361

Universal far-from equilibrium
dynamics of a holographic
superconductor

Sonner, Campo, and Zurek

arXiv:1406.2329



Dual gravity theory

$$S_{\text{grav}} = \frac{1}{16\pi G_{\text{Newton}}} \int d^4x \sqrt{-G} \left[R + \Lambda + \frac{1}{e^2} \left(-\frac{1}{4} F_{MN} F^{MN} - |D\Phi|^2 - m^2 |\Phi|^2 \right) \right]$$

$$\Lambda = -3 \quad m^2 = -2$$

Herzog, Horowitz, Hartnoll, Gubser

*AdS*₄

Eddington-Finkelstein
coordinates

$$ds^2 = r^2 g_{\mu\nu}(t, \mathbf{x}, r) dx^\mu dx^\nu + 2drdt$$

Probe limit

$$0 = \nabla_M F^{NM} - J^M,$$

$$0 = (-D^2 + m^2)\Phi.$$

EOM's:

Boundary conditions:

$$r \rightarrow \infty$$

Drive:

No solution of Einstein equations but do not worry, Hubeny 2008

Dictionary:

PDE's in x, y, r, t

$$\Psi = \frac{\psi^{(1)}(x, y, t)}{r} + \frac{\psi^{(2)}(x, y, t)}{r^2} + \dots$$

$$A_t = \mu - \rho/r$$

hep-th/9905104v2

1309.1439

Science 2013

$$\begin{aligned} \epsilon(t) &= t/\tau_Q & t_i &= (1 - T_i/T_c)\tau_Q \\ t &\in (t_i, t_f) & t_f &= (1 - T_f/T_c)\tau_Q \end{aligned}$$

$$\langle O_2 \rangle \sim \psi_2$$

Stochastic driving

$$\psi^{(1)} = \varphi(t, x)$$

$$\langle \varphi^*(t, x) \varphi(t', x') \rangle = \xi \delta(t - t') \delta(x - x')$$

Field theory:

$$\xi(T, \nu)$$

Quantum/thermal fluctuations

Gravity:

$$\xi \propto 1/N^2$$

Hawking radiation

Predictions

Mean field critical exponents

Slow quenches:

$$C(t, r) \sim |\psi|^2(t) e^{-\frac{r^2}{2\ell_{\text{co}}^2(t)}}, \quad |\psi|^2(t) \sim \tilde{\epsilon} t_{\text{freeze}} \bar{t} e^{a_2 \bar{t}^2}, \quad \ell_{\text{co}}(t) \sim \xi_{\text{freeze}} \sqrt{\bar{t}}$$

$$\frac{t_{\text{eq}}}{t_{\text{freeze}}} \sim \sqrt{\log \frac{N^2}{\sqrt{\tau_Q}}}$$

$$\rho \sim \frac{1}{\sqrt{\log \frac{N^2}{\sqrt{\tau_Q}}}} \rho_{\text{KZ}}$$

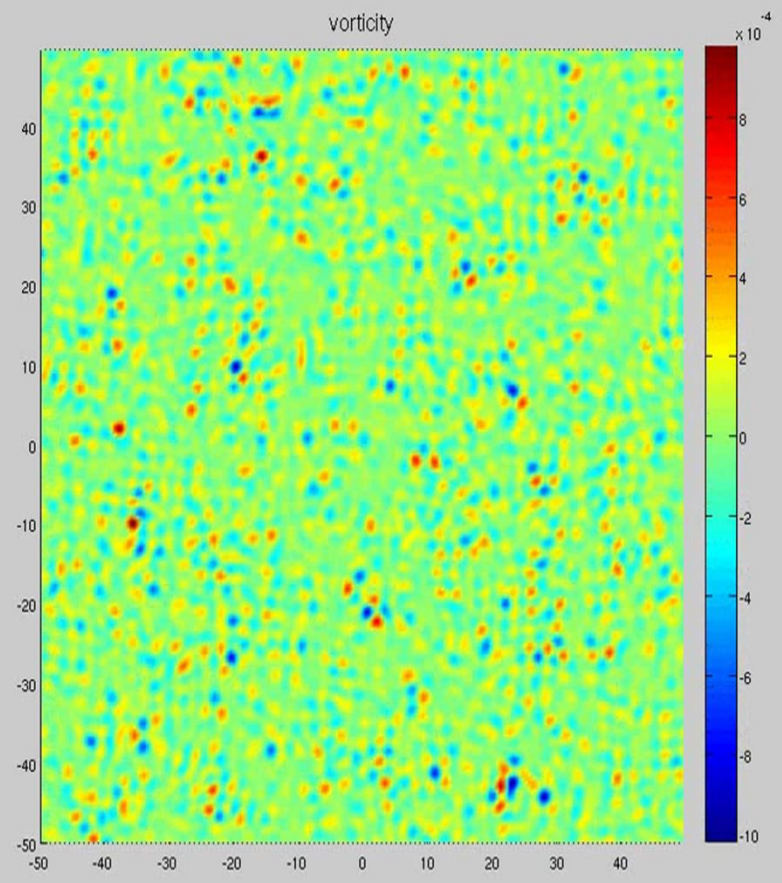
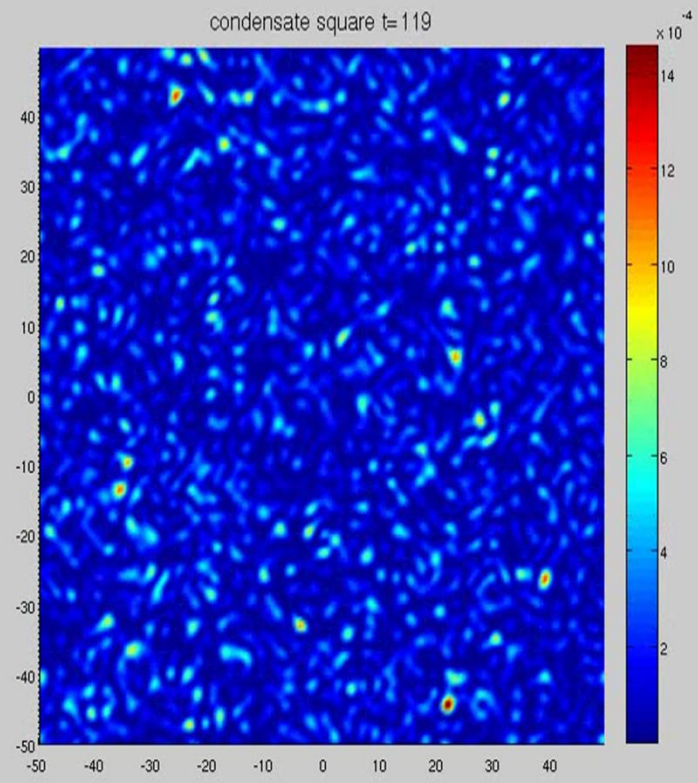
Fast quenches:

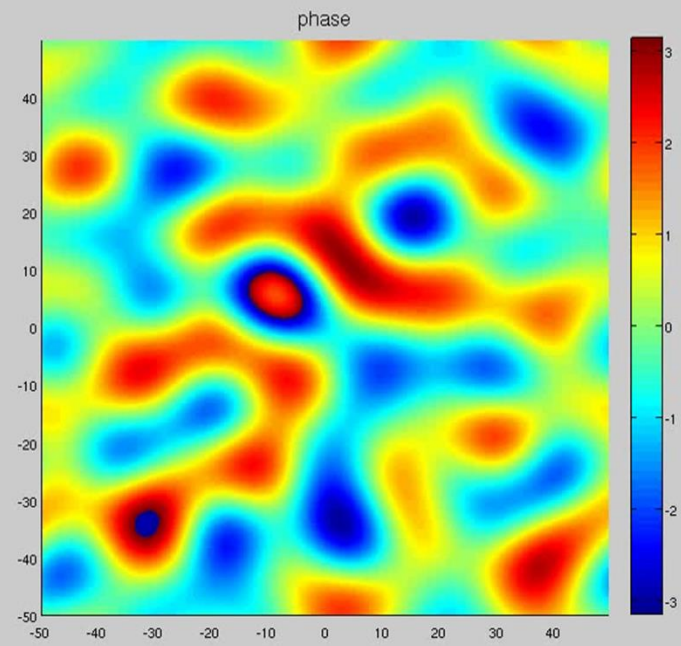
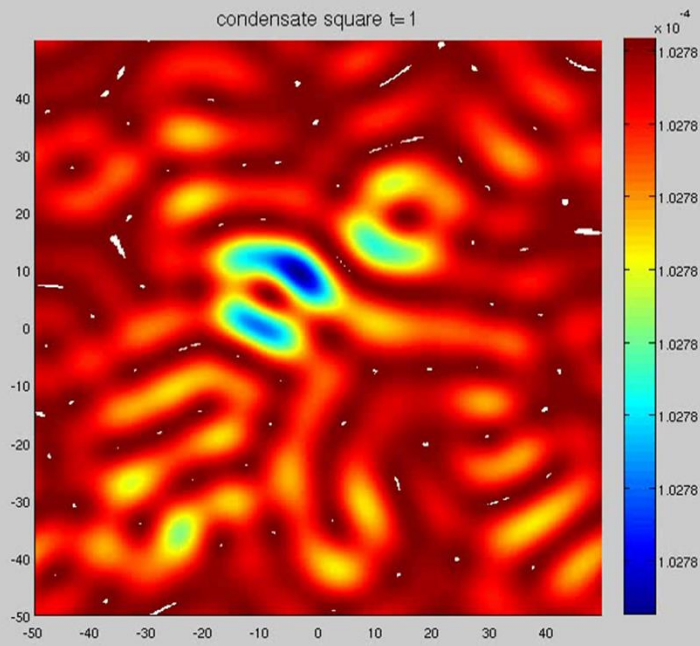
$$C(t, r) = |\psi|^2(t) e^{-\frac{r^2}{2\ell_{\text{co}}^2(t)}}, \quad |\psi|^2(t) \sim \zeta \exp [2b(t - t_{\text{freeze}})\epsilon_f]$$

$$\ell_{\text{co}}^2(t) = 4a(t - t_{\text{freeze}})$$

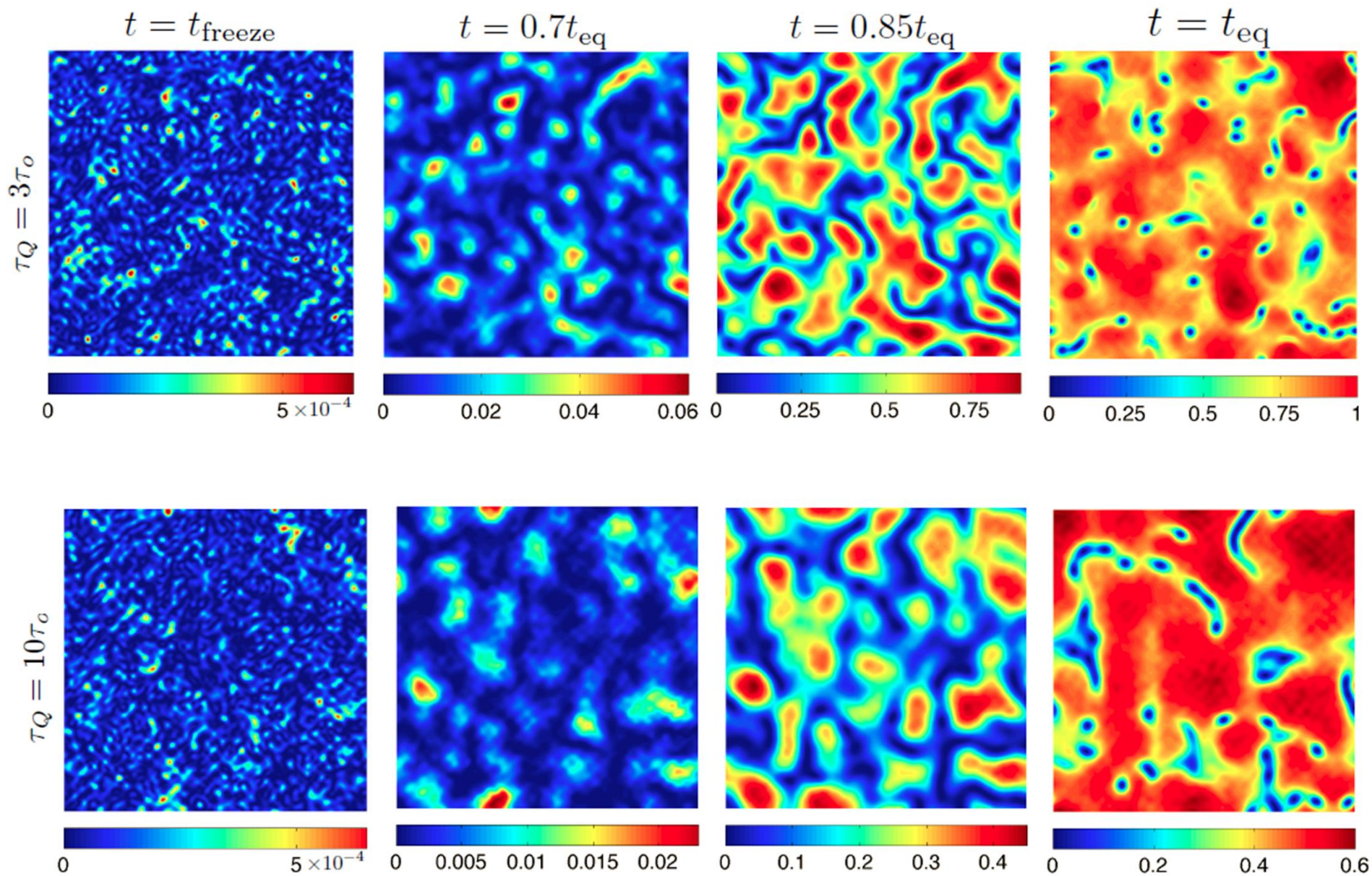
$$\rho \sim \frac{\epsilon_f}{\log \frac{N^2}{\epsilon_f}}$$

Movies!!





Slow

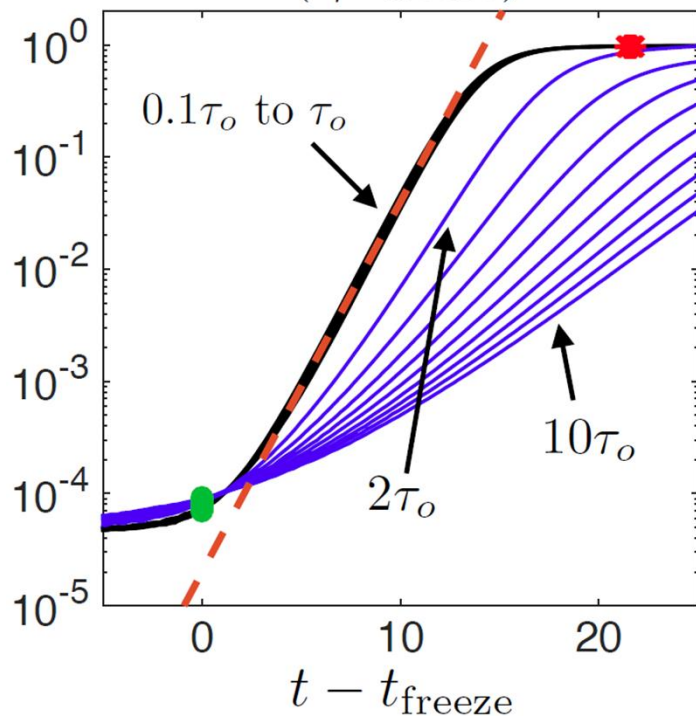
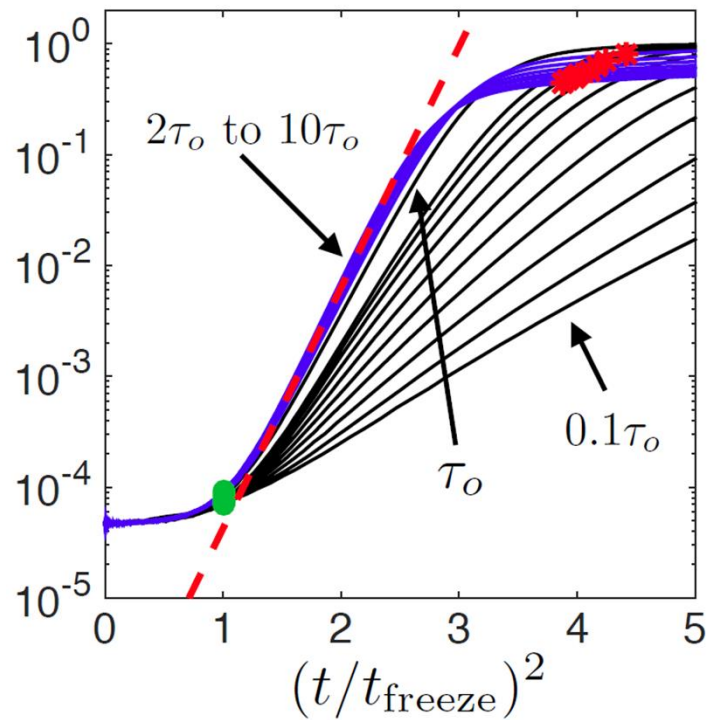
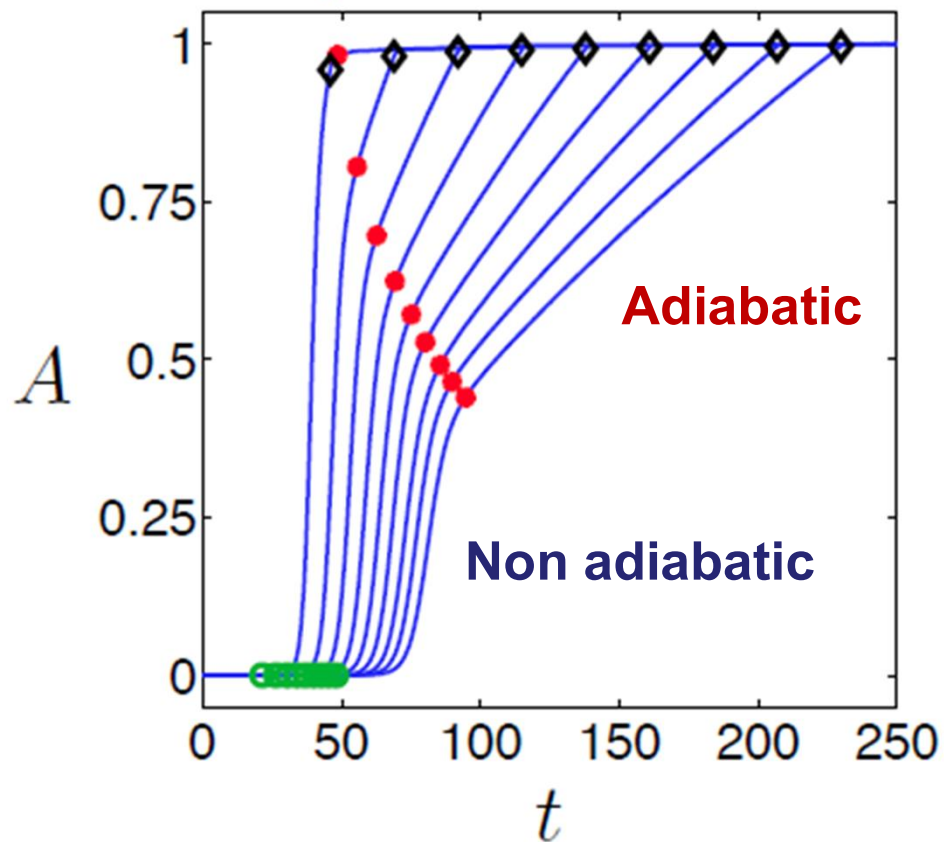


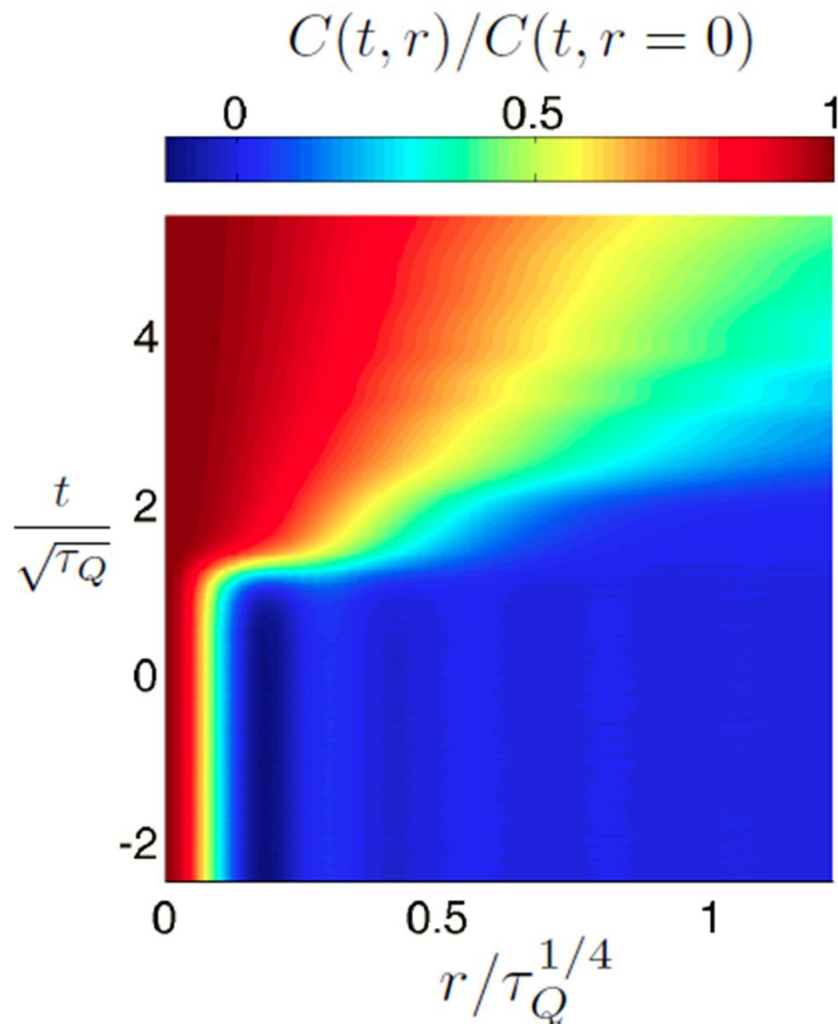
t_{eq} is the relevant scale

$$A(t) = \frac{1}{M} \sum_{i=1}^M \frac{a_i(t)}{a_i(\infty)}$$

$$a_i(t) \equiv \int d^2x |\psi_i(t, \mathbf{x})|^2$$

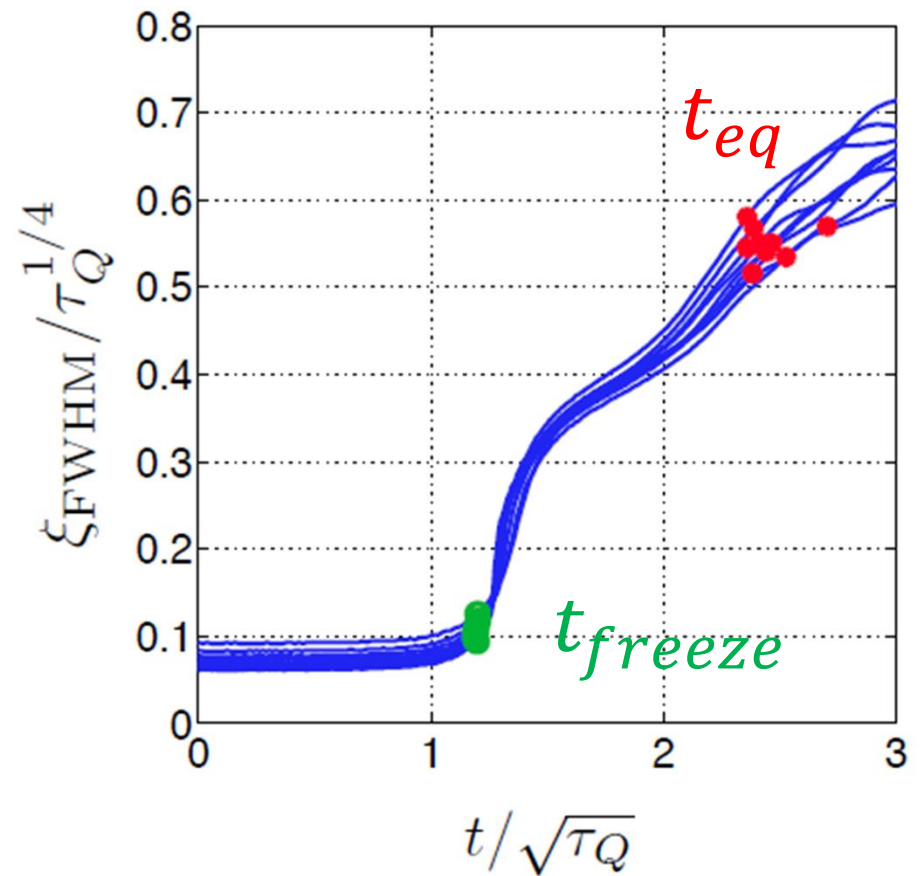
$$\tau_Q = 2\tau_o, \dots, 10\tau_o$$





Strong coarsening
 $t > t_{freeze}$

Full width half max of $C(t, r)$

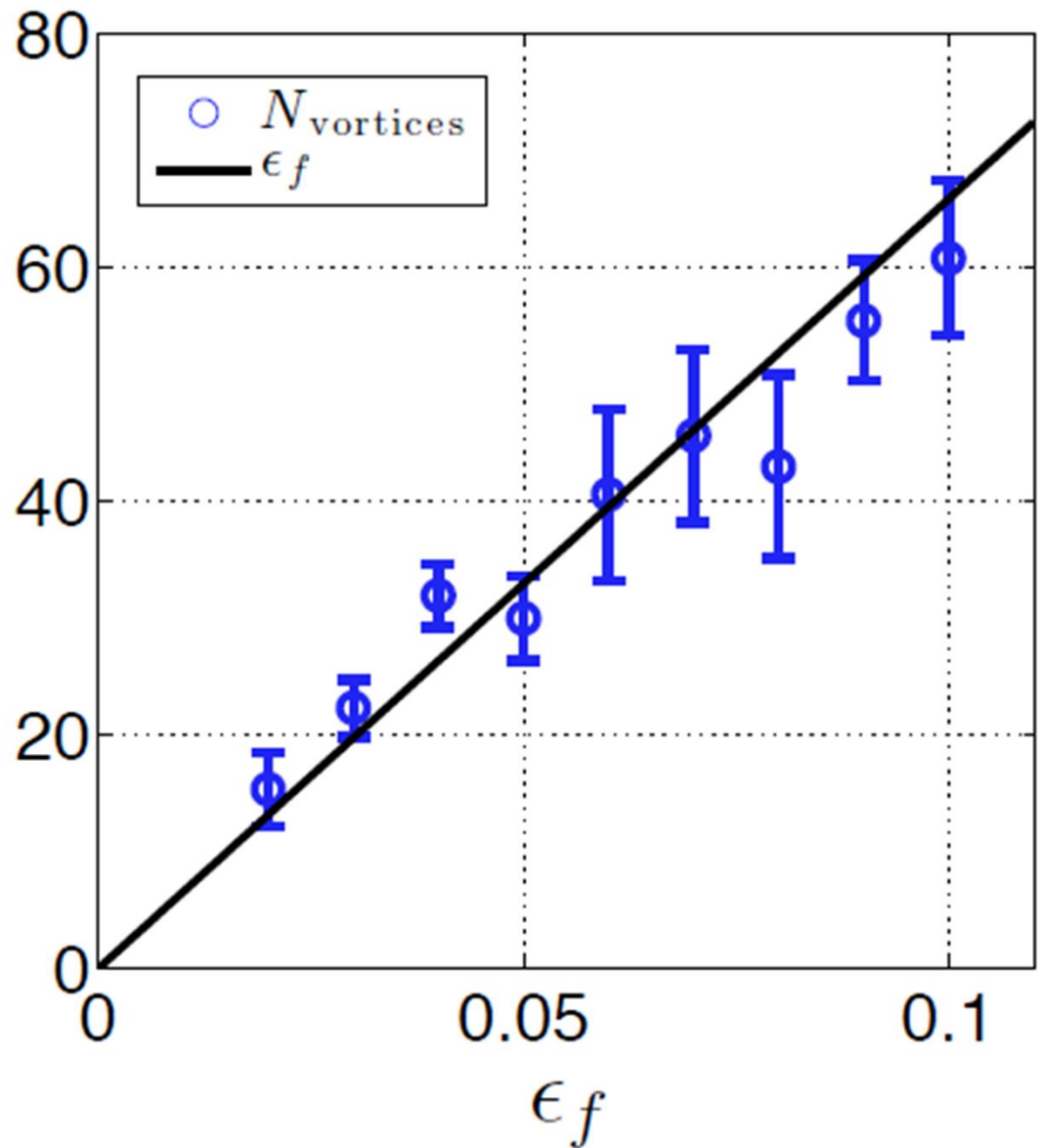


$$l_{co}(t) \sim \xi_{freeze} \sqrt{t/t_{freeze}}$$

Fast
quenches

$$\rho \sim \frac{\epsilon_f}{\log\left(\frac{N^2}{\epsilon_f}\right)}$$

$T < T_c$ dynamic
irrelevant



Slow

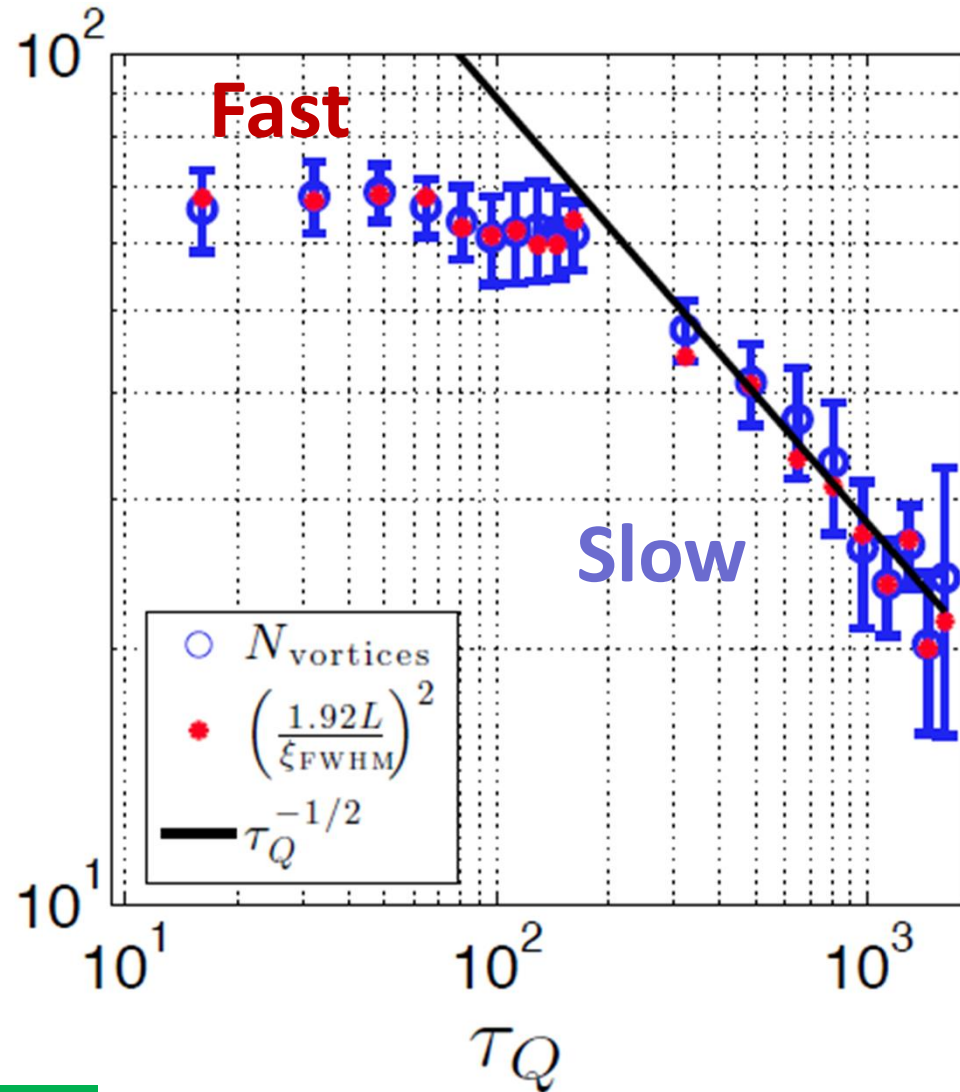
$$\rho \sim \frac{\rho_{KZ}}{(\log(N^2/\tau_Q^{1/2}))^{1/2}}$$

Fast

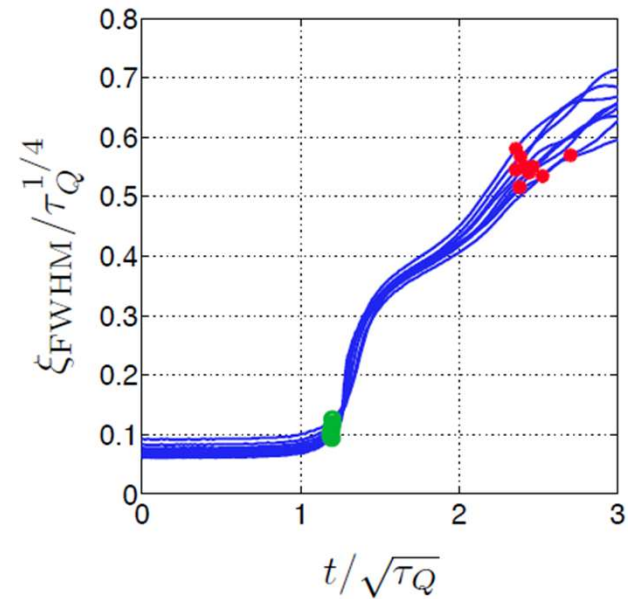
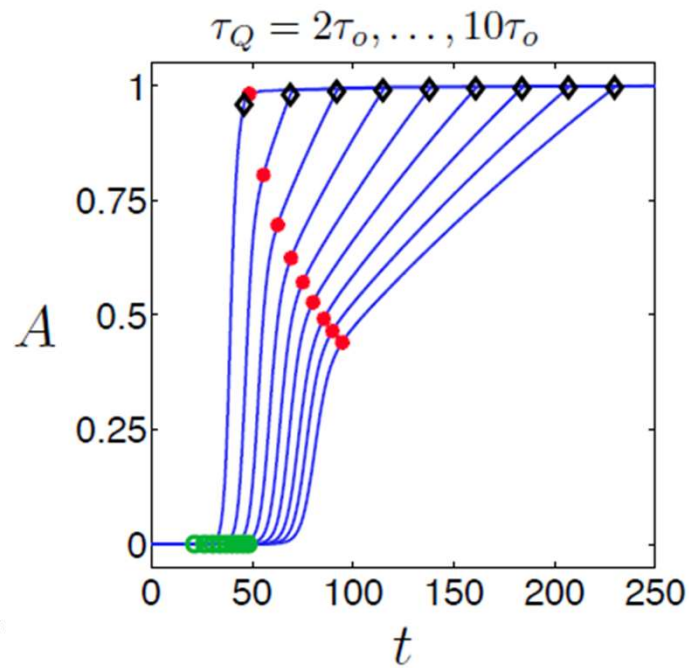
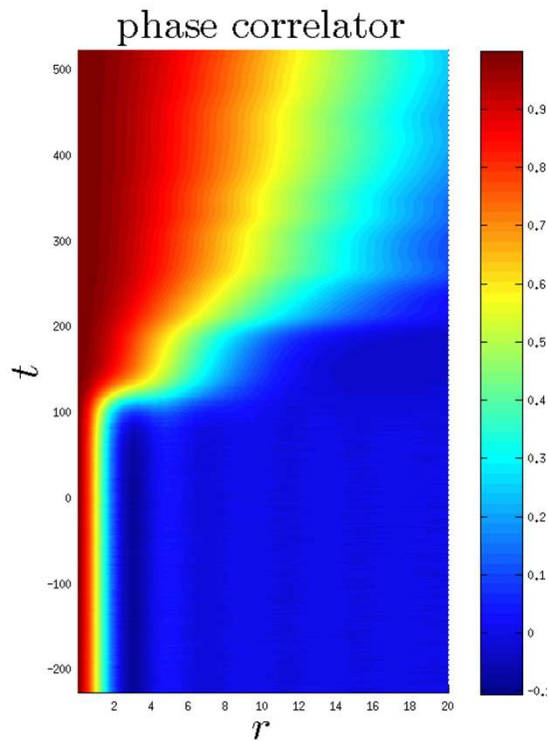
$$\rho \sim \frac{\epsilon_f}{\log\left(\frac{N^2}{\epsilon_f}\right)}$$

Relevant for ^4He ?

$$t_{\text{eq}} \sim [\log \tau_Q]^{1/(1+\nu z)} t_{\text{freeze}}$$



~25 times less defects than KZ prediction!!



time



Freezing

Condensate formation

Defect generation

Phase coherence ?

Physics beyond Kibble-Zurek

Novel dynamical region

$$t_{\text{eq}} > t > t_{\text{freeze}}$$

Holographic duality
helpful to discover and
model this region

More efficient than
SGPE?

NEXT

Cracking thermalization?

^4He ?

Cosmology

Vortex physics

BKT transition

Thanks!